

## Seiberg–Witten invariants – An expository account

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**Abstract.** The Seiberg–Witten monopole equations, and a new invariant for 4-manifolds which results from these equations, are introduced in this paper.

**Keywords.** Seiberg–Witten invariants; monopoles; 4-manifolds;  $\text{Spin}_c$  structures; Dirac operator.

### 1. Introduction

In this paper, we will introduce the Seiberg–Witten moduli space, which is a moduli space to a system of equations known as the ‘monopole equations’ (taken modulo a gauge group). This leads to invariants for 4-manifolds, which are more manageable than the earlier  $SU(2)$  instanton invariants, and have had some notable new consequences, such as the proof of Thom’s conjecture due to Kronheimer and Mrowka (see the accompanying paper [P]).

The break-up of this paper is as follows. In § 2, the linear algebra related to  $\text{Spin}$  and  $\text{Spin}_c$  groups is introduced. Section 3 discusses  $\text{Spin}_c$  structures on 4-manifolds, and a proof of the theorem of Hirzebruch–Hopf, i.e. that every compact orientable Riemannian 4-manifold admits a  $\text{Spin}_c$  structure is presented. Thereafter, in § 4,  $\text{Spin}_c$  connections, the Dirac operator and monopole equations, and the relevant gauge group is introduced, and on quotienting out by its action, the moduli space for the space of solutions (the Seiberg–Witten moduli space) is constructed, and its virtual dimension calculated from Fredholm considerations. Section 5 contains a proof of the Weitzenböck formula for the  $\text{Spin}_c$ -Dirac operator, and is included for the sake of completeness (even though it is standard). Section 6 contains the proof of the crucial properness result which results in compactness of the moduli space constructed, and the dependence upto cobordism of the moduli space on the Riemannian metric on the 4-manifold is proved (for a more detailed discussion of this matter, see [P]). The last § 7 contains calculations for the case of a Kahler manifold.

For a rapid primer on Fredholm theory, the reader may look at the appendix of the accompanying paper [P].

### 2. Spin groups

We recall some constructions of spin groups in low dimensions.

#### 2.1 Dimension 3

Let  $W$  be a vector space of dimension 2. Consider the representation of  $GL(W)$  on  $\text{End}^0(W)$ , the space of all traceless endomorphisms of  $W$ . There is a natural non-degenerate form  $\{ , \}$

on  $\text{End}^0(W)$  given by

$$\{f, g\} = \text{Trace}_W f \circ g.$$

Moreover, we have a sequence of isomorphisms of representations of  $GL(W)$ ,

$${}^3\wedge \text{End}^0(W) = {}^4\wedge \text{End}(W) = {}^4\wedge(W^* \otimes W) = ({}^2\wedge W^*)^{\otimes 2} \otimes ({}^2\wedge W)^{\otimes 2} = \mathbf{1},$$

where  $\mathbf{1}$  denotes the trivial representation. Thus we obtain a natural homomorphism  $GL(W) \rightarrow SO(\text{End}^0(W))$ . Over the complex numbers this identifies  $GL(2)$  with the ‘Cspin’ group of  $SO(3)$ . The subgroup  $SL(2)$  is identified with the spin group.

### 2.2 Dimension 4

Let  $W_+$  and  $W_-$  be two vector spaces of dimension 2 and let  $\phi : {}^2\wedge W_- \rightarrow {}^2\wedge W_+$  be an isomorphism. Then the vector space  $U = \text{Hom}(W_+, W_-)$  is isomorphic to its dual via a map  $B : U \rightarrow U^* = \text{Hom}(W_-, W_+)$  defined by the identity

$$\phi(f(w_+) \wedge w_-) = w_+ \wedge B(f)(w_-).$$

Thus we have a non-degenerate pairing

$$(f, g) = \text{Trace}_{W_+} B(f) \circ g = \text{Trace}_{W_-} g \circ B(f)$$

which can be seen to be a symmetric form. The group of automorphisms of the triple  $(W_+, W_-, \phi)$  is

$$S(GL(W_+) \times GL(W_-)) = \{(g, h) \mid \det(g) = \det(h)\}.$$

We have a sequence of isomorphisms of representations of this group

$${}^4\wedge(U) = {}^4\wedge(W_+^* \otimes W_-) = ({}^2\wedge W_+^*)^{\otimes 2} \otimes ({}^2\wedge W_-)^{\otimes 2} \xrightarrow{\phi \otimes \phi} \mathbf{1}.$$

Thus we obtain a morphism

$$S(GL(W_+) \times GL(W_-)) \rightarrow SO(\text{Hom}(W_+, W_-)).$$

Over the complex numbers this identifies the group  $S(GL(2) \times GL(2))$  with the ‘Cspin’ group of  $SO(4)$  and the subgroup  $SL(2) \times SL(2)$  is identified with the spin group of  $SO(4)$ .

