

Group representations and the method of sections

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MS received 29 July 1987

Abstract. The unitary representations of the Euclidean and Poincaré groups are analysed, using viewpoints suggested by the method of sections, as applied to the monopole problem, and by the method of induced representations.

Keywords. Poincaré group; induced representations; coset spaces; fibre bundles.

PACS Nos 02-20; 02-40; 11-30

1. Introduction

The quantum mechanical motion of a charged particle in the field of a static magnetic monopole was first discussed by Dirac (Dirac 1931). The vector potential of such a magnetic field has an unavoidable string singularity at least semi-infinite in extent. Dirac showed that as a kinematic requirement every wave function for the charged particle must necessarily vanish along the string. Consistency of this requirement and of behaviour near the string with quantum mechanical principles led Dirac to the quantization condition

$$eg = \frac{n}{2} \hbar. \quad (1)$$

In an effort to eliminate the string, Wu and Yang (Wu and Yang 1975) developed another way of describing quantum mechanical states in the above situation. The three-dimensional space around and excluding the monopole, $S^2 \times \mathbb{R}_+$, is expressed as the union of two open, overlapping, topologically trivial regions R_a and R_b . For instance, R_a omits the negative z -axis while R_b omits the positive part. In each of the regions R_a and R_b a corresponding singularity free vector potential A_a , A_b can be chosen; in the overlap $R_a \cap R_b$ where both are defined, they are related by a gauge transformation. Each quantum mechanical state is then described by a pair of singularity-free wave functions ψ_a , ψ_b defined on R_a , R_b respectively. Over $R_a \cap R_b$, ψ_a and ψ_b are related by the same gauge transformation that relates A_a to A_b . This transition rule connecting ψ_a and ψ_b in the overlap is kinematic in nature, being the same for all states (in the given monopole field). The requirement that it be well-defined, i.e. single valued, leads us back to the quantization condition (1). This so-called method

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of sections has been extended to treat the unitary irreducible representations of the rotation group $SO(3)$, leading to the monopole harmonics (Wu and Yang 1976; Biedenharn and Louck 1981). The non-Abelian generalization of the magnetic monopole field is discussed by Coleman, and by Goddard and Olive (Coleman 1981; Goddard and Olive 1978).

In another independent approach to the string problem, it has been shown (Balachandran *et al* 1980) that the string can be avoided altogether if one works with a degenerate Lagrangian based on an augmented configuration space $SU(2) \times \mathbb{R}_+$ rather than the physically accessible $S^2 \times \mathbb{R}_+$. Here $SU(2)$ is viewed as a principal fibre bundle, with structure group $U(1)$, on the base S^2 , the sphere of all possible directions emanating from the monopole in physical space. S^2 is itself the space of (right) cosets $SU(2)/U(1)$; and any attempt to reduce the degenerate Lagrangian to a nondegenerate one with $S^2 \times \mathbb{R}_+$ as configuration space automatically reintroduces the string singularity because $SU(2)$ is a nontrivial $U(1)$ principal fibre bundle on S^2 .

These developments bring out the usefulness and clarifying effect of using topological and geometrical notions in such problems and also direct our attention to other problems where, given a Lie group G and a closed Lie subgroup $H \subset G$, the fact that G is a principal fibre bundle on the space of (right) cosets G/H , with H as structure group, plays a basic role.

A well-known method of building up unitary representations (UR) of Lie groups, particularly relevant and useful in quantum mechanics, involves considerations of just the above kind. This is the method of induced representations as developed by Mackey (Mackey 1968; Isham 1983). Given G and H as above, let $D(\cdot)$ be a unitary irreducible representation (UIR) of H in a Hilbert space \mathcal{V} carrying an inner product $(\cdot, \cdot)_{\mathcal{V}}$. One builds up a Hilbert space \mathcal{H} , say, consisting of elements written abstractly as $|\Psi\rangle$, $|\Psi'\rangle$, \dots , and carrying a unitary representation $U(\cdot)$ of G , in this way. Each $|\Psi\rangle$ in \mathcal{H} corresponds to a function on G with values in \mathcal{V} :

$$|\Psi\rangle \rightarrow \psi(g) \in \mathcal{V} \quad \text{for all } g \in G, \quad (2)$$

obeying a covariance condition with respect to H :

$$\psi(gh) = D(h^{-1})\psi(g) \quad \text{for all } h \in H. \quad (3)$$

On such $|\Psi\rangle$, we define a representation $U(\cdot)$ of G by the action

$$U(g_1)|\Psi\rangle = |\Psi'\rangle: \quad \psi'(g) = \psi(g_1^{-1}g). \quad (4)$$

It is immediate that this action preserves the covariance condition (3) and possesses the representation property. Moreover, the covariance condition ensures that $(\psi(g), \psi(g))_{\mathcal{V}}$ is constant over each (right) coset gH in G . At this stage we introduce the symbol Σ to generically denote a (right) coset space, and p for a general point in Σ :

$$\Sigma = G/H, \quad p \in \Sigma. \quad (5)$$

The distinguished point in Σ corresponding to the coset H containing the identity element will be denoted by $p^{(0)}$. One can then make the representation $U(\cdot)$ of G given by (4) a unitary one by defining a suitable invariant inner product

$$\langle \Psi | \Psi \rangle = \int_{\Sigma} d\mu(p) (\psi(p), \psi(p))_{\mathcal{V}}, \quad (6)$$

and thus also completing the definition of the Hilbert space \mathcal{H} . This unitary representation of G is said to be induced from the UIR $D(\cdot)$ of H originally given on \mathcal{V} .

The UIR's of the Euclidean group $E(3)$ (Pauli 1965), i.e. the group of rigid motions in three-dimensional space and of the inhomogeneous Lorentz group, or Poincaré group \mathcal{P} (Wigner 1939; Wightman 1960) of special relativity, are of basic significance in quantum mechanics. Even though these representations have been known and are familiar since a long time, the method of induced representations seems the most elegant and economical way of arriving at them.

