

LATTICE-THEORY OF ALKALINE EARTH CARBONATES.

Part IV. Elasticity Constants of Calcite.

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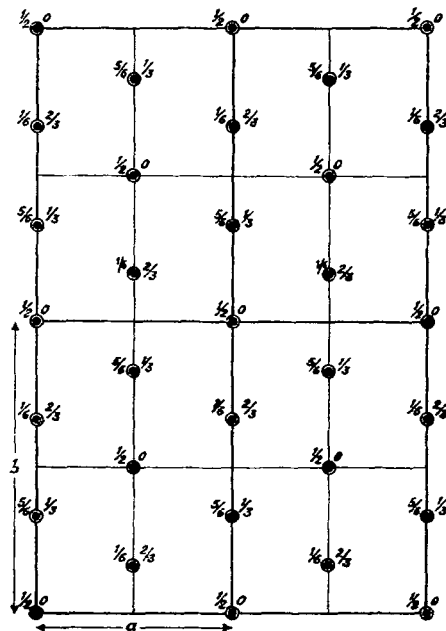
(Communicated by Sir C. V. Raman, kt., F.R.S., N.I.)

Introduction.

THE three previous papers,¹ under the same title, will be quoted herein as I, II and III respectively. In this part we proceed with the calculation of such elasticity constants of calcite, as do not essentially depend on ionic deformations. The method is exactly the same as used in II for the calculation of the elasticity constants of aragonite. Thus c_{ij} consists of two parts c_{ij}^e and c_{ij}^r due, respectively, to the electro-static and the repulsive forces. These two parts will be calculated separately.

§1. Systems of Parallel Neutral Planes.

Fig. 1 gives the projection of the crystal on x - y plane and Fig. 2 on the



The numbers on the right refer to Calcium \circ
 " " " left " \circ
 CO_3

Fig. 1.

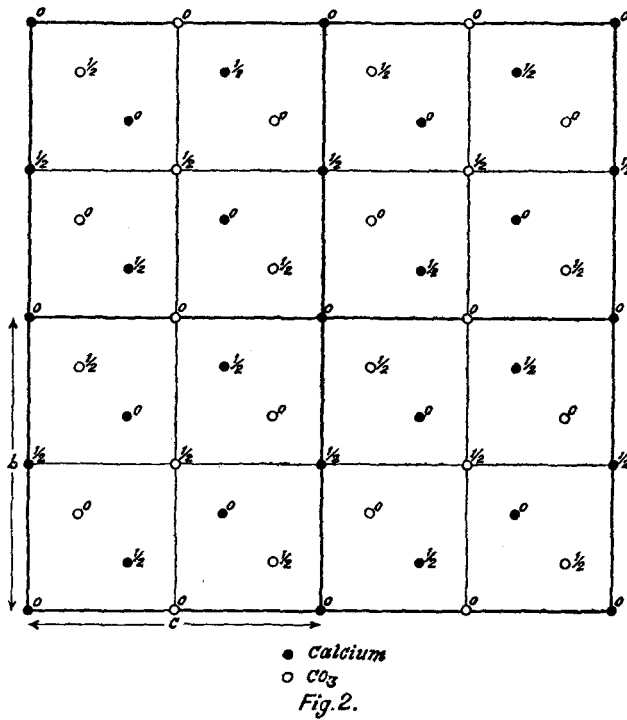
¹ B. Y. Oke, *Proc. Ind. Acad. Sci.*, 1936, 4, 1, 514 and 525.

y - z plane. The numbers in the figures denote the "heights" as multiples c and a [III (1, 1)], above the planes $z = 0$ and $x = 0$ respectively.

We take the following three systems of parallel neutral planes :

- (1) all planes perpendicular to the x -axis,
 - (2) all planes perpendicular to the y -axis,
- and (3) all planes perpendicular to the line in the plane yoz making an angle θ with the z -axis, where θ is given by

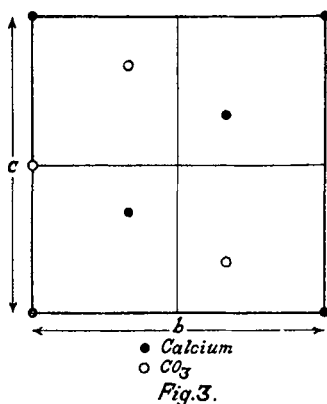
$$\tan \theta = \frac{b}{c} \quad (1, 1)$$