### 2.3 Dimension 6

Let  $U$  be a four dimensional vector space and let  $\psi : {}^4\wedge U \rightarrow \mathbf{1}$  be a chosen isomorphism so that the group of automorphisms of the pair  $(U, \psi)$  is  $SL(U)$ . Consider the pairing  $\langle, \rangle$  on  ${}^2\wedge U$  given by the composite

$${}^2\wedge U \otimes {}^2\wedge U \xrightarrow{\wedge} {}^4\wedge U \xrightarrow{\psi} \mathbf{1}.$$

This is symmetric and non-degenerate. Moreover, we have a natural sequence of isomorphisms of representations of  $SL(U)$

$${}^6\wedge({}^2\wedge U) = ({}^4\wedge U)^{\otimes 6} \xrightarrow{\psi^{\otimes 6}} \mathbf{1}.$$

Thus we have a representation of  $SL(U)$  in  $SO({}^2\wedge U)$ . Over the complex numbers this identifies  $SL(4)$  with the spin group of  $SO(6)$ .

### 2.4 Combination of the above

Now consider the situation of (2.3) where  $U = \text{Hom}(W_+, W_-)$ . In this situation  $U$  carries a non-degenerate pairing  $(,)$  as described above and hence there is an induced pairing on  $\overset{2}{\wedge}U$  which we also denote by  $(,)$ . We then have an automorphism  $*$  on  $\overset{2}{\wedge}U$  defined by the identity  $(\alpha, \beta) = (\alpha, *\beta)$ . Now the fact that  $\psi = \phi \otimes \phi$  satisfies  $(\psi, \psi) = 1$  implies that  $*^2 = 1$ . Moreover, one can see that the positive (resp. negative) eigenspace  $\Lambda^+$  (resp.  $\Lambda^-$ ) of  $*$  is of dimension 3. Thus the combined representation

$$S(GL(W_+) \times GL(W_-)) \rightarrow SO(U) \hookrightarrow SL(U) \rightarrow SO(\overset{2}{\wedge}U)$$

gives a morphism into  $SO(\Lambda^+) \times SO(\Lambda^-)$ . Now we have natural maps

$$S(GL(W_+) \times GL(W_-)) \rightarrow GL(W_{\pm})$$

and hence we have representations of  $S(GL(W_+) \times GL(W_-))$  into  $SO(\text{End}^0(W_{\pm}))$ . Consider the homomorphisms of representations of  $S(GL(W_+) \times GL(W_-))$

$$\overset{2}{\wedge}U \rightarrow \text{End}^0(W_+) \text{ where } f \wedge g \mapsto B(f) \circ g - B(g) \circ f$$

and similarly

$$\overset{2}{\wedge}U \rightarrow \text{End}^0(W_-) \text{ where } f \wedge g \mapsto f \circ B(g) - g \circ B(f).$$

These induce isomorphisms of  $\text{End}^0(W_{\pm})$  with  $\Lambda^{\pm}$ .

### 2.5 Compact forms

Let us fix hermitian structures  $h_{\pm}$  on  $W_{\pm}$  so that  $\phi$  is an isometry. The group of automorphisms then becomes  $S(U(W_+) \times U(W_-))$ . We define a  $\mathbb{C}$ -anti-linear automorphism  $f \mapsto f^{\dagger}$  defined by the identity

$$h_+(f^{\dagger}(w), w') = h_+(w, B(f)(w')).$$

One sees that  $f^{\dagger\dagger} = f$ . Thus we obtain a real vector space  $T$  so that  $U = T + \iota T$ . Moreover, one sees that the form  $(,)$  restricts to a positive definite form on  $T$ ; hence we obtain a representation  $S(U(W_+) \times U(W_-)) \rightarrow SO(T)$ . The above discussion then gives us a decomposition of  $\overset{2}{\wedge}T$  into  $\Lambda_{\mathbb{R}}^{\pm}$ .

We have a  $\mathbb{C}$ -anti-linear endomorphism  $f \mapsto f^{\dagger}$  of  $\text{End}^0(W_{\pm})$  given by

$$h_{\pm}(f^{\dagger}w, w') = h_{\pm}(w, f(w')).$$

One shows that under the isomorphism between  $\text{End}^0(W_{\pm})$  and  $\Lambda^{\pm} = \Lambda_{\mathbb{R}}^{\pm} + \iota\Lambda_{\mathbb{R}}^{\pm}$ , we obtain identifications of  $\Lambda_{\mathbb{R}}^{\pm}$  with the spaces  $\text{End}^0(W_{\pm})^{ah}$  consisting of  $f = -f^{\dagger}$ .

We note that for any pair of elements  $\Phi, \Psi$  of  $W_+$  we have an element  $\sigma(\Phi, \Psi)$  of  $\text{End}^0(W_+)$  given by

$$w \mapsto i(h_+(w, \Psi) \cdot \Phi - \frac{1}{2} h_+(\Phi, \Psi) \cdot w).$$

When  $\Phi = \Psi$  this is an element of  $\text{End}^0(W_+)^{ah}$ . We identify this with an element of  $\Lambda_{\mathbb{R}}^+$ .