It is clear that each $\psi(g)$ obeying the covariance condition (3), though not constant over a coset, is determined all over a coset if its "value" at one representative point in the coset is known. Now G acts on $\Sigma = G/H$ transitively by left multiplication:

$$g \in G, p \in \Sigma \rightarrow p' = gp \in \Sigma. \tag{7}$$

Elements of H are distinguished by the property that they leave $p^{(0)}$ invariant:

$$hp^{(0)} = p^{(0)} \quad \text{for all } h \in H. \tag{8}$$

To choose a coset representative for each coset in G is to pick an element $l(p) \in G$ for each $p \in \Sigma$ such that

$$l(p)p^{(0)} = p \tag{9}$$

in the sense of the action (7) of G on Σ . The "independent information" contained in $\psi(g)$ is then the collection of its values at the elements $l(p) \in G$:

$$\psi(l(p)) \equiv \varphi(p) \in \mathcal{V}. \tag{10}$$

In this sense, each $\psi(g)$ is essentially a function $\varphi(p)$ on Σ with values in \mathcal{V} , and the action of $U(g)$ on $|\Psi\rangle$ can be expressed directly in terms of $\varphi(p)$:

$$\begin{aligned} U(g)|\Psi\rangle &= |\Psi'\rangle \Rightarrow \\ \varphi'(p) &= D(l(p)^{-1}gl(g^{-1}p))\varphi(g^{-1}p). \end{aligned} \tag{11}$$

If an orthonormal basis for \mathcal{V} is indexed by letters α, β, \dots , the components $\varphi_\alpha(p)$ of $\varphi(p) \in \mathcal{V}$ can be written as

$$\varphi_\alpha(p) = \langle p, \alpha | \varphi \rangle, \tag{12}$$

so that $|p, \alpha\rangle$ is an (ideal) orthogonal basis for \mathcal{H} . Then (11) can be reexpressed as

$$\begin{aligned} U(g)|p, \alpha\rangle &= \sum_{\beta} D_{\beta\alpha}(h(g, p))|gp, \beta\rangle, \\ h(g, p) &= l(gp)^{-1}gl(p) \in H. \end{aligned} \tag{13}$$

The inner products among the (ideal) basis vectors $|p, \alpha\rangle$ are fixed by (6). This form of an induced group representation is familiar from the Wigner construction of the UIR's of \mathcal{P} , in which context the element $h(g, p)$ in H in (13) is called the "Wigner rotation".

Given the above general structure, the following questions naturally arise: When can the coset representatives $l(p)$ be chosen in a smooth and singularity-free way all over Σ ?

Since the Wigner rotation is a function of both g and p in general, what is the maximum simplification that can be achieved in its form, at least when g is an element $h \in H$? In this paper we address ourselves to these questions in the context of the representations of $E(3)$ and \mathcal{P} .

The attempt to choose $l(p)$ for each $p \in \Sigma$ obeying (9) is evidently an attempt to find a section when G is viewed as a principal fibre bundle on $\Sigma = G/H$ as base and with H as fibre. It follows that a globally smooth choice of $l(p)$ is possible when and only when this bundle is trivial. If it is nontrivial, any attempt to find an $l(p)$ for all p is bound to encounter singularities, just like the vector potential in the monopole problem. As with the Wu-Yang treatment of quantum mechanical states in the presence of a monopole, we can avoid such "kinematic singularities" in the induced representations of G by expressing the base Σ as the union of (at least) two open topologically trivial regions R_a, R_b :

$$\Sigma = R_a \cup R_b, \quad (14)$$

and choosing $l_a(p), l_b(p)$ over R_a, R_b in a singularity-free way. This means that over each of R_a and R_b , the bundle structure can be trivialised; and the portions of G sitting on top of R_a and R_b, G_a and G_b say, are homeomorphic to the products $R_a \times H, R_b \times H$ respectively. For p in the overlap we necessarily have:

$$\begin{aligned} p \in R_a \cap R_b: l_b(p) &= l_a(p) h_T(p), \\ h_T(p) &\in H. \end{aligned} \quad (15)$$

Then $|\Psi\rangle \in \mathcal{H}$ is specified by a pair of functions $\varphi_a(p), \varphi_b(p)$ defined on R_a, R_b respectively and taking values in \mathcal{V} :

$$\begin{aligned} \varphi_a(p) &\equiv \psi(l_a(p)), \quad p \in R_a; \\ \varphi_b(p) &\equiv \psi(l_b(p)), \quad p \in R_b. \end{aligned} \quad (16)$$

In the overlap we have the kinematic transition rule

$$p \in R_a \cap R_b: \varphi_b(p) = D(h_T(p))^{-1}(\varphi_a(p)). \quad (17)$$

In principle it is straightforward to derive equations which give the effect of $U(g)$ on a pair $(\varphi_a(p), \varphi_b(p))$ to give a new pair $(\varphi'_a(p), \varphi'_b(p))$ also obeying (17). The important point is that in order to avoid string-like singularities in $l(p)$ and the accompanying Wigner rotations, one must pay the price of working with sections rather than globally defined wave functions, at least if the intention is to operate directly with the independent information in $\psi(g)$.

When G is a nontrivial bundle over G/H , one can pose the following further questions: (a) Can the two (or more) open sets R_a, R_b in Σ be chosen so that each is carried into itself under action by elements of H ? (b) If the answer is in the affirmative, can the choices of $l_a(p), l_b(p)$ be made so that at least for elements $h \in H$ the p -dependence in the Wigner rotation is eliminated:

$$\begin{aligned} l_a(hp) &= h l_a(p) h^{-1}, \quad p \in R_a; \\ l_b(hp) &= h l_b(p) h^{-1}, \quad p \in R_b? \end{aligned} \quad (18)$$

The motivation is that if this can be achieved, then for elements $h \in H$ the action of $U(h)$ on $|\varphi\rangle$ would be simple and would use just the UIR $D(\cdot)$ of H originally chosen:

$$\begin{aligned} U(h)|\varphi\rangle &= |\varphi'\rangle: \\ \varphi'_a(p) &= D(h)\varphi_a(h^{-1}p), \\ \varphi'_b(p) &= D(h)\varphi_b(h^{-1}p). \end{aligned} \tag{19}$$

This form, when it can be achieved, is called the Shirokov-Foldy form (see Mukunda 1970).

We shall study the UIR's of $E(3)$ and \mathcal{P} from the point of view of answering the questions raised above. Since our intention is not to derive these well known representations by the method of induction but rather to examine them to answer the questions posed above, we shall be brief in the identifications of G , H and Σ in each case, and shall go on directly to study the relevant topological aspects. Furthermore, we shall be more interested in studying the way the operators $U(h)$ act on $\varphi_a(p)$, $\varphi_b(p)$ for $h \in H$, and not in the manner in which $U(g)$ for general $g \in G$ mix these φ 's.

The paper is organized as follows. In §2 we study the UIR's of $E(3)$, or rather of its two-fold covering group. The results are very similar to what are known for the treatments of the monopole problem. Section 3, devoted to the UIR's of the Poincaré group, begins with notational preliminaries and general considerations common to all the three types of UIR's; we then study in detail the timelike, spacelike, and lightlike UIR's in that sequence, since that is the way the complexities increase. It will be seen that from the topological point of view, the lightlike case is the most intricate. Section 4 is devoted to a summary of our results and some remarks.

2. Representations of the Euclidean group

In this section G will be the (two-fold) covering group of the Euclidean group $E(3)$. It is the semidirect product of $SU(2)$ by the group $T(3)$ of translations in three-dimensional space:

$$G = SU(2) \times T(3). \tag{20}$$

A general element in G is

$$\begin{aligned} g &= (u, \mathbf{a}), \\ u &\in SU(2), \quad \mathbf{a} \in \mathbb{R}^3, \end{aligned} \tag{21}$$

and the composition law is

$$(u', \mathbf{a}')(u, \mathbf{a}) = (u'u, \mathbf{a}' + R(u')\mathbf{a}). \tag{22}$$

The rotation $R(u) \in SO(3)$ determined by $u \in SU(2)$ is defined as usual by

$$u \boldsymbol{\sigma} \cdot \mathbf{x} u^{-1} = \boldsymbol{\sigma} \cdot R(u)\mathbf{x}, \tag{23}$$

where $\boldsymbol{\sigma}$ is the set of Pauli matrices and \mathbf{x} is a real three-vector.

