

Thermal effect and equation of state in solids

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MS received 15 June 1998

Abstract. In the present paper, an attempt has been made to include the thermal effect into an isothermal equation of state using the Debye approximation representing volume as a function of pressure. The calculations are done in case of NaCl, CsCl, Mo, and W. The present calculations are in good agreement with the reported results.

Keywords. Thermal effect; equation of state; Debye approximation.

PACS Nos 64.10; 64.30; 62.50; 62.20

1. Introduction

Two types of isothermal equation of state (EOS) are available in the literature. In one, the pressure is expressed as a function of volume, i.e. $P = f(V)$. The second in which the volume is expressed as a function of pressure, i.e., $V = f(P)$. The thermal effect into the first type of EOS due to the Debye approximation is very well known [1,2], somehow no attempt has been made to include the thermal effect due to the Debye approximation into the second type of EOS. The possible reason may be that (i) the pressure dependence of γ is not known and (ii) the pressure dependence of the Debye temperature, Θ , is not known. But at present the pressure dependence of γ and Θ are available in the literature. Therefore, the aim of the present paper is to include the thermal effect into a second type of isothermal EOS to make it as temperature dependent EOS.

2. Theory

An EOS expressing the volume in terms of pressure really corresponds to the experimental situation and hence such EOS is considered to be better than the other form of EOS, i.e., expressing pressure in terms of volume. This goal is achieved by differentiating free energy [3] of the solids with respect to pressure considering the Debye temperature a pressure dependent, $\Theta(P)$. The EOS of a solid, in Debye approximation, has been found as

$$V(P, T) = V_P + \frac{\gamma(P, T)}{B_T(P, T)} \left[\frac{9}{8} nR\Theta + 3nRTD(X) \right] \quad (1)$$

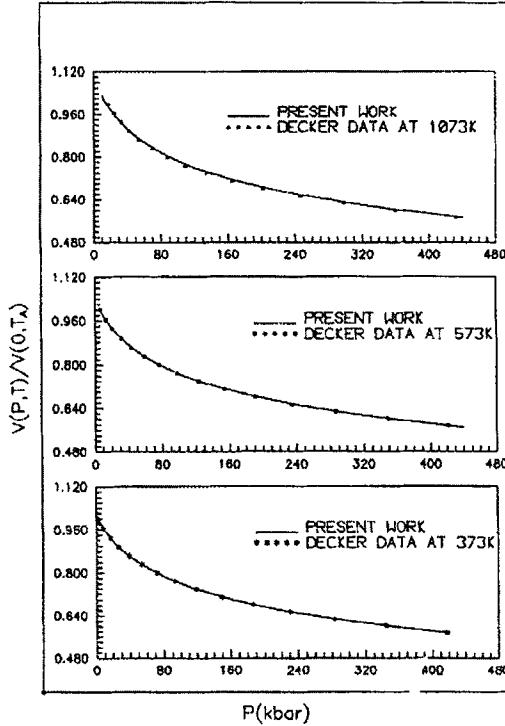


Figure 1. Comparison of $V(P, T)/V(0, T_A)$ with pressure at different temperatures for NaCl.

where

$$D(X) = \frac{3}{X^3} \int_0^X \frac{X^3}{e^X - 1} dX \quad \text{and} \quad X = \frac{\Theta}{T}$$

V_P , Θ , n and γ denote, respectively, the volume in static lattice, the Debye temperature, the number of atoms in the unit cell and the thermodynamic Gruensien parameter. The second term, i.e., zero point volume can be neglected because of its smallness.

Equation (1) can also be written as

$$V(P, T) = V(P, T_A) + \frac{\gamma(P, T_A)}{B_T(P, T_A)} [E(\Theta_T) - E(\Theta_{T_A})] \quad (2)$$

In obtaining eq. (2), we make use of the following assumption:

$$\frac{\gamma(P, T)}{B_T(P, T)} \simeq \frac{\gamma(P, T_A)}{B_T(P, T_A)} \quad (3)$$

Expression for the variation of $\gamma(P, T_A)$ and $\Theta(P, T_A)$ with pressure have been taken from Kumari and Dass [4] and are given below

$$\gamma(P, T_A) = \gamma(0, T_A) - \mu P \quad (4)$$

$$\Theta(P, T_A) = \Theta(0, T_A) \left[1 + \frac{B'_T(0, T_A)}{B_T(0, T_A)} P \right]^{+\gamma(0, T_A)/B'_T(0, T_A)} \quad (5)$$

where μ has been taken as an adjustable parameter.