### 2.6 Unitary group case

We now further specialise to the case when  $W_+ = \mathbf{1} \oplus \det W_-$ . For ease of notation we use  $W$  for  $W_-$ . In this case, we have a natural sequence of identifications

$$\text{Hom}_{\mathbb{C}}(W_+, W_-) = W \oplus W^* = W \oplus \overline{W} = W \otimes_{\mathbb{R}} \mathbb{C}.$$

Thus we can identify the special orthogonal representation  $T$  with the underlying real vector space of  $W$ . Now let  $\overset{(2,0)}{\wedge} T$  denote the underlying real vector space to  $\overset{2}{\wedge}_{\mathbb{C}} W$  and let  $\overset{(1,1)}{\wedge} T$  the real vector space such that  $W \otimes \overline{W} = \overset{(1,1)}{\wedge} T \otimes_{\mathbb{R}} \mathbb{C}$ . We have a natural decomposition

$$\overset{2}{\wedge} T = \overset{(2,0)}{\wedge} T \oplus \overset{(1,1)}{\wedge} T.$$

The imaginary part of the hermitian metric on  $W$  gives a natural element  $\omega$  of the latter space. One then computes that

$$\lambda_{\mathbb{R}}^+ = \overset{(2,0)}{\wedge} T \oplus \mathbb{R} \cdot \omega \text{ and } \lambda_{\mathbb{R}}^- = \omega^\perp \cap \overset{(1,1)}{\wedge} T.$$

Moreover, under the identification between  $\Lambda_{\mathbb{R}}^+$  and  $\text{End}(W_+)^{ah}$  we obtain identifications

$$\overset{(2,0)}{\wedge} T = \text{Hom}_{\mathbb{C}}(\mathbf{1}, \det W) = \det W \text{ and } \mathbb{R} = \mathbb{R} \cdot \omega = \mathbb{R}i \cdot \mathbf{1}_{\det W}.$$

### 3. Spin structures on four manifolds

Let  $X$  be a compact oriented four manifold. For any metric  $g$  on  $X$  we have the principal  $SO(4)$  bundle  $P$  on  $X$  which consists of oriented orthonormal frames. This corresponds to a class  $[P]$  in  $H^1(X, SO(4))$ . Using the exact sequence

$$1 \rightarrow U(1) \rightarrow \text{Spin}_c(4) \rightarrow SO(4) \rightarrow 1$$

we see that we have an exact sequence

$$H^1(X, U(1)) \rightarrow H^1(X, \text{Spin}_c(4)) \rightarrow H^1(X, SO(4)) \rightarrow H^2(X, U(1)).$$

We see that the obstruction to giving a reduction of structure group from  $SO(4)$  to  $\text{Spin}_c(4)$  is given by a class in  $H^2(X, U(1))$ . Moreover, from the exact sequence

$$1 \rightarrow \mathbb{Z}/2\mathbb{Z} \rightarrow \text{Spin}(4) \rightarrow SO(4) \rightarrow 1$$

we see that the obstruction to giving a spin structure lies in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$ . Under the natural inclusion of  $\mathbb{Z}/2\mathbb{Z}$  in  $U(1)$ , the obstruction for spin maps to the obstruction for  $\text{Spin}_c$ . In fact consider the diagram

$$\begin{array}{ccccccc} & & 1 & & 1 & & \\ & & \downarrow & & \downarrow & & \\ 1 & \rightarrow & \mathbb{Z}/2\mathbb{Z} & \rightarrow & \text{Spin}(4) & \rightarrow & SO(4) \rightarrow 1 \\ & & \downarrow & & \downarrow & & \parallel \\ 1 & \rightarrow & U(1) & \rightarrow & \text{Spin}_c(4) & \rightarrow & SO(4) \rightarrow 1 \\ & & \downarrow & & \downarrow & & \\ & & U(1) & = & U(1) & & \\ & & \downarrow & & \downarrow & & \\ & & 1 & & 1 & & \end{array}$$

By the associated diagram of cohomologies

$$\begin{array}{ccccccc}
 & & \downarrow & & \downarrow & & \downarrow \\
 \rightarrow & H^1(X, \mathbb{Z}/2\mathbb{Z}) & \rightarrow & H^1(X, \text{Spin}(4)) & \rightarrow & H^1(X, \text{SO}(4)) & \rightarrow & H^2(X, \mathbb{Z}/2\mathbb{Z}) \\
 & \downarrow & & \downarrow & & \parallel & & \downarrow \\
 \rightarrow & H^1(X, U(1)) & \rightarrow & H^1(X, \text{Spin}_c(4)) & \rightarrow & H^1(X, \text{SO}(4)) & \rightarrow & H^2(X, U(1)) \\
 & \downarrow & & \downarrow & & & & \downarrow \\
 & H^1(X, U(1)) & = & H^1(X, U(1)) & & & & H^2(X, U(1)) \\
 & \downarrow & & \downarrow & & & & \\
 & & & & & & & 
 \end{array}$$

we see that the distinct lifts of a given  $\text{SO}(4)$  bundle to a  $\text{Spin}_c(4)$  bundle correspond exactly to the different lifts of the  $\text{Spin}(4)$  obstruction class in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$  to a class in  $H^1(X, U(1))$ . We note that the latter is the group of metrised complex line bundles.