The subgroup $H \subset G$ is the semidirect product of $U(1)$ by $T(3)$, where $U(1)$ is the diagonal sub-group in $SU(2)$:

$$U(1) = \left\{ h(\theta) = \exp\left(\frac{i}{2}\theta\sigma_3\right) \mid 0 \leq \theta < 4\pi \right\} \subset SU(2),$$

$$H = U(1) \times T(3) = \{(h, \mathbf{a}) \mid h \in U(1), \mathbf{a} \in \mathbb{R}^3\}. \tag{24}$$

(Note that we have departed from the notation in the introduction and are using h for an element in the homogeneous part H_0 of H , rather than in H ; this is convenient and adequate in dealing with $E(3)$ as well as with \mathcal{P}). Therefore the coset space $\Sigma = G/H$ is

$$\Sigma = SU(2) \times T(3)/U(1) \times T(3) \simeq SU(2)/U(1) = S^2. \tag{25}$$

Thus the general point $p \in \Sigma$ can be taken to be a real three-vector \mathbf{p} of fixed (say unit) length. The distinguished point $\mathbf{p}^{(0)}$ is the “north pole” of S^2 :

$$\mathbf{p}^{(0)} = (0, 0, 1). \tag{26}$$

This corresponds to setting up the canonical map $\pi: G \rightarrow \Sigma$ by the rule

$$\pi: (u, \mathbf{a}) \in G \rightarrow u\sigma_3 u^{-1} = \boldsymbol{\sigma} \cdot \mathbf{p}, \quad \mathbf{p} \in S^2. \tag{27}$$

The (left) action of $g = (u, \mathbf{a}) \in G$ on Σ amounts to rotating \mathbf{p} by the homogeneous part u of g :

$$\mathbf{p}' = g\mathbf{p} = R(u)\mathbf{p}, \text{ independent of } \mathbf{a}. \tag{28}$$

This can be seen by writing out the subset of elements of G making up the coset g_0H for some $g_0 \in G$, and seeing how this subset changes upon left multiplication by g .

For the inducing procedure one takes the space \mathcal{V} to be one-dimensional, and the UIR $D(\cdot)$ of H to be

$$(h(\theta), \mathbf{a}) \rightarrow \exp(im\theta + i\mathbf{a} \cdot \mathbf{p}^{(0)}),$$

$$m = 0, \pm 1/2, \pm 1, \dots \tag{29}$$

The representation $U(\cdot)$ of G that results is irreducible. It is characterized by helicity m , and magnitude of momentum equal to unity. Here the momentum operators are identified, as usual, as the hermitian generators \mathbf{p} of the translations $T(3) \subset G$.

For all practical purposes, in dealing with the topological aspects and considering choices of $R_{a,b}, l_{a,b}(p)$, etc., one can restrict oneself to the homogeneous parts G_0, H_0 i.e. $SU(2), U(1)$ of G and H . Thus $l_{a,b}(p)$ can be chosen to lie in $SU(2)$. Since $SU(2)$ is a nontrivial $U(1)$ bundle over S^2 , we do need (at least) two open subsets R_a, R_b of $\Sigma = S^2$ to avoid singularities in $l(p)$. Now the action of H_0 on Σ is seen from (28) to consist of rotations about the z -axis. As in the Wu-Yang treatment of the monopole problem, we therefore choose

$$R_a = S^2 - \{(0, 0, -1)\},$$

$$R_b = S^2 - \{(0, 0, 1)\}. \tag{30}$$

This secures invariance of each of R_a, R_b under H_0 action. Therefore it is meaningful to ask for $l_a(\mathbf{p}) \in \text{SU}(2)$ for all $\mathbf{p} \in R_a$ such that

$$l_a(\mathbf{p})\mathbf{p}^{(0)} = \mathbf{p}, \tag{a}$$

$$l_a(h(\theta)\mathbf{p}) = h(\theta)l_a(\mathbf{p})h(\theta)^{-1}, \quad 0 \leq \theta < 4\pi. \tag{b}$$

As is known, such $l_a(\mathbf{p})$ can be found, for example

$$l_a(\mathbf{p}) = [2(1 + p_3)]^{-1/2} (\mathbb{1} + \mathbf{p} \cdot \boldsymbol{\sigma} \sigma_3),$$

$$\mathbf{p} \in R_a. \tag{32}$$

We next ask if an $l_b(\mathbf{p}) \in \text{SU}(2)$ can be found obeying the same (31(a), (b)) but now for $\mathbf{p} \in R_b$. The answer is that this is impossible, and for the following reason. Consider the point $(0, 0, -1) \in R_b$, which has the same (homogeneous) stability group $U(1) \subset \text{SU}(2)$ as $\mathbf{p}^{(0)}$. If a singularity-free choice of $l_b(\mathbf{p}) \in \text{SU}(2)$ over R_b were available obeying (31b), then at $\mathbf{p} = (0, 0, -1)$ we would have

$$l_b((0, 0, -1)) = h(\theta)l_b((0, 0, -1))h(\theta)^{-1}, \quad 0 \leq \theta < 4\pi. \tag{33}$$

This forces $l_b((0, 0, -1))$ to belong to the commutant of $U(1)$ in $\text{SU}(2)$, which is $U(1)$ itself. But then $l_b((0, 0, -1))$ cannot satisfy (31a):

$$l_b((0, 0, -1)) \in U(1) \Rightarrow$$

$$l_b((0, 0, -1))\mathbf{p}^{(0)} = \mathbf{p}^{(0)} \neq (0, 0, -1). \tag{34}$$

Thus the twin conditions (31) can be obeyed on R_a but not on R_b , and it is not possible to bring the action of $U(1)$ to the simple form indicated in (19).