First System.—In this the distance between consecutive planes is given by

$$d = \frac{a}{2} \quad (1, 2)$$

In a unit cell of each neutral plane, there are six lattice-points. Fig. 3 shows the arrangement of the ions in the planes at $x = 0, \pm a, \pm 2a, \dots, \pm na, \dots$



The y and z co-ordinates of the 6 lattice-points are given by

$$\left. \begin{array}{ll}
 k = 1, \text{ calcium ion } \left(0, 0\right); & k = 4, \text{ CO}_3 \text{ ion } \left(0, \frac{c}{2}\right) \\
 k = 2, \text{ ,, ,, } \left(\frac{b}{3}, \frac{c}{3}\right); & k = 5, \text{ ,, ,, } \left(\frac{b}{3}, \frac{5c}{6}\right) \\
 k = 3, \text{ ,, ,, } \left(\frac{2b}{3}, \frac{2c}{3}\right); & k = 6, \text{ ,, ,, } \left(\frac{2b}{3}, \frac{c}{6}\right)
 \end{array} \right\} (1, 3)$$

The ionic arrangement in the planes at

$$x = \pm \frac{a}{2}, \pm \frac{3a}{2}, \dots \pm \frac{2n+1}{2} a, \dots$$

is similar to this except that there is a displacement of the ions through a distance $\frac{b}{2}$ along Oy .

Second System.—In this the distance between consecutive planes is given by

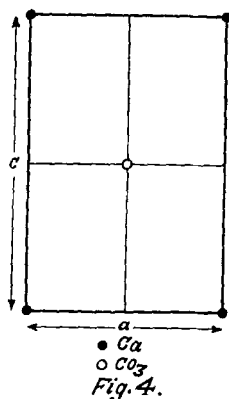
$$d = \frac{b}{6} \tag{1, 4}$$

The unit cell of each plane consists of only two lattice-points.

Fig. 4 shows the arrangement of the ions in the plane $y = 0$. The x and z co-ordinates of the lattice-points are

$$k = 1, \text{ calcium ion } (0, 0) \text{ and } k = 2, \text{ CO}_3 \text{ ion } \left(0, \frac{c}{2}\right) \tag{1, 5}$$

The successive planes are similar in the ionic arrangement except that each is displaced through distances $\frac{a}{2}$ and $\frac{2c}{3}$ along Ox and Oz respectively.



Third System.—In this the distance between consecutive planes is given by

$$d_{\theta} = \frac{bc}{2\sqrt{b^2 + c^2}} \quad (1, 6)$$

For the representation of this system we choose a system of co-ordinates $x_{\theta}, y_{\theta}, z_{\theta}$; such that Ox_{θ} coincides with Ox and Oz_{θ} is perpendicular to the system of parallel planes. The dimensions of a unit cell of the plane lattice are given by

$$\left. \begin{aligned} a_{\theta} &= a \\ b_{\theta} &= \frac{\sqrt{b^2 + c^2}}{3} \end{aligned} \right\} \quad (1, 7)$$

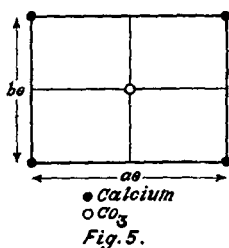


Fig. 5 shows the arrangement of the ions in the plane $Z_{\theta} = 0$. A unit cell of the lattice consists of two points only, whose x_{θ} and y_{θ} co-ordinates are given by

$$k = 1, \text{ calcium ion } (0, 0); \quad k = 2, \text{ CO}_3 \text{ ion } \left(\frac{a_{\theta}}{2}, \frac{b_{\theta}}{2}\right); \quad (1, 8)$$

The ionic arrangement in successive planes is similar except that each is displaced through distances $\frac{a_{\theta}}{2}$ and $\frac{b_{\theta}}{4}$ along Ox_{θ} and Oy_{θ} respectively.