Table 1. Input data.

| Solids | NaCl | CsCl | Mo | W |
|--|--------------|--------------|--------------|--------------|
| $B_T(0, T_A)$ (kbar) | 240.14 [7] | 172.51 [7] | 2709.2 [5] | 3138.04 [5] |
| $B'_T(0, T_A)$ | 4.54 [7] | 4.96 [7] | 3.58 [5] | 3.68 [5] |
| $\gamma(0, T_A)$ | 1.6275 [6] | 1.99 [6] | 1.58 [2] | 1.65 [2] |
| $\Theta(0, T_A)$ (K) | 279.0 [6] | 151.0 [6] | 411.0 [2] | 340.0 [2] |
| $\mu(X10^{-3} \text{ kbar}^{-1})$ | 2.2 | 1.2 | 0.175 | 0.175 |
| $V(0, T_A)$ (cm ³ /mole) | 26.99 | 42.18 | 9.39 | 9.55 |
| Temperature range (K) | 273.0–1073.0 | 273.0–1073.0 | 293.0–7605.0 | 293.0–6905.0 |
| Pressure range (kbar) | 0.0–320.0 | 0.0–432.0 | 0.0–3000.0 | 0.0–3000.0 |

Table 2. Comparison of density as a function of pressure and temperature for Mo and W.

| P (Gpa) | Mo | | | W | | |
|-----------|---------|--|--|---------|--|--|
| | T (K) | ρ (gm/cm ³) From [2] | ρ (gm/cm ³) Calculated | T (K) | ρ (gm/cm ³) From [2] | ρ (gm/cm ³) Calculated |
| 0.0 | 293.0 | 10.215 | 10.215 | 293.0 | 19.256 | 19.256 |
| 10.0 | 311.0 | 10.577 | 10.573 | 309.0 | 19.846 | 19.841 |
| 20.0 | 338.0 | 10.906 | 10.901 | 332.0 | 20.388 | 20.383 |
| 30.0 | 378.0 | 11.208 | 11.205 | 336.0 | 20.891 | 20.887 |
| 40.0 | 435.0 | 11.489 | 11.487 | 413.0 | 21.360 | 21.358 |
| 50.0 | 510.0 | 11.752 | 11.751 | 477.0 | 21.802 | 21.801 |
| 60.0 | 604.0 | 11.999 | 11.998 | 556.0 | 22.219 | 22.219 |
| 70.0 | 716.0 | 12.233 | 12.232 | 652.0 | 22.614 | 22.615 |
| 80.0 | 847.0 | 12.455 | 12.454 | 765.0 | 22.991 | 22.991 |
| 90.0 | 996.0 | 12.666 | 12.665 | 894.0 | 23.351 | 23.350 |
| 100.0 | 1163.0 | 12.868 | 12.866 | 1040.0 | 23.695 | 23.694 |
| 110.0 | 1347.0 | 13.062 | 13.058 | 1202.0 | 24.026 | 24.023 |
| 120.0 | 1549.0 | 13.248 | 13.243 | 1380.0 | 24.345 | 24.340 |
| 130.0 | 1767.0 | 13.427 | 13.420 | 1573.0 | 24.652 | 24.645 |
| 140.0 | 2002.0 | 13.599 | 13.592 | 1582.0 | 24.948 | 24.939 |
| 150.0 | 2252.0 | 13.766 | 13.757 | 2006.0 | 25.235 | 25.224 |
| 160.0 | 2517.0 | 13.927 | 13.917 | 2244.0 | 25.513 | 25.500 |
| 170.0 | 2798.0 | 14.083 | 14.072 | 2496.0 | 25.782 | 25.768 |
| 180.0 | 3092.0 | 14.235 | 14.223 | 2762.0 | 26.044 | 26.029 |
| 190.0 | 3401.0 | 14.382 | 14.370 | 3042.0 | 26.298 | 26.282 |
| 200.0 | 3723.0 | 14.525 | 14.513 | 3334.0 | 26.546 | 26.530 |
| 210.0 | 4058.0 | 14.664 | 14.652 | 3639.0 | 26.787 | 26.771 |
| 220.0 | 4406.0 | 14.800 | 14.789 | 3957.0 | 27.022 | 27.007 |
| 230.0 | 4766.0 | 14.932 | 14.922 | 4286.0 | 27.251 | 27.238 |
| 240.0 | 5139.0 | 15.061 | 15.053 | 4628.0 | 27.475 | 27.465 |
| 250.0 | 5523.0 | 15.187 | 15.181 | 4980.0 | 27.693 | 27.687 |
| 260.0 | 5918.0 | 15.310 | 15.307 | 5344.0 | 27.907 | 27.905 |
| 270.0 | 6324.0 | 15.430 | 15.430 | 5719.0 | 28.116 | 28.119 |
| 280.0 | 6741.0 | 15.548 | 15.552 | 6104.0 | 28.320 | 28.330 |
| 290.0 | 7168.0 | 15.663 | 15.672 | 6499.0 | 28.521 | 28.537 |
| 300.0 | 7605.0 | 15.776 | 15.790 | 6905.0 | 28.717 | 28.742 |