Now we have a natural exact sequence (the exponential sequence) of sheaves

$$0 \rightarrow \mathbb{Z} \rightarrow C^\infty \rightarrow U(1) \rightarrow 1$$

which gives the natural isomorphisms  $H^{i+1}(X, \mathbb{Z}) = H^i(X, U(1))$ . Moreover, under these isomorphisms the exact sequence

$$\rightarrow H^1(X, U(1)) \rightarrow H^1(X, U(1)) \rightarrow H^2(X, \mathbb{Z}/2\mathbb{Z}) \rightarrow H^2(X, U(1)) \rightarrow H^2(X, U(1))$$

is the same as the exact sequence

$$\rightarrow H^2(X, \mathbb{Z}) \rightarrow H^2(X, \mathbb{Z}) \rightarrow H^2(X, \mathbb{Z}/2\mathbb{Z}) \rightarrow H^3(X, \mathbb{Z}) \rightarrow H^3(X, \mathbb{Z}).$$

To summarise, the obstruction to giving a  $\text{Spin}_c(4)$  structure is the image in  $H^3(X, \mathbb{Z})$  of the obstruction to a  $\text{Spin}(4)$  which lies in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$ . If the former is zero then the different  $\text{Spin}_c(4)$  structures correspond to the different lifts of the  $\text{Spin}(4)$  obstruction class to  $H^2(X, \mathbb{Z})$ .

In the case when the principal bundle is the one associated with the metrised tangent bundle as above, we have the result that the obstruction to having a spin structure is given by  $w_2(X)$  in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$ ; the second Stiefel–Whitney class of  $X$ . Then we have Wu’s formula which implies that for any  $y$  in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$  we have  $w_2(X) \cap y = y \cap y$ . Now consider the image  $w$  of  $w_2(X)$  in  $H^3(X, \mathbb{Z})$ ; this is a 2-torsion class. Let  $H^2(X, \mathbb{Z})_\tau$  denote the group of torsion elements in  $H^2(X, \mathbb{Z})$ . There is a natural duality between the 2-torsion in  $H^3(X, \mathbb{Z})$  and the group  $H^2(X, \mathbb{Z})_\tau \otimes \mathbb{Z}/2\mathbb{Z}$ ; this duality is given as follows. Let  $a \in H^2(X, \mathbb{Z})_\tau$  be a torsion class and let  $b \in H^3(X, \mathbb{Z})$  be a 2-torsion class. Let  $b'$  be a class in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$  whose image is  $b$ . Let  $a'$  be the image of  $a$  in  $H^2(X, \mathbb{Z}/2\mathbb{Z})$ , then  $\langle a, b \rangle = \langle a', b' \rangle$ . By this identification we have

$$\langle a, w \rangle = \langle a', w_2(X) \rangle = \text{Trace}(a' \cap w_2(X)) = \text{Trace}(a' \cap a') = 0$$

for all  $a$  in  $H^2(X, \mathbb{Z})_\tau$ . But then by the duality we see that  $w$  is 0. Hence, in this case we obtain that  $w_2(X)$  is the reduction modulo 2 of an integral cohomology class; in other words *an oriented compact Riemannian four manifold always has a  $\text{Spin}_c(4)$  structure.*

#### 4. Monopole equations and their moduli space

In this section we describe the monopole moduli spaces and compute the expected dimension.

4.1 Connections for Spin structures

Let  $(X, g, c)$  be a compact oriented Riemannian four manifold with a  $\text{Spin}_c$  structure (denoted by  $c$ ). Let  $Q$  denote the corresponding principal  $\text{Spin}_c(4)$  bundle over  $X$ . Then the principal bundle of oriented orthonormal frames on  $X$  is given by  $P = Q/U(1)$ . We have a natural torsion free connection on this bundle called the Riemannian connection. The pull-back of this to  $Q$  gives us a 1-form on  $Q$  with values in  $\text{Lie}(SO(4))$  which is invariant for the action of  $\text{Spin}_c(4)$ . Now consider the principal  $U(1)$  bundle  $Q/\text{Spin}(4)$  associated with  $Q$  which is just the space of all unit vectors in the line bundle  $L = \det(W_\pm)$ . Let  $A$  be connection on this line bundle. We can pull this back to a form on  $Q$ . Adding the above two forms together we obtain a connection on  $Q$  which we shall denote by  $\nabla_A$  since the Riemannian connection is unique whereas  $A$  can be varied.