Denoting the elasticity constants in directions perpendicular to these three systems by c_{11} , c_{22} and c_{33}^θ respectively ; we have as in II (1, 9)

$$\left. \begin{aligned} c_{11} &= c_{11}^e + c_{11}^\theta \\ c_{22} &= c_{22}^e + c_{22}^\theta \\ \text{and } c_{33}^\theta &= c_{33}^e + c_{33}^\theta \end{aligned} \right\} \quad (1, 9)$$

§2. *Evaluation of c_{11}^e , c_{22}^e and c_{33}^e .*

As derived in II (2, 12) we have

$$c_{11}^\theta = \frac{4\pi^2}{b^2c^2} 4e^2 \sum_{j=1}^{\infty} \left\{ \frac{\alpha_{l,m} |jd|}{2\pi} e^{-\alpha_{l,m} |jd|} \sum_k [F_j^{l,m}(k)] \right\} \quad (2, 1)$$

where
$$\alpha_{l,m} = 2\pi \left| \frac{l}{b} + \frac{m}{c} \right| \quad (2, 2)$$

The function $F_j^{l,m}(k)$ depends, as shown in II [§2, refer equation (2, 5) — (2, 10)] on the ionic arrangement of the planes.

The functions $\alpha_{l,m}$ and $F_j^{l,m}$ will be different in the 3 cases.

In the first system, only two adjoining planes need to be considered. The bracketed expression on summation is found to be

$$- 0.267 ; \quad (2, 3)$$

In the second case, only five adjoining planes have to be considered.

$$- 0.085 ; \quad (2, 4)$$

In the third case, the summation yields the value

$$- 0.1048 ; \quad (2, 5)$$

Substituting these in the proper formula we get

$$c_{11}^e = \frac{16\pi^2e^2}{b^2c^2} (- 0.267) = - 0.176 \times 10^{12} \text{ dynes/cm.}^2 ; \quad (2, 6)$$

$$c_{22}^e = \frac{16\pi^2e^2}{c^2a^2} (- 0.085) = - 0.169 \times 10^{12} \text{ dynes/cm.}^2 ; \quad (2, 7)$$

and
$$c_{33}^{\theta} = \frac{16\pi^2e^2}{a_\theta^2b_\theta^2} (- 0.1048) = - 0.916 \times 10^{12} \text{ dynes/cm.}^2 ; \quad (2, 8)$$

§3. *The Nature of the Repulsive Forces.*

The change in the potential due to the repulsive forces is given by III (2, 1)

$$2 \frac{\beta}{\delta^n} ; \quad (3, 1)$$

where, as in I (3, 4)

$$\beta = \frac{ae^2}{n} \delta^{n-1} ; \quad (3, 2)$$

δ is defined in III (1, 24).

The force between two dissimilar ions is $\frac{b_{12}}{r^n}$ and $\frac{b_{11}}{r^n}$ or $\frac{b_{22}}{r^n}$ between two similar ions. Let

$$b_{11} = b_{22} = -kb_{12}; \quad (3, 3)$$

As in II (§3) we consider only the close neighbours of any ion; in fact we take only such neighbours of which the distances from the ions are less than δ , defined in III (1, 24). Further from III (1, 1) we get

$$b = 8.63 \quad (3, 4)$$

Thus b may be taken roughly equal to c . Putting

$$\left. \begin{aligned} r_1^2 &= \frac{b^2}{4}; & s_1^2 &= a^2; \\ r_2^2 &= b^2 \left(\frac{1}{9} + \frac{1}{3^{\frac{1}{6}}} \right); & s_2^2 &= b^2 \left(\frac{1}{9} + \frac{1}{9} \right); \\ r_3^2 &= b^2 \left(\frac{1}{9} + \frac{1}{3^{\frac{1}{6}}} \right); & s_3^2 &= b^2 \left(\frac{1}{9} + \frac{1}{9} \right); \\ r_4^2 &= b^2 \left(\frac{1}{9} + \frac{1}{4} \right); \end{aligned} \right\} \quad (3, 5)$$

The potential due to the repulsive forces is given by

$$V_{\Delta} = 12 b_{12} \left\{ \frac{2}{r_1^n} + \frac{6}{r_2^n} + \frac{6}{r_3^n} + \frac{12}{r_4^n} - k \left(\frac{6}{s_1^n} + \frac{6}{s_2^n} + \frac{12}{s_3^n} \right) \right\} \quad (3, 6)$$

n has been found in III (Table II) to be 6.5.