Under these circumstances, a possible natural choice for $l_b(\mathbf{p})$ is

$$l_b(\mathbf{p}) = l_a(-\mathbf{p}) \cdot i\sigma_2, \quad \mathbf{p} \in R_b. \tag{35}$$

This is based on the fact that

$$\mathbf{p} \in R_b \Leftrightarrow -\mathbf{p} \in R_a \tag{36}$$

and that the rotation $R(i\sigma_2)$ carries $\mathbf{p}^{(0)}$ to $-\mathbf{p}^{(0)}$. With this choice for $l_b(\mathbf{p})$ we find that the “transition element” $h_T(\mathbf{p}) \in U(1)$ determined by (15) in $R_a \cap R_b$ is

$$h_T(\mathbf{p}) = l_a(\mathbf{p})^{-1} l_b(\mathbf{p})$$

$$= \exp(-i\phi\sigma_3)$$

$$= h(-2\phi),$$

$$\phi = \text{azimuth of } \mathbf{p} \in R_a \cap R_b. \tag{37}$$

And instead of (31b) obeyed by $l_a(\mathbf{p})$, $l_b(\mathbf{p})$ obeys

$$l_b(h(\theta)\mathbf{p}) = h(\theta)l_b(\mathbf{p})h(\theta), \quad 0 \leq \theta < 4\pi. \tag{38}$$

Therefore if the helicity m UIR of G is set up in the language of sections, the functions

$\varphi_a(\mathbf{p})$, $\varphi_b(\mathbf{p})$ defined over R_a , R_b respectively will be related by the transition rule

$$\begin{aligned} \varphi_b(\mathbf{p}) &= \exp(2im\phi) \varphi_a(\mathbf{p}), \\ \phi &= \text{azimuth of } \mathbf{p} \in R_a \cap R_b. \end{aligned} \tag{39}$$

And the effect of $U(h(\theta))$ on these sections is

$$\begin{aligned} U(h(\theta))|\varphi\rangle &= |\varphi'\rangle \\ \varphi'_a(\mathbf{p}) &= \exp(im\theta) \varphi_a\left(R\left(\exp\left(\frac{-i}{2}\theta\sigma_3\right)\right)\mathbf{p}\right), \\ \varphi'_b(\mathbf{p}) &= \exp(-im\theta) \varphi_b\left(R\left(\exp\left(\frac{-i}{2}\theta\sigma_3\right)\right)\mathbf{p}\right). \end{aligned} \tag{40}$$

The essential point we have brought out is that it is impossible to make $\varphi_a(\mathbf{p})$ and $\varphi_b(\mathbf{p})$ transform identically under the distinguished subgroup $U(1) \subset SU(2)$, provided $m \neq 0$.

We conclude with the remark that in this problem the analogue of the Dirac quantization condition (1) is that, in order that the transition rule (39) be well defined and single-valued over $R_a \cap R_b$, m is restricted to the values $0, \pm 1/2, \pm 1, \dots$

3. Representations of the Poincaré group

3.1 General considerations

The UIR's of the covering group $\bar{\mathcal{P}}$ of the Poincaré group \mathcal{P} have a greater variety than in the case of $E(3)$. As shown by Wigner (Wigner 1939; Wightman 1960), they are of three distinct types: timelike (t), spacelike (s), and lightlike (l), depending on the nature of the eigenvalues of the energy momentum operators. (We ignore those UIR's of \mathcal{P} where the translations are trivially realized). We write p^μ for these eigenvalues, and for convenience we normalize $p^\mu p_\mu$ to -1 in the t and to $+1$ in the s cases. (We use the metric $g_{00} = -1$). In the t and l cases, p^0 may be taken to be positive definite.

The group $\bar{\mathcal{P}}$ is the semidirect product of $SL(2, C)$ with the space-time translation group $T(4)$: (our notations for \mathcal{P} are standard; see, for example, Mukunda 1970):

$$\bar{\mathcal{P}} = SL(2, C) \times T(4). \tag{41}$$

Elements of $\bar{\mathcal{P}}$ and their composition law are:

$$\begin{aligned} g &= (A, a^\mu) \in \bar{\mathcal{P}}, \\ A &\in SL(2, C), \quad a^\mu \in \mathbb{R}^4; \\ (A', a'^\mu)(A, a^\mu) &= (A'A, a'^\mu + \Lambda(A')^\mu_\nu a^\nu). \end{aligned} \tag{42}$$

The homogeneous proper Lorentz transformation $\Lambda(A) \in SO(3, 1)$ determined by A is fixed by

$$\begin{aligned} A \sigma_\mu x^\mu A^\dagger &= \sigma_\mu \Lambda(A)^\mu_\nu x^\nu, \\ \sigma \cdot x \equiv \sigma_\mu x^\mu &= x^0 \cdot \mathbf{1} + \mathbf{x} \cdot \boldsymbol{\sigma}. \end{aligned} \tag{43}$$

Each type of UIR of \mathcal{P} is associated with a particular subgroup of $SL(2, C)$:

$$\begin{aligned}
 t: SU(2) &= \left\{ \begin{pmatrix} \lambda & \mu \\ -\mu^* & \lambda^* \end{pmatrix} \middle| |\lambda|^2 + |\mu|^2 = 1 \right\}; \\
 s: SU(1, 1) &= \left\{ \begin{pmatrix} \lambda & \mu \\ \mu^* & \lambda^* \end{pmatrix} \middle| |\lambda|^2 - |\mu|^2 = 1 \right\}; \\
 l: E(2) &= \left\{ \begin{pmatrix} \lambda & \mu \\ 0 & \lambda^* \end{pmatrix} \middle| |\lambda| = 1 \right\}.
 \end{aligned} \tag{44}$$

For the inducing construction, the necessary subgroups of \mathcal{P} are the semidirect products of these subgroups of $SL(2, C)$ with $T(4)$:

$$H_{t \text{ or } s \text{ or } l} = (SU(2) \text{ or } SU(1, 1) \text{ or } E(2)) \times T(4) \tag{45}$$

The coset spaces $\Sigma_{t,s,l} = \mathcal{P}/H_{t,s,l}$ can be realized in Minkowski space as

$$\begin{aligned}
 \Sigma_t &= \{p^\mu/p^\mu p_\mu = -1, p^0 > 0\}, \\
 \Sigma_s &= \{p^\mu/p^\mu p_\mu = 1\}, \\
 \Sigma_l &= \{p^\mu/p^\mu p_\mu = 0, p^0 > 0\}.
 \end{aligned} \tag{46}$$

\mathcal{P} is a principal fibre bundle over $\Sigma_{t,s,l}$ with structure group $H_{t,s,l}$. The canonical projection map in each case can be given in a manner similar to (27), and once again only the homogeneous part A in the general element (A, a^μ) is involved:

$$\begin{aligned}
 \pi: (A, a^\mu) \in \mathcal{P} &\rightarrow A(\mathbb{1} \text{ or } \sigma_3 \text{ or } \mathbb{1} + \sigma_3)A^+ = \sigma_\mu p^\mu, \\
 &p^\mu \in \Sigma_{t \text{ or } s \text{ or } l}; \\
 p^{(0)\mu} &= (1, 0, 0, 0) \text{ or } (0, 0, 0, 1) \text{ or } (1, 0, 0, 1).
 \end{aligned} \tag{47}$$

The distinguished points $p^{(0)}$ in the three cases are also listed. The action of $g = (A, a^\mu) \in \mathcal{P}$ on a point $p^\mu \in \Sigma_{t,s,l}$ involves only A :

$$g(A, a^\mu): p^\mu \rightarrow p'^\mu = \Lambda(A)^\mu_\nu p^\nu. \tag{48}$$

Therefore the stability group of $p^{(0)}$ in each case is the corresponding H given in (45).