4.2 Dirac equation

Fixing  $A$  for the time being we have the differential operator  $\nabla_A : W_+ \rightarrow W_+ \otimes T^*X$  induced by the connection as above. On the one hand the Riemannian structure gives us a natural (flat) identification between  $T^*X$  and  $TX$  and on the other we have seen that  $TX$  can be thought of as a subspace of  $\text{Hom}_\mathbb{C}(W_+, W_-)$ ; moreover, this identification is also invariant under the connection (flat). Thus by contraction we obtain the composite differential operator of order 1

$$D_A = D_{A,+} : W_+ \rightarrow W_+ \otimes T^*X \rightarrow W_+ \otimes TX \rightarrow W_-.$$

This is called the Dirac operator. As seen earlier we have a natural identification

$$\text{Hom}_\mathbb{C}(W_+, W_-) = \text{Hom}_\mathbb{C}(W_-, W_+).$$

Thus we also obtain an operator  $D_A^* = -D_{A,-} : W_- \rightarrow W_+$ . We have the identity (see § 5)

$$\int_X h_-(-D_{A,+}\Phi, \Psi) = \int_X h_+(\Phi, D_{A,-}\Psi)$$

so that we see that  $D_A^*$  is the adjoint of  $D_A$ . This justifies the notation.

The first monopole equation is the Dirac equation  $D_A(\Phi) = 0$ .

4.3 Second monopole equation

Consider the curvature  $F_A$  of the connection  $A$  on  $L$ . This gives a two form with values in the Lie algebra of  $U(1)$  which is just  $\mathbb{R}$ . Let  $F_A^+$  denote the projection into  $\Lambda^+$ . We have also defined the map  $\sigma : W_+ \otimes \overline{W_+} \rightarrow \text{End}(W_+)^{ah}$ . Moreover, we have obtained an identification of  $\Lambda^+$  with the space of skew-Hermitian endomorphisms of  $W_+$ . The second monopole equation is

$$F_A^+ = \sigma(\Phi, \Phi).$$

4.4 Gauge group

Let  $\mathcal{G} = \text{Map}(X, U(1))$  and consider the action of this group on the space  $\mathcal{N}$  of pairs  $(A, \Phi)$  where  $A$  is a connection on the line bundle  $L$  and  $\Phi$  a section of  $W_+$  given by

$$g \cdot (A, \Phi) = (g \cdot A, g \cdot \Phi) = (A - (1/2\pi i)g^{-1}dg, g\Phi).$$

We see easily that if  $(A, \Phi)$  satisfies the monopole equation then so does  $g \cdot (A, \Phi)$ . In fact we have

$$D_{g \cdot A}(g \cdot \Phi) = g \cdot D_A \Phi \text{ and } F_{g \cdot A} = F_A \text{ and } \sigma(g\Phi, g\Phi) = \sigma(\Phi, \Phi).$$

Thus we may consider the ‘moduli space’ of monopoles

$$M = M_c = \{(A, \Phi) \mid D_A \Phi = 0 \text{ and } F_A^+ = \sigma(\Phi, \Phi)\} / \mathcal{G}.$$

We will show that for ‘good’ metrics this is a compact orientable manifold. We shall also find out how it depends on this choice of metric.

Now let  $\mathcal{W}_\pm$  be the spaces of sections of  $W_\pm$ . We have a map

$$\nu : \mathcal{N} \rightarrow \mathcal{W}_- \text{ given by } (A, \Phi) \mapsto (D_A \Phi).$$

Let  $\mathcal{M}$  denote the inverse image  $\nu^{-1}(0)$  and let  $\mathcal{M}^*$  denote the open subset consisting of pairs  $(A, \Phi)$  where  $\Phi \neq 0$ . The differential of the map  $(A, \Phi) \mapsto D_A \Phi$  is given by

$$(a, \phi) \mapsto D_A \phi + 2\pi i a \circ \Phi.$$

Suppose  $\Psi$  is orthogonal to the image. Then we obtain the equations

$$D_A^* \Psi = 0 \text{ and } \Phi \otimes \Psi = 0$$

by orthogonality with the image of vectors of the form  $(a, 0)$  and  $(0, \phi)$  respectively. Now a solution of an elliptic operator vanishes on an open set only if it is identically 0. Thus we see that  $\Phi = 0$ ; in other words  $\nu$  is a submersion when restricted to the space  $\mathcal{N}^*$  consisting of pairs  $(A, \Phi)$  where  $\Phi \neq 0$ . Thus  $\mathcal{M}^*$  is a manifold (albeit of infinite dimension).

The group  $\mathcal{G}$  acts freely on  $\mathcal{M}^*$  since a solution of an elliptic operator cannot vanish on an open set unless it is 0. Consider the space  $\Omega^{2+}$  consisting of 2-forms invariant under  $*$  with the trivial action of  $\mathcal{G}$ . The map  $\mathcal{M} \rightarrow \Omega^{2+}$  given by  $F_A^+ - \sigma(\Phi, \Phi)$  factors through the quotient  $\mathcal{M}/\mathcal{G}$ . We thus obtain a ‘complex’  $\mathcal{G} \rightarrow \mathcal{M}^* \rightarrow \Omega^{2+}$ . The moduli space can be thought of as being its ‘cohomology’.