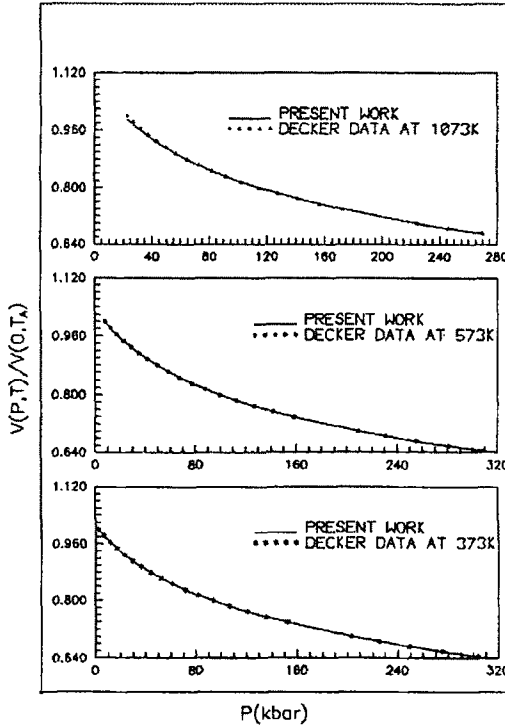


Figure 2. Comparison of $V(P, T)/V(0, T_A)$ with pressure at different temperatures for CsCl.

As far as $V(P, T_A)$ and $B_T(P, T_A)$ are concerned, we have been able to reduce the three parameter EOS [5] into two parameter EOS as given below

$$\frac{V(P, T_A)}{V(0, T_A)} = [(1 + \beta) \exp(ZP) - \beta]^{-1/\eta} \quad (6)$$

$$B_T(P, T_A) = B_T(0, T_A)[1 + \beta(1 - \exp(-ZP))] \quad (7)$$

where

$$\beta = 3B'_T(0, T_A), \quad \eta = [3B_T(0, T_A) + 1]/3 \quad \text{and} \quad Z = [3B_T(0, T_A)]^{-1}.$$

3. Calculations and results

Thus, it is clear from eqs (2)–(8) that $V(P, T)$ can be calculated once the value of $B_T(0, T_A)$, $B'(0, T_A)$, $\gamma(0, T_A)$, μ and $\Theta(0, T_A)$ become known. The best fitted value of μ and other relevant parameters needed for calculation are reported in table 1. By making use of eq. (2) and taking the values of relevant parameters from table 1, the volume is computed for NaCl, CsCl, Mo, and W as a function of pressure and temperature.

The calculated values of $V(P, T)/V(0, T_A)$ for NaCl and CsCl are compared with the available data of Decker [6] in figures 1 and 2. The agreement is very good as the discrepancy lies within $\pm 1.0\%$ in the whole range of pressure and temperature.

The computation of density is done with the help of eq. (2) for Mo, and W. The results are compared with the Hugonit result of Hixon and Fritz [2] in table 2. The discrepancy is $\pm 0.087\%$ for Mo and $\pm 0.09\%$ for W in the whole range of pressure and temperature. Thus, very good agreement is observed.

Hence, it can be concluded that the inclusion of thermal effect into an isothermal EOS is quite successful in representing the volume compression as a function of pressure and temperature.

Acknowledgements

One of the authors (PK) is grateful to the University Grants Commission, New Delhi for the award of senior research fellowship.

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