Expressing the distances in terms of δ

$$V_{\Delta} = \frac{12 b_{12}}{\delta^{6.5}} \{1275.6 - 332 k\} \quad (3, 6')$$

Comparing this with (3, 1)

$$\frac{2 a e^2 \delta^{5.5}}{6.5} = 12 b_{12} (1275.6 - 332 k) \quad (3, 7)$$

whence

$$b_{12} = \frac{a e^2 \delta^{5.5}}{6 \times 6.5 (1275.6 - 332 k)} \quad (3, 8)$$

Substituting this in (3, 6)

$$\begin{aligned} V_{\Delta} &= 2 \frac{a e^2 \delta^{5.5}}{6.5 (1275.6 - 332 k)} \left\{ \frac{2}{r_1^{6.5}} + \frac{6}{r_2^{6.5}} + \frac{6}{r_3^{6.5}} + \frac{12}{r_4^{6.5}} \right. \\ &\quad \left. - k \left(\frac{6}{s_1^{6.5}} + \frac{6}{s_2^{6.5}} + \frac{12}{s_3^{6.5}} \right) \right\} \end{aligned} \quad (3, 9)$$

V , the potential per unit volume, is given by

$$V = \frac{V_{\Delta}}{\delta^3} \quad (3, 10)$$

For c_{11}' , c_{22}' and c_{33}'' we take formula similar to those given in II (3, 11) and (3, 12).

The value of k is determined, as in the case of aragonite, by comparison with the observed value of c_{11} .

§4. *Evaluation of k and c_{33}^{θ} .*

Carrying out the differentiations and then substituting the numerical values of the r 's and s 's and their components in the respective directions we get :—

$$c_{11}^r = \frac{\alpha e^2}{\delta^4} \frac{1870 \cdot 3 - 363 \cdot 8 k}{1275 \cdot 6 - 332 k} \quad (4, 1)$$

$$c_{22}^r = \frac{\alpha e^2}{\delta^4} \frac{1870 \cdot 3 - 363 \cdot 8 k}{1275 \cdot 6 - 332 k} \quad (4, 2)$$

$$c_{33}^{r\theta} = \frac{\alpha e^2}{\delta^4} \frac{2478 \cdot 6 - 455 \cdot 1 k}{1275 \cdot 6 - 332 k} \quad (4, 3)$$

From (4, 1) and (4, 2)

$$c_{11}^r = c_{22}^r$$

and since from (2, 6) and (2,7) we have c_{11}^e and c_{22}^e approximately equal, we obtain by calculation the observed result

$$c_{11} = c_{22} \quad (4, 4)$$

The value of c_{11} , as experimentally determined by Voigt,² is $1 \cdot 37 \times 10^{12}$ dynes/cm.²

Hence from (1, 9) and (2, 6)

$$c_{11}^r = 1 \cdot 546 \times 10^{12} \text{ dynes/cm.}^2 \quad (4, 5)$$

$$\therefore \frac{\alpha e^2}{\delta^4} \frac{1920 \cdot 8 - 327 \cdot 8 k}{1215 \cdot 6 - 332 k} = 1 \cdot 546 \times 10^{12} \text{ dynes/cm.}^2 \quad (4, 6)$$

Substituting the proper values of e and δ , and the value of α from III (1, 28) we get

$$1 \cdot 03 \frac{1870 \cdot 3 - 363 \cdot 8 k}{1275 \cdot 6 - 332 k} = 1 \cdot 546 \quad (4, 7)$$

whence

$$k = 0 \cdot 353 \quad (4, 8)$$

Substituting this value of k in (4, 3)

$$c_{33}^{r\theta} = 2 \cdot 06 \cdot \times 10^{12} \text{ dynes/cm.}^2 \quad (4, 9)$$

The value of the elasticity constant in the direction Oz_{θ} is not directly measured. As in II (5, 1) this can be obtained with the transformation formula given by Voigt. We get

$$c_{33}^{\theta} = 1 \cdot 109 \times 10^{12} \text{ dynes/cm.}^2 \quad (4, 10)$$

Substituting from (4, 9) and (2, 8) in (1, 9) the calculated value is given by

$$c_{33}^{\theta} = 1 \cdot 148 \times 10^{12} \text{ dynes/cm.}^2 \quad (4, 11)$$

The agreement between the two values is quite satisfactory.

² W. Voigt, *Lehrbuch der Kristallphysik*, p. 754.

Summary.

The electro-static parts of the elasticity constants c_{11} , c_{22} and c_{33}^{θ} are calculated by the same method as used for aragonite. The nature of the 'repulsive' forces is further completely determined by comparison with the observed value of c_{11} . The value of c_{33}^{θ} is then determined and found to agree well with the experimental value.