To see when we have a nontrivial bundle structure, it suffices to look at the homogeneous parts. The topological structures of $SL(2, C)$, its relevant subgroups and the Σ 's are:

$$\begin{aligned}
 SL(2, C) &\simeq S^3 \times \mathbb{R}^3; \\
 SU(2) &\simeq S^3, \quad \Sigma_t \simeq \mathbb{R}^3; \\
 SU(1, 1) &\simeq S^1 \times \mathbb{R}^2, \quad \Sigma_s \simeq S^2 \times \mathbb{R}; \\
 E(2) &\simeq S^1 \times \mathbb{R}^2, \quad \Sigma_l \simeq S^2 \times \mathbb{R}.
 \end{aligned} \tag{49}$$

It is immediately clear that \mathcal{P} is a trivial H_t bundle over Σ_t , but a nontrivial $H_{s \text{ or } l}$ bundle over $\Sigma_{s \text{ or } l}$. This is the same as the statement that $SL(2, C)$ is a trivial $SU(2)$ bundle over Σ_t , and a nontrivial $SU(1, 1)$ ($E(2)$) bundle over Σ_s (Σ_l). Therefore in the s and l cases we will have to express Σ as the union of (at least) two open subsets R_a and R_b , as in the $E(3)$ case, in order to have singularity-free expressions for the coset representatives $l(p)$; but Σ_t need not be split up in this way. Even in the t case, however, we can ask if $l(p)$ can be chosen so as to obey (18) with respect to $SU(2)$.

As in the previous section, we use h, h', \dots , to denote elements in the homogeneous parts $SU(2), \dots$, of $H \dots$, rather than in $H \dots$. To obtain the various UIR's of \mathcal{P} via the inducing construction, we pick a UIR $D(\cdot)$ of $SU(2), \dots$ in a Hilbert space \mathcal{V} and extend it to a UIR of $H \dots$ by:

$$(h, a^\mu) \in H \rightarrow D(h) \exp(ia^\mu p_\mu^{(0)}). \tag{50}$$

Here the standard momenta $p^{(0)}$ are as given in (47). For the t case, \mathcal{V} is of finite dimension; for the s case, \mathcal{V} is of infinite dimension for a nontrivial $D(\cdot)$; and for the l case, \mathcal{V} is of dimension one or infinity according as the "translations" in $E(2)$ are realized trivially or nontrivially.

With this general background we can examine now the three types of UIR's of \mathcal{P} from the point of view of the questions posed in § 1. The characteristic variations in the answers as we look at the cases t, s and l in sequence will be brought out.

3.2 The timelike representations

In this case there are no topological obstructions to the choice of a singularity-free coset representative $l(p)$ for all $p \in \Sigma_t$, so the method of sections is unnecessary. If $D(\cdot)$ is a UIR of $SU(2)$ in the space \mathcal{V} (of dimension $(2s + 1)$ corresponding to spin s), the induced UIR of \mathcal{P} operates on functions $\varphi(\mathbf{p})$ with values in \mathcal{V} , i.e. on $(2s + 1)$ -component wave functions $\varphi_\alpha(\mathbf{p})$. The inner product is

$$\begin{aligned} \langle \varphi | \varphi \rangle &= \int_{\Sigma_t} \frac{d^3 p}{p^0} \varphi(\mathbf{p})^\dagger \varphi(\mathbf{p}), \\ p^0 &= (1 + \mathbf{p} \cdot \mathbf{p})^{1/2}. \end{aligned} \tag{51}$$

Correspondingly the (ideal) basis vectors for \mathcal{H} are $|\mathbf{p}, \alpha\rangle$ and obey

$$\langle \mathbf{p}', \beta | \mathbf{p}, \alpha \rangle = p^0 \delta^{(3)}(\mathbf{p}' - \mathbf{p}) \delta_{\beta\alpha}. \tag{52}$$

The action of $U(A, a)$ is given by

$$\begin{aligned} U(A, a) |\mathbf{p}, \alpha\rangle &= \exp(-ia^\mu p'_\mu) \sum_{\beta} D(l(p')^{-1} A l(p))_{\beta\alpha} |\mathbf{p}', \beta\rangle, \\ p'^\mu &= \Lambda(A)^\mu_\nu p^\nu. \end{aligned} \tag{53}$$

The only interesting question is whether $l(p)$ can be chosen so as to satisfy (18):

$$l(hp) = h l(p) h^{-1}, \quad h \in SU(2). \tag{54}$$

The answer is well known: the choice

$$l(p) = [2(1 + p^0)]^{-1/2} (\mathbb{1} + \sigma_\mu p^\mu) \tag{55}$$

does fulfil (54). Therefore if we restrict A in (53) to elements of $SU(2)$, we have the simple action

$$U(h)|\mathbf{p}, \alpha\rangle = \sum_{\beta} D(h)_{\beta\alpha} |\mathbf{p}', \beta\rangle, \\ \mathbf{p}' = R(h)\mathbf{p}, \quad h \in SU(2). \tag{56}$$

3.3 The spacelike representations

Unlike Σ_t which is topologically like \mathbb{R}^3 , Σ_s has a nontrivial structure. As coordinates for Σ_s , we can use the unit vector $\hat{\mathbf{p}} \in S^2$ and $p^0 \in \mathbb{R}$: then the 3-vector \mathbf{p} is

$$\mathbf{p} = (1 + (p^0)^2)^{1/2} \hat{\mathbf{p}}. \tag{57}$$

We know in advance that Σ_s must be expressed as the union $R_a \cup R_b$ of two open, topologically trivial regions (two will do) to avoid singularities in $l(p)$. The first question is whether R_a and R_b can be chosen to be invariant under $SU(1, 1)$. For $h \in SU(1, 1)$, the Lorentz transformation $\Lambda(h)$ acts on p^0, p^1, p^2 , (as an element of $SO(2, 1)$) and leaves p^3 invariant. Therefore we make the choices

$$R_a = \{p^\mu / p^\mu p_\mu = 1, p^3 > -1\}, \\ R_b = \{p^\mu / p^\mu p_\mu = 1, p^3 < 1\}; \tag{58}$$

then each of them is invariant under $SU(1, 1)$. The distinguished point $p^{(0)} = (0, 0, 0, 1)$ lies in R_a . We can now ask for an $l_a(p)$ defined free of singularities all over R_a obeying

$$l_a(p)p^{(0)} = p, \tag{a} \\ l_a(hp) = h l_a(p) h^{-1}, \quad h \in SU(1, 1) \tag{b}$$

It has been shown elsewhere that such $l_a(p)$ can be constructed (Mukunda 1970): an example, strikingly similar to the solution (32) for the E(3) problem, is

$$l_a(p) = [2(1 + p_3)]^{-1/2} (\mathbb{1} + p^\mu \sigma_\mu \sigma_3). \tag{60}$$