#### 4.5 Virtual dimension of the moduli space

To compute the dimension of the moduli space we need to compute the cohomology of the complex of differentials of the complex  $\mathcal{G} \rightarrow \mathcal{M}^* \rightarrow \Omega^{2+}$ . The tangent space to  $\mathcal{G}$  at identity can be identified with  $\Omega^0$  the space of functions and the tangent space to  $\Omega^{2+}$  can be identified with itself since it is a vector space. We have an exact sequence

$$0 \rightarrow T\mathcal{M}^* \rightarrow \Omega^1 \oplus \mathcal{W}_+ \rightarrow \Omega^{2+} \rightarrow 0,$$

where we have identified the tangent space of  $\mathcal{A}$  with  $\Omega^1$  the space of 1-forms. Thus the complex of differentials

$$T\mathcal{G} \rightarrow T\mathcal{M}^* \rightarrow \Omega^{2+}$$

is quasi-isomorphic to the complex

$$\Omega^0 \rightarrow \Omega^1 \oplus \mathcal{W}_+ \rightarrow \Omega^{2+} \oplus \mathcal{W}_-,$$

where the maps are

$$h \mapsto (-dh, 2\pi i h \Phi) \text{ and } (a, \psi) \mapsto (d^+ a - \Im \sigma(\dot{\Phi}, \phi), D_A \phi + 2\pi i a \circ \Phi)$$

for  $h \in \Omega^0$ ,  $a \in \Omega^1$  and  $\phi \in \mathcal{W}_+$ . Here  $d^+$  denotes the exterior derivative combined with the projection to  $\Omega^{2+}$  and  $\Im \sigma(\dot{\Phi}, \phi)$  denotes the skew-Hermitian part of  $\sigma(\Phi, \phi)$ . This complex is homotopic to the complex where the first map is  $h \mapsto (-dh, 0)$  and the second is  $(a, \phi) \mapsto (d^+ a, D_A \phi)$  since the difference between these two complexes is given by compact operators. Thus the index of our complex of differentials is the index of the

complex

$$\Omega^0 \rightarrow \Omega^1 \oplus \mathcal{W}_+ \rightarrow \Omega^{2+} \oplus \mathcal{W}_-,$$

where the maps are

$$h \mapsto (-dh, 0) \text{ and } (a, \psi) \mapsto (d^+a, D_A\psi).$$

This is a topological invariant for the pair  $(X, c)$  by the Atiyah-Singer Index theorem; we call this the virtual dimension of the moduli space. In case we can find a point  $\delta \in \Omega^{2+}$  which is a regular value for the map  $\mathcal{M}^* \rightarrow \Omega^{2+}$  we see that this index will be the dimension of

$$M_{c,\delta} = \{(A, \Phi) \mid D_A\Phi = 0 \text{ and } F_A^+ = \sigma(\Phi, \Phi) + \delta\}/cG.$$

We call this the perturbed moduli space. We will show that such a value of  $\delta$  exists and that  $M_{c,\delta}$  is a compact orientable manifold whose dimension is the virtual dimension.

### 5. Differential calculus

We derive various identities among differential operators in the context of  $\text{Spin}_c$  connections.

#### 5.1 The adjoint of the Dirac operator

We have defined the Dirac operator as the composite

$$D_{A,+} = D_A : W_+ \xrightarrow{\nabla_A} T^*X \otimes W_+ \rightarrow W_-,$$

where the latter map is the contraction under the identification of  $T^*X$  with  $TX \subset \text{Hom}_{\mathbb{C}}(W_+, W_-)$ . We have similarly the Dirac operator  $D_{A,-} : W_- \rightarrow W_+$  since we have an identification of  $\text{Hom}_{\mathbb{C}}(W_+, W_-)$  with its dual space  $\text{Hom}_{\mathbb{C}}(W_-, W_+)$ . In terms of an orthonormal frame of tangent vectors  $\{e_i\}$  we obtain a sequence of identities:

$$(D_A\Phi, \Psi) = \sum_i (e_i \circ \nabla_{e_i}\Phi, \Psi)$$

and since  $(f \circ \Phi, \Psi) = (f^\dagger \circ \Phi, \Psi) = (\Phi, f \circ \Psi)$  for all  $f$  in  $TX$ ,

$$(D_A\Phi, \Psi) = \sum_i (\nabla_{e_i}\Phi, e_i \circ \Psi).$$

Now the fact that  $\nabla$  is a metric connection means that

$$(\nabla_{e_i}\Phi, e_i \circ \Psi) = e_i(\Phi, e_i \circ \Psi) - (\Phi, \nabla_{e_i}(e_i \circ \Psi)).$$

