

An identity for 4-spacetimes embedded into E_5

JOSÉ L LÓPEZ-BONILLA and H N NÚÑEZ-YÉPEZ*

Departamento de Ciencias Básicas, Universidad Autónoma Metropolitana-Azcapotzalco, Apartado Postal 21-726, Coyoacán 04000 D.F., México

*Departamento de Física, Universidad Autónoma Metropolitana-Iztapalapa, Apartado Postal 21-726, Coyoacán 04000 D.F., México

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Abstract. We show that if a 4-spacetime V_4 can be embedded into E_5 then, if b_{ij} is the second fundamental form tensor associated with V_4 , the quantity $(\text{trace } b) \cdot b_{ij}^{-1}$ depends only on intrinsic geometric properties of the spacetime. Such fact is used to obtain a necessary condition for the embedding of a V_4 into E_5 .

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1. Introduction

Let us consider a 4-spacetime local and isometrically embedded into E_5 . This means that there exists the second fundamental form tensor $b_{ij} = b_{ji}$ which fulfils the Gauss and the Codazzi equations [1–3]

$$R_{ijkc} = \varepsilon(b_{ik}b_{jc} - b_{ic}b_{jk}) \quad (1)$$

and

$$b_{ij;k} = b_{ik;j}, \quad (2)$$

where $\varepsilon = \pm 1$, R_{ijkc} is the V_4 Riemann tensor and r stands for covariant derivative. It is well-known that whenever $\det b_j^i \neq 0$ then (1) implies (2); furthermore, it is not difficult to obtain the following relationship [3–7]

$$-24 \det(b_c^r) = K_2 \equiv *R^{*ijk} R_{ijk}, \quad (3)$$

K_2 being one of the Lanczos invariants [8, 9] defined in terms of the double dual [10] of Riemann tensor

$$*R^{*ij}_{rc} \equiv \frac{1}{4} \eta^{ijaq} R_{aq}{}^{me} \eta_{merc}, \quad (4)$$

where η_{abcd} is a Levi-Civita symbol. This work deals with the case in which $K_2 \neq 0$; according to (3), this implies that the inverse matrix to the second fundamental form, b_{ij}^{-1} , exists. We offer a proof that it is possible to find a relation between $*R^{*ijk}$ and b_{ar}^{-1} which is analogous to (1). Among other things this implies that $(\text{trace } b) \cdot b^{-1}$ depends only on the intrinsic geometry of V_4 .

2. Gauss equation for $*R_{ijk}^*$

For any class-one spacetime we can use (1) to obtain the identity [2, 11]

$$pb_{ij} = \frac{K_2}{48} g_{ij} - \frac{1}{2} R_{imnj} G^{mn}, \quad (5)$$

where $p \equiv \varepsilon b^{ar} G_{ar}/3$, $G_{ac} = R_{ac} - Rg_{ac}/2 = *R^{*i}_{aci}$ is the Einstein tensor, $R_{ij} = R^r_{ijr}$ is the Ricci tensor, and $R = R^c_c$ is the scalar curvature. In [2] González *et al* have shown that p is determined by the intrinsic geometry of V_4 . Furthermore, when $p = 0$, using (5) is trivial to establish that \mathbf{b} is determined by the intrinsic geometry of the spacetime and, therefore, that $(\text{trace } \mathbf{b})\mathbf{b}^{-1}$ is also an intrinsic quantity. The problem appears when the case $p \neq 0$ is considered [6, 12, 13] for then the previous statements become non-trivial. It is the purpose of this work to show that $(\text{trace } \mathbf{b})\mathbf{b}^{-1}$ is in fact intrinsic irrespective of the value of p . To this end an equation of the type (1) but relating $*R_{ijk}^*$ to b_{ar}^{-1} is needed.

Yakupov [14, 15] has shown that for every class-one V_4 the following equation holds good

$$*R^{*imnj} R_{acnj} = \frac{K_2}{12} (\delta_a^i \delta_c^m - \delta_c^i \delta_a^m). \quad (6)$$

Substituting (1) into (6) we get

$$\varepsilon \frac{K_2}{24} (\delta_a^i \delta_c^m - \delta_c^i \delta_a^m) = *R^{*ijmnp} b_{an} b_{jc}; \quad (7)$$

assuming $K_2 \neq 0$ then, on multiplying (7) by $b_r^{-1a} b_q^{-1c}$, we easily get

$$\frac{24}{K_2} *R^{*imrq} = \varepsilon (b_{ir}^{-1} b_{mq}^{-1} - b_{iq}^{-1} b_{mr}^{-1}). \quad (8)$$

This expression has the same structure as Gauss equation (1) hence illustrating the analogous role, the curvature tensor and its double dual play. The problem of embedding for a class-one 4-spacetime is thus reduced to analyzing (1) or (8).

González *et al* [2] have studied how (1) implies (5), in an analogous way from (8) we may get

$$b_r^i b_a^{-1j} = -\frac{1}{2} \delta_a^j + \frac{12}{K_2} *R^{*amnj} R^{rn}. \quad (9)$$

This exhibits that $(\text{trace } \mathbf{b})\mathbf{b}^{-1}$ is a quantity which only depends on the intrinsic geometry of V_4 irrespective of the value of p .

Equation (9) is interesting since it allows establishing a new necessary condition for the embedding of a V_4 (but with $K_2 \neq 0$) into E_5 . This condition can be obtained by first contracting i with j in (5), in this way we get

$$p \text{ trace } b = \frac{1}{2} \left(\frac{K_2}{6} - R^{cq} G_{cq} \right); \quad (10)$$

then, on multiplying (5) and (9), we obtain

$$24^* R_{imj}^* R_{ri}^c R^{mn} G^{ri} + K_2 \left[3 \left(\frac{K_2}{8} + R^2/2 - R^{mn} R_{mn} \right) \delta_j^c - R R_j^c + 2 R_n^c R_j^n \right] = 0, \quad (11)$$

where we used (10). We have established that every class-one spacetime with $K_2 \neq 0$ should comply with (11), which is thus a necessary condition for embedding that we have not found previously published.

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