

## SOME SOLUTIONS OF THE WAVE EQUATION

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1. For rectilinear boundaries a number of solutions of the two-dimensional wave equation

$$\nabla^2 \phi = c^2 \frac{\partial^2 \phi}{\partial t^2} \quad (1)$$

are known, in which the boundary condition satisfied by  $\phi$  is

$$\phi = 0, \quad (2.1)$$

or

$$\frac{\partial \phi}{\partial n} = 0, \quad (2.2)$$

$\partial n$  denoting a normal element drawn to the boundary. In the transverse vibrations of membranes the condition (2.1) amounts to a fixed edge, and in Hydro-dynamics of inviscid liquids (2.2) is to be satisfied over a fixed boundary. Some such solutions satisfying (2.1) have been given by B. R. Seth.<sup>1</sup> For hexagonal and equilateral boundaries\* incomplete solutions have been obtained by D. G. Christopherson<sup>2</sup> and B. B. Sen.<sup>3</sup> Christopherson's solution has been completed by P. N. Sharma.<sup>4</sup>

2. When the boundary condition is nonhomogeneous solutions of the wave equation for rectilinear boundaries have not received much attention, though formal solutions can be obtained through the Green's function. In a recent investigation of some problems in the bending of elastic plates it was required to determine solutions of (1) satisfying the boundary condition

$$\phi = k, \text{ a constant.} \quad (2.3)$$

It is proposed to give some of these solutions. It is found that for equilateral, some rhombus and pentagonal boundaries the solution can be obtained in a finite number of trigonometrical terms.

3. Assuming that  $\theta \propto \cos(pt/c)$  we can rewrite (1) as

$$\nabla^2 \phi + p^2 \phi = 0. \quad (3.1)$$

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\* The complete solution for an equilateral boundary was originally obtained by Lamé in *Leçons sur l'Elasticité*.

A typical solution of (3.1) is

$$\phi = \frac{\sin}{\cos} (p \sqrt{\cos \theta}) \frac{\sin}{\cos} (py \sin \phi). \quad (3.2)$$

It may be noted that when applied to the vibrations of a membrane the condition (2.3) means that its edge is vibrating harmonically.

4. For a rectangular section with sides  $x = \pm a$ ,  $y = \pm b$ , we see that  $k \cos px \sec pa$  satisfies the boundary condition (2.3) over  $x = \pm a$ . Using the Fourier's expansion for  $\cos px$  which vanishes for  $x = \pm a$ , we get

$$\phi = k \left[ \cos px \sec pa - \sum A_n \sec b \sqrt{p^2 - \alpha_n^2} \cos \alpha_n x \cos \sqrt{p^2 - \alpha_n^2} y \right], \quad (4)$$

where  $\alpha_n = (2n + 1) \pi / 2a$ ,

$$A_n = \frac{2(-1)^{n+1} k p^2}{a \alpha_n (p^2 - \alpha_n^2)}.$$

For a right-angled isosceles triangular section with sides  $x = a$ ,  $y = a$ ,  $y + x = 0$ , we see that the functions to be used for the individual sides can be taken to be of the form  $\sin px$ ,  $\sin py$  and  $\cos \frac{p}{\sqrt{2}} (x + y)$ , respectively.

On account of the symmetry which exists about the line  $y = x$  we now take two Fourier series which are such that their sum vanishes over  $y = -x$  and in which terms can be obtained from one another by interchanging  $x$  and  $y$ . Thus we get

$$\begin{aligned} \phi = & k \cos \frac{p}{\sqrt{2}} (x + y) + \frac{1}{2} k (1 - \cos pa \sqrt{2}) \left[ \frac{\sin px}{\sin pa} + \frac{\sin py}{\sin pa} \right] \\ & - \sum A_n \operatorname{cosec} a \sqrt{p^2 - \alpha_n^2} \left[ \cos \alpha_n x \sin \sqrt{p^2 - \alpha_n^2} y + \cos \alpha_n y \sin \sqrt{p^2 - \alpha_n^2} x \right] \\ & - \sum B_n \sec a \sqrt{p^2 - \beta_n^2} \left[ \sin \beta_n x \cos \sqrt{p^2 - \beta_n^2} y + \sin \alpha_n y \cos \sqrt{p^2 - \beta_n^2} x \right] \end{aligned} \quad (5)$$

where  $\alpha_n = (2n + 1) \pi / 2a$ ,  $\beta_n = n\pi/a$ , and

$$A_n = \frac{(-1)^{n+1} k p^2 \cos^2 \frac{1}{2} pa \sqrt{2}}{a \alpha_n (\frac{1}{2} p^2 - \alpha_n^2)}, \quad B_n = \frac{(-1)^{n+1} k \beta_n p^2 \sin^2 \frac{1}{2} pa \sqrt{2}}{a (\frac{1}{2} p^2 - \beta_n^2) (p^2 - \beta_n^2)}.$$

For an equilateral section with sides  $y = a$ ,  $y = \pm x \sqrt{3}$ , we should take functions of the type  $\cos py$ ,  $\sin py$ ,  $\cos \frac{1}{2} px \sqrt{3} \cos \frac{1}{2} py$  and  $\cos \frac{1}{2} px \sqrt{3} \sin \frac{1}{2} py$ . It is found that no Fourier expansions are required, and we get

$$\begin{aligned} \phi = & k [2 \cos \frac{1}{2} py \cos \frac{1}{2} px \sqrt{3} - \cos py] \\ & - k \cot \frac{1}{2} pa [2 \sin \frac{1}{2} py \cos \frac{1}{2} px \sqrt{3} - \sin py]. \end{aligned} \quad (6)$$

The result for a rhombus with sides  $y = \pm x \sqrt{3} \pm 4\pi n/p$ ,  $n$  being an integer, can be readily deduced from (6) and we get

$$\phi = k [2 \cos \frac{1}{2} py \cos \frac{1}{2} px \sqrt{3} - \cos py.] \tag{7}$$

In like manner for a pentagonal section with sides  $y = a, y = \pm x \sqrt{3} \pm 4\pi n/p$ , the solution given in (6) holds good.

REFERENCES

1. Seth, B. R. .. *Proc. Ind. Acad. Sci.*, 1940, **12A**, 487-90.  
..... .. *Loc. cit.*, 1941, **13**, 390-94.
2. Christopherson, D. G. .. *Quart. J. of Math. (Oxford)*, 1940, **11**, 65.
3. Sen, B. B. .. *Bull. Cal. Math. Soc.*, 1934, **26**, 69.
4. Sharma, P. N. .. *Mathematics Student*, 1946, **14**, 63-64.