Let  $d\tau$  denote the volume form. Then for any function  $f$  and any vector field  $v$  we have,

$$v(f)d\tau = d(f(v \lrcorner d\tau)) - fd(v \lrcorner d\tau).$$

Thus we obtain

$$\begin{aligned} (D_A\Phi, \Psi)d\tau &= \sum_i d((\Phi, e_i \circ \Psi)e_i \lrcorner d\tau) - \sum_i (\Phi, e_i \circ \Psi)d(e_i \lrcorner d\tau) \\ &\quad - \sum_i (\Phi, \nabla_{e_i}(e_i \circ \Psi)). \end{aligned}$$



For any vector field  $v$  we have the identity

$$d(v \lrcorner d\tau) = \sum_j (e_j, \nabla_{e_j} v) d\tau.$$

Moreover, since  $(e_j, e_j) = \delta_{i,j}$  is a constant we have

$$\begin{aligned} \sum_i (\Phi, e_i \circ \Psi) d(e_i \lrcorner d\tau) &= \sum_{i,j} (\Phi, e_i \circ \Psi) (e_j, \nabla_{e_j} e_i) d\tau \\ &= - \sum_{i,j} (\Phi, e_i \circ \Psi) (\nabla_{e_j} e_j, e_i) d\tau. \end{aligned}$$

The other term can be written as follows

$$\nabla_{e_i} (e_i \circ \Psi) = (\nabla_{e_i} e_i) \circ \Psi + e_i \circ \nabla_{e_i} \Psi.$$

and

$$(\nabla_{e_i} e_i) \circ \Psi = \sum_j (\nabla_{e_i} e_i, e_j) e_j \circ \Psi.$$

Combining the above identities we obtain

$$(D_A \Phi, \Psi) = \sum_i d((\Phi, e_i \circ \Psi) e_i \lrcorner d\tau) - \sum_i (\Phi, e_i \circ (\nabla_{e_i} \Psi)) d\tau.$$

Hence

$$\int_X (-D_{A,+} \Phi, \Psi) = \int_X (\Phi, D_{A,-} \Psi)$$

and  $-D_{A,-}$  is the adjoint operator of  $D_{A,+}$ .

By an entirely similar chain of reasoning we show that the adjoint  $\nabla^*: TX \otimes W_- \rightarrow W_+$  of  $\nabla$  on  $W_+$  is given by

$$\nabla(v \otimes \Phi) = - \left( \sum_i (e_i, \nabla_{e_i} v) \Phi + \nabla_v \Phi \right).$$

In invariant terms, we can describe this as the composite

$$TX \otimes W_+ \xrightarrow{-\nabla} T^*X \otimes TX \otimes W_+ \xrightarrow{\text{Trace} \otimes 1} W_+.$$

### 5.2 The Weitzenbock formula

We now compute the composite  $D_A^* D_A \Phi$ . As before we choose a local orthonormal frame  $\{e_i\}$  for  $X$ . We then have

$$D_A^* D_A \Phi = \sum_i -D_A (e_i \circ \nabla_{e_i} \Phi) = - \sum_{i,j} e_j \circ \nabla_{e_j} (e_i \circ \nabla_{e_i} \Phi).$$

We expand the summand to obtain

$$e_j \circ \nabla_{e_j} e_i \circ \nabla_{e_i} \Phi + e_j \circ e_i \nabla_{e_j} \nabla_{e_i} \Phi.$$

As above the first term can be expanded again as

$$\sum_k (e_k, \nabla_{e_j} e_i) e_j \circ e_k \circ \nabla_{e_i} \Phi = - \sum_k (\nabla_{e_j} e_k, e_i) e_j \circ e_k \circ \nabla_{e_i} \Phi.$$

We obtain the formula

$$D_A^* D_A \Phi = \sum_{i,j,k} (\nabla_{e_j} e_k, e_i) e_j \circ e_k \circ \nabla_{e_i} \Phi - \sum_{i,j} e_j \circ e_i \circ \nabla_{e_j} \nabla_{e_i} \Phi.$$

Now defining  $\nabla_{V,W}^2 = \nabla_V \nabla_W - \nabla_{\nabla_V W}$ ,

$$D_A^* D_A \Phi = - \sum_{i,j} e_j \circ e_i \circ \nabla_{e_j, e_i}^2 \Phi.$$

Similar calculations yield the formula

$$\nabla_A^* \nabla_A \Phi = - \sum_i \nabla_{e_i, e_i}^2 \Phi.$$

Now the difference gives us

$$D_A^* D_A \Phi - \nabla_A^* \nabla_A \Phi = - \sum_{i \neq j} e_j \circ e_i \circ \nabla_{e_j, e_i}^2 \Phi.$$

From the definition of  $\nabla_{\cdot, \cdot}^2$ , we have

$$\nabla_{V,W}^2 - \nabla_{V,W}^2 = \nabla_V \nabla_W - \nabla_{\nabla_V W} - \nabla_W \nabla_V - \nabla_{\nabla_W V} = \nabla_V \nabla_W - \nabla_W \nabla_V - \nabla_{[V,W]}$$

using the fact that the connection is torsion free. Since we have an orthonormal basis we have  $e_i \circ e_j = -e_j \circ e_i$  so that we obtain

$$D_A^* D_A \Phi - \nabla_A^* \nabla_A \Phi = - \sum_{i < j} e_j \circ e_i \circ R(e_j, e_i) \Phi,$$

where  $R(V, W) = \nabla_V \nabla_W - \nabla_W \nabla_V - \nabla_{[V,W]}$  is the curvature tensor.