Next we show by an argument like the one used in § 2 that it is impossible to find an $l_b(p)$, smoothly defined all over R_b , obeying (59). If such an l_b existed, then consider the point $-p^{(0)} = (0, 0, 0, -1) \in R_b$ which has the same (homogeneous) stability group $SU(1, 1) \subset SL(2, C)$ as $p^{(0)}$. Then at the point $-p^{(0)}$, (59) for l_b would lead to:

$$l_b((0, 0, 0, -1)) = h l_b((0, 0, 0, -1)) h^{-1}, \quad h \in SU(1, 1), \tag{61}$$

i.e. $l_b((0, 0, 0, -1))$ must belong to the commutant of $SU(1, 1)$ in $SL(2, C)$. But this commutant is trivial, consisting just of the two elements $Z_2 = \{\mathbb{1}, -\mathbb{1}\}$; which then means that $l_b((0, 0, 0, -1))$ cannot carry $p^{(0)}$ to $-p^{(0)}$:

$$l_b((0, 0, 0, -1)) \in Z_2 \Rightarrow l_b((0, 0, 0, -1)) p^{(0)} = p^{(0)} \neq -p^{(0)}. \tag{62}$$

The two desirable conditions (59) can therefore be obeyed on R_a but not on R_b . Under these circumstances, a choice for l_b that suggests itself, in analogy to (35) in the $E(3)$ case, is:

$$l_b(p) = l_a(-p) \cdot i\sigma_2, \quad p \in R_b, -p \in R_a. \tag{63}$$

With this $l_b(p)$ the transition element $h_T(p) \in \text{SU}(1, 1)$ can be computed:

$$\begin{aligned} h_T(p) &= l_a(p)^{-1} l_b(p) \\ &= l_a(-p^0, -p^1, -p^2, p^3) l_a(-p) \cdot i\sigma_2 \\ &= \begin{pmatrix} \lambda(p) & \mu(p) \\ \mu(p)^* & \lambda(p)^* \end{pmatrix} \in \text{SU}(1, 1), \\ \lambda(p) &= (-p_1 + ip_2)/(1 - p_3^2)^{1/2}, \quad \mu(p) = -p^0/(1 - p_3^2)^{1/2}, \\ & p \in R_a \cap R_b. \end{aligned} \tag{64}$$

The behaviour of $l_b(p)$ under $\text{SU}(1, 1)$ can also be found; in place of (59), we have:

$$\begin{aligned} l_b(hp) &= l_a(-hp) \cdot i\sigma_2 \\ &= h \cdot l_a(-p) \cdot h^{-1} \cdot i\sigma_2 \\ &= h l_b(p) \tau(h^{-1}), \\ \tau(h) &= \sigma_2 h \sigma_2 \in \text{SU}(1, 1). \end{aligned} \tag{65}$$

Here, $\tau(\cdot)$ is an outer automorphism on $\text{SU}(1, 1)$, the effect on a general element being

$$\tau\left(\begin{pmatrix} \lambda & \mu \\ \mu^* & \lambda^* \end{pmatrix}\right) = \begin{pmatrix} \lambda^* & -\mu^* \\ -\mu & \lambda \end{pmatrix}. \tag{66}$$

Putting together the above results we see that when the spacelike UIR of $\bar{\mathcal{P}}$, based on the UIR $D(\cdot)$ of $\text{SU}(1, 1)$ on \mathcal{V} , is described in a singularity-free manner using the method of sections, each $|\varphi\rangle$ in \mathcal{H} is represented by a pair of functions $\varphi_a(p), \varphi_b(p)$ defined respectively on R_a, R_b and taking values in \mathcal{V} ; they obey the transition rule

$$\varphi_b(p) = D(h_T(p))^{-1} \varphi_a(p), \quad p \in R_a \cap R_b; \tag{67}$$

and with respect to $\text{SU}(1, 1)$ we can at best achieve

$$\begin{aligned} U(h)|\varphi\rangle &= |\varphi'\rangle: \\ \varphi'_a(p) &= D(h) \varphi_a(h^{-1}p), \\ \varphi'_b(p) &= D(\tau(h)) \varphi_b(h^{-1}p), \quad h \in \text{SU}(1, 1). \end{aligned} \tag{68}$$

It is impossible to achieve (19).

The analogues to (51, 52, 53, 56) can be worked out, but we omit the details.

3.4 The lightlike representations

The coset space Σ_l , which is just the positive light cone, can be described in three dimensions as

$$\Sigma_l = \{ \mathbf{p} / p^0 = |\mathbf{p}| > 0 \} = \mathbb{R}^3 - \{ \mathbf{0} \}. \tag{69}$$

It is the omission of the origin from \mathbb{R}^3 that makes the topology of Σ_l nontrivial. Since $SL(2, C)$ is a non-trivial $E(2)$ bundle over Σ_l , we do need (at least) two open topologically trivial regions R_a, R_b with $\Sigma_l = R_a \cup R_b$, to avoid kinematic singularities in $l(p)$. The first question we ask is whether R_a and R_b can be chosen to be both $E(2)$ invariant. Surprisingly, the answer is in the negative.

To see this, we analyse in detail the action of $E(2)$ on Σ_l . A general element in $E(2)$ is

$$h = \begin{pmatrix} \lambda & \mu \\ 0 & \lambda^* \end{pmatrix}, \quad |\lambda| = 1, \quad \mu \in \mathbb{C}. \tag{70}$$

The phase λ produces spatial rotations about the z -axis, which act very simply on Σ_l . The complex parameter μ is just the translation part, $T(2)$, in $E(2)$. If we write $\mu = a_1 - ia_2$ and treat $a_\perp = (a_1, a_2)$ as a transverse 2-vector, we find the following action on a general $p^\mu \in \Sigma_l$:

$$\begin{aligned} h &= \begin{pmatrix} 1 & a_1 - ia_2 \\ 0 & 1 \end{pmatrix} \in E(2), \quad p' = \Lambda(h)p \\ p^{0'} - p^{3'} &= p^0 - p^3, \\ p^{0'} + p^{3'} &= p^0 + p^3 + 2a_\perp \cdot p_\perp + (p^0 - p^3)a_\perp^2, \\ p'_\perp &= p_\perp + (p^0 - p^3)a_\perp. \end{aligned} \tag{71}$$

In the three-dimensional picture of Σ_l given by (69) we see: each point on the positive z -axis is $E(2)$ invariant; each point not on the positive z -axis is either on the negative z -axis or can be transformed to such a point uniquely by a suitable choice of $T(2)$ parameters a_\perp . This motivates the definition of two disjoint subsets of Σ_l , together making up Σ_l , in this way:

$$\begin{aligned} \Sigma_l &= \Sigma_l^{(1)} \cup \Sigma_l^{(2)}, \\ \Sigma_l^{(1)} \cap \Sigma_l^{(2)} &= \phi; \\ \Sigma_l^{(1)} &= \{ (0, 0, \kappa) / \kappa > 0 \}, \\ \Sigma_l^{(2)} &= \{ (p_\perp, p_3) / \text{either } |p_\perp| > 0 \text{ or } p_\perp = 0, p_3 < 0 \}. \end{aligned} \tag{72}$$

In the Minkowski sense, $\Sigma_l^{(1)}$ and $\Sigma_l^{(2)}$ correspond respectively to $p^0 - p^3 = 0$ and $p^0 - p^3 > 0$; so each of them is $E(2)$ invariant. However, while $\Sigma_l^{(2)}$ is an open subset of Σ_l , $\Sigma_l^{(1)}$ is not, and this is the source of the problem. Four-vectors $p^\mu \in \Sigma_l^{(1)}$ are individually $E(2)$ invariant:

$$p^\mu \in \Sigma_l^{(1)} \Leftrightarrow p^\mu = \kappa p^{(0)\mu} = (\kappa, 0, 0, \kappa), \quad \kappa > 0. \tag{73}$$

On $\Sigma_l^{(2)}$ however $E(2)$ acts nontrivially. From (71) we see that the choice

$$a_{\perp} = -p_{\perp}/(p^0 - p^3) \tag{74}$$

gives a $T(2)$ transformation that carries $p^{\mu} \in \Sigma_l^{(2)}$ into

$$p'^{\mu} = \frac{1}{2}(p^0 - p^3, 0, 0, p^3 - p^0). \tag{75}$$

Conversely, starting with $(\kappa, 0, 0, -\kappa)$ for $\kappa > 0$ we can get any desired value for p_{\perp} via a suitable $T(2)$ element:

$$(\kappa, 0, 0, -\kappa) \in \Sigma_l^{(2)} \xrightarrow{T(2)} p^{\mu} = ((p_{\perp}^2 + 4\kappa^2)/4\kappa, p_{\perp}, (p_{\perp}^2 - 4\kappa^2)/4\kappa), \quad \kappa > 0. \tag{76}$$

Returning to the three-dimensional picture of Σ_l : each point on the positive z -axis is $E(2)$ invariant; each point on the negative z -axis expands under $E(2)$ action into a paraboloid of revolution about the z -axis, with vertex on the negative z -axis, and opening out in the direction of increasing z . The open $E(2)$ invariant region $\Sigma_l^{(2)}$ is just the union of all these paraboloids.

With this detailed geometrical picture of the $E(2)$ action on Σ_l viewed as $\mathbb{R}^3 - \{0\}$, the claim made earlier can be proved. If Σ_l is the union of two overlapping open sets R_a and R_b , each $E(2)$ invariant, then $\Sigma_l^{(1)}$ must be contained in (at least) one of them, say in R_b . (The possibility that only part of $\Sigma_l^{(1)}$ is in R_b can be handled in a manner similar to the ensuing argument). Since R_b is open, some open region surrounding $\Sigma_l^{(1)}$ must also be contained in R_b . By the $E(2)$ action, since R_b is assumed to be $E(2)$ invariant, some open collection of the paraboloids of revolution, forming an open region containing $\Sigma_l^{(1)}$, must be in R_b . But that means that R_b is topologically nontrivial, since the origin is excluded. In other words, R_b has the same degree of topological nontriviality as $\mathbb{R}^3 - \{0\}$, if not more. That is, the bundle cannot be trivialized over \mathbb{R}_b . Any local trivializations of the bundle must use topologically trivial open sets R_a and R_b of which at least one is not $E(2)$ invariant.

In view of the above result, let us accept the choice of topologically trivial open regions R_a, R_b according to

$$\begin{aligned} R_a &= \{p^{\mu} / p^{\mu} p_{\mu} = 0, p^0 > 0, p^0 - p^3 > 0\}, \\ R_b &= \{p^{\mu} / p^{\mu} p_{\mu} = 0, p^0 > 0, p^0 + p^3 > 0\}. \end{aligned} \tag{77}$$

We do have $\Sigma_l = R_a \cup R_b$, and

$$p \in R_a \cap R_b \Leftrightarrow p_{\perp} \neq 0. \tag{78}$$

The open set R_a is $\Sigma_l^{(2)}$, and is $E(2)$ invariant. However R_b , which contains $\Sigma_l^{(1)}$ and in particular the distinguished point $p^{(0)} = (1, 0, 0, 1)$, is not $E(2)$ invariant. Coset representatives $l_a(p), l_b(p)$ can be chosen smoothly over R_a, R_b respectively. While for $l_b(p)$ such a question cannot be posed, for $l_a(p)$ we can ask whether it can obey

$$l_a(hp) = h l_a(p) h^{-1}, \quad h \in E(2). \tag{79}$$

It turns out, however, that this is impossible! R_a contains four-momenta of the form

$(\kappa, 0, 0, -\kappa)$ for all $\kappa > 0$. If a smooth choice of $l_a(p)$ over R_a obeying (79) were possible, we could take $p = (\kappa, 0, 0, -\kappa)$ and $h \in U(1)$ in (79) to get:

$$l_a((\kappa, 0, 0, -\kappa)) = \exp\left(\frac{i}{2}\theta\sigma_3\right) l_a((\kappa, 0, 0, -\kappa)) \exp\left(\frac{-i}{2}\theta\sigma_3\right),$$

$$0 \leq \theta < 4\pi. \tag{80}$$

This restricts $l_a((\kappa, 0, 0, -\kappa))$ to be diagonal, and so some spatial rotation around the z-axis followed by a pure Lorentz transformation in the z-direction; it cannot then take $p^{(0)} = (1, 0, 0, 1)$ into $(\kappa, 0, 0, -\kappa)$!

We have thus the result that (79) cannot be obeyed in the $E(2)$ invariant region R_a of Σ_l over which it can be meaningfully posed. One can also see that a change in the specific choice of R_a will not help: any open $E(2)$ invariant region in Σ_l will contain an open collection of the paraboloids of revolution, which then include some momenta of the form $(\kappa, 0, 0, -\kappa)$, and the argument based on (80) will again apply. We must therefore content ourselves with smooth choices of $l_a(p)$ and $l_b(p)$ over R_a and R_b , and a computation of the transition element in $R_a \cap R_b$. No simplification in the action of $E(2)$ on the individual sections along the lines hoped for can be achieved, though the reasons for R_a and for R_b are quite different.

The following are possible choices for coset representatives over R_a and R_b :

$$l_a(p) = [2(p^0 - p^3)]^{-1/2} (\mathbb{1} + \sigma_3 - p \cdot \sigma) \sigma_2, \quad p^0 - p^3 > 0;$$

$$l_b(p) = [2(p^0 + p^3)]^{-1/2} (\mathbb{1} - \sigma_3 + p \cdot \sigma), \quad p^0 + p^3 > 0. \tag{81}$$

We find that this particular choice of $l_a(p)$, while it cannot obey (79), does obey an equation with respect to $E(2)$ involving an outer automorphism τ' on $E(2)$:

$$l_a(hp) = h l_a(p) \tau'(h^{-1}),$$

$$\tau'(h) = \sigma_2 h^{-1 \dagger} \sigma_2 = h^*, \quad h \in E(2). \tag{82}$$

Finally, we calculate the transition element $h_T(p) \in E(2)$:

$$h_T(p) = l_a(p)^{-1} l_b(p)$$

$$= \begin{pmatrix} \lambda(p) & \mu(p) \\ 0 & \lambda(p)^* \end{pmatrix},$$

$$\lambda(p) = (-p_2 + ip_1)/|p_\perp|,$$

$$\mu(p) = 2i(1 - p_3)/|p_\perp|, \quad p \in R_a \cap R_b. \tag{83}$$

Since by (78) p_\perp does not vanish in the overlap, $h_T(p)$ is well-defined.

Putting together the above results we have the following picture. If we start with a UIR $D(\cdot)$ of $E(2)$ in a Hilbert space \mathcal{V} and the standard momentum $p^{(0)} = (1, 0, 0, 1)$, and by the inducing construction build up a lightlike UIR of $\overline{\mathcal{P}}$, then an element $|\varphi\rangle \in \mathcal{H}$ corresponds to the pair of functions $\varphi_a(p)$, $\varphi_b(p)$ defined over R_a , R_b and with values in \mathcal{V} . The transition rule is the same in appearance as (67) in the spacelike case,

but with the definitions of R_a , R_b and $h_\tau(p)$ appropriate for the lightlike situation. Under the action of an element of $E(2)$, we have:

$$U(h)|\varphi\rangle = |\varphi'\rangle:$$

$$\varphi'_a(p) = D(\tau'(h))\varphi_a(h^{-1}p), \quad h \in E(2). \quad (84)$$

For φ'_b we have no simple formula since $E(2)$ does not leave R_b invariant; in fact φ'_b involves both φ_a and φ_b , reminding one of the indecomposability of the lightlike representations of \mathcal{P} found in another context (Matthews *et al* 1974).

4. Concluding remarks

Motivated by the quantum mechanical description of states in the field of a magnetic monopole, we have analysed the UIR's of the (covering groups of the) Euclidean group $E(3)$ and Poincaré group \mathcal{P} , to see how kinematic singularities can be avoided in the wave functions occurring in these UIR's. We have answered the following questions which come up naturally when these UIR's are constructed by the inducing procedure:

(a) In those cases where the group $G(E(3)$ or \mathcal{P}) is a nontrivial principal fibre bundle on the relevant coset space $\Sigma = G/H$, what are natural choices of open topologically trivial regions of Σ over each of which the bundle can be trivialized?

(b) Can each of these regions be chosen to be individually invariant under H_0 , the homogeneous part of H ?

(c) In the UIR of G that results by inducing, starting from a UIR $D(h)$ of H_0 , can the action of elements of H_0 be made as simple as possible and to involve $D(h)$ itself? Can the choices of coset representatives be adjusted to achieve this?

Leaving aside the case of $E(3)$, the UIR's of \mathcal{P} have been found to behave as follows: In the timelike representations, as is well known, the bundle is trivial, and globally smooth choices of coset representatives $l(p)$ do exist. Moreover $l(p)$ can be chosen so that for elements $h \in H_0 = SU(2)$ the Wigner rotation is not momentum dependent. Then the UIR $D(h)$ of $SU(2)$ used in the inducing construction directly describes the $SU(2)$ behaviour of wave functions. For spacelike representations, two regions R_a and R_b in Σ are needed, as in the monopole and $E(3)$ problems, and we can choose them to be individually $SU(1, 1)$ invariant. Over R_a the coset representative $l_a(p)$ can be chosen to have the simplifying property (18); the one over R_b cannot. The lightlike representations show further complexity. Two regions R_a , R_b are needed to trivialize the bundle, but of these only R_a is $E(2)$ invariant. Moreover, even on R_a we cannot pick a smooth coset representative $l_a(p)$ obeying (18). Thus from the view point of this paper we see a gradual increase in the intricacy of the topological structure and group theoretical behaviour as we go from t to s to l .

We have not given general formulae, in the s and l cases, for the effect of a general $U(g)$ on φ_a and φ_b . One expects a mixing of these sections, easy in principle, but tedious in practice to work out. For a comparison, the case of monopole harmonics has been worked out in detail by Wu and Yang.

Finally we would like to collect together and draw attention to the specific coset representatives we have found in the case of the Poincaré group, for the s and l UIR's.

Including also the well-known choice in the timelike case, the complete list is:

$$\begin{aligned}
 t: \quad l(p) &= [2(1+p^0)]^{-1/2} (\mathbb{1} + p \cdot \sigma); \\
 s: \quad l_a(p) &= [2(1+p_3)]^{-1/2} (\mathbb{1} + p \cdot \sigma \sigma_3), \quad p_3 > -1, \\
 & \quad l_b(p) = [2(1-p_3)]^{-1/2} (\mathbb{1} - p \cdot \sigma \sigma_3) i \sigma_2, \quad p_3 < 1; \\
 l: \quad l_a(p) &= [2(p^0 - p^3)]^{-1/2} (\mathbb{1} + \sigma_3 - p \cdot \sigma) \sigma_2, \quad p^0 - p^3 > 0, \\
 & \quad l_b(p) = [2(p^0 + p^3)]^{-1/2} (\mathbb{1} - \sigma_3 + p \cdot \sigma), \quad p^0 + p^3 > 0.
 \end{aligned}$$

It would be interesting to see if there is some general argument leading to such similar forms in these widely differing situations, which then might be meaningful in higher dimensions.

References

- Balachandran A P, Marmo G, Skagerstam B S and Stern A 1980 *Nucl. Phys.* **B162** 385
- Biedenharn L C and Louck J D 1981 *The Racah-Wigner Algebra in Quantum Theory Encyclopaedia of Mathematics and its Applications* (Reading, Mass.: Addison-Wesley) Vol. 9, Chap. 5, Topic 2
- Coleman S 1981 *The magnetic monopole fifty years later*, Lectures delivered at the International School of Subnuclear Physics, Ettore Majorana, Harvard University Preprint HUTP-82/A032
- Dirac P A M 1931 *Proc. R. Soc. London* **A133** 60.
- Goddard P and Olive D I 1978 *Rep. Prog. Phys.* **41** 1357
- Isham C J 1983 *Relativity, Groups and Topology II*, Les Houches, Session XL, Course 11, Section 5.8 (Eds B S Dewitt and R Stora (Amsterdam: North-Holland) p. 1266
- Mackey G W 1968 *Induced representations of groups and quantum mechanics* (N. York and Amsterdam: W. A. Benjamin Inc., Torino: Editore Boringhieri)
- Matthews P M, Seetharaman M and Simon M T 1974 *Phys. Rev.* **D9** 1706
- Mukunda N 1970 *Ann. Phys. (N.Y.)* **61** 329
- Pauli W 1965 *Ergebnisse der Exacten Naturwissenschaften Band 37* p. 85 (Berlin, Heidelberg, New York: Springer-Verlag)
- Wightman A S 1960 *Relations de dispersion et particules elementaires*, Les Houches (Paris: Hermann) p. 159
- Wigner E 1939 *Ann. Math.* **40** 149
- Wu T T and Yang C N 1975 *Phys. Rev.* **D12** 3845
- Wu T T and Yang C N 1976 *Nucl. Phys.* **B107** 365