### 5.3 The curvature tensors

The  $\text{Spin}_c$  connection has been expressed as a sum of the Riemannian connection and the  $U(1)$  connection  $A$  on  $L$ . Thus the curvature tensor  $R$  is also the sum of the Riemann curvature tensor  $S$  and the curvature of  $A$ . The former can be expressed as

$$S(V, W) = \sum_{k,l} (S(V, W) e_l, e_k) e_k \circ e_l.$$

Thus we obtain

$$\sum_{i,j} e_j \circ e_i \circ S(e_j, e_i) = \sum_{i,j,k,l} (S(e_j, e_i) e_l, e_k) e_j \circ e_i \circ e_k \circ e_l.$$

By the orthonormality of  $e_i$ 's we easily resolve the latter to obtain  $\sum_{i,j} (S(e_j, e_i) e_i, e_j)$  which is the negative of the scalar curvature  $s$ . The curvature of  $A$  considered as an operator on  $W_+$  acts as  $2\pi i F_A^+$ . Thus the final (Weitzenbock) formula reads

$$D_A^* D_A - \nabla_A^* \nabla_A = s - 2\pi i F_A^+.$$

### 5.4 Extrema

Let  $x$  be a point of our manifold where  $(\Phi, \Phi)$  attains a maximum. Then for any vector  $v$  at  $x$  we have  $v((\Phi, \Phi))(x) = 0$ . Thus consider the following identity (where  $\mathfrak{R}$  denotes the

real part)

$$\begin{aligned} \Re(\nabla_A^* \nabla_A \Phi, \Phi) &= - \sum_i \frac{1}{2} (e_i e_i(\Phi, \Phi) - (\nabla_{e_i} \Phi, \nabla_{e_i} \Phi)) \\ &\quad + \sum_{i,j} (e_j, \nabla_{e_i} e_i) \Re(\nabla_{e_i} \Phi, \Phi). \end{aligned}$$

Since  $e_i(\Phi, \Phi)(x) = 0$  the last term vanishes at  $x$ . Moreover, since  $x$  is a local maximum for  $(\Phi, \Phi)$  the term  $e_i e_i(\Phi, \Phi)(x)$  is negative. Thus we see that  $\Re(\nabla_A^* \nabla_A \Phi, \Phi)$  is positive at  $x$ .

### 6. The Seiberg–Witten invariants

In this section we construct the Seiberg–Witten invariants. First of all we fix a four manifold  $X$ , a Riemannian metric  $g$  and a  $\text{Spin}_c$  structure  $c$ . At the end of the section we will discuss the independence of the invariants on the metric considered.

#### 6.1 Statement of the basic construction

Let  $M_\delta$  denote the fibre of  $\mathcal{M}/\mathcal{G} \rightarrow \Omega^{2+}$  over the point  $\delta$ . We wish to show that there is a  $\delta$  such that this is a compact manifold. To show this we need to show

1. There are regular values for  $\mathcal{M}^*/\mathcal{G} \rightarrow \Omega^{2+}$ .
2. There are regular values as above such that the fibre of  $\mathcal{M}/\mathcal{G} \rightarrow \Omega^{2+}$  is contained in  $\mathcal{M}^*/\mathcal{G}$ .
3. The map  $\mathcal{M}/\mathcal{G} \rightarrow \Omega^{2+}$  is proper.

#### 6.2 Properness

Let  $\delta_i$  be a convergent sequence of elements in  $\Omega^{2+}$ . This means that the sequence converges in the  $L_k^2$  Sobolev norm for every  $k$ . Let  $(A_i, \Phi_i)$  be such that  $D_{A_i} \Phi_i = 0$  and  $F_{A_i}^+ - \sigma(\Phi_i, \Phi_i) = \delta_i$ . To show properness we need to find a convergent subsequence of  $(A_i, \Phi_i)$ ; for which it is enough to show that this sequence is bounded in the  $L_k^2$  Sobolev norm for every  $k$ .

Let  $B$  be a fixed smooth connection on  $L$ ; we express  $A_i = B + a_i$  where  $a_i$  are 1-forms. Consider the function  $h_i = G * d * a_i$  and let  $g_i^0 = \exp(2\pi i h_i)$ . Then  $g_i^0 \cdot A_i = B + a_i - dh_i$  and we obtain  $* d * (a_i - dh_i) = 0$ . Now we can choose  $g_i$  so that the harmonic part of  $a_i$  lies in the fundamental domain for  $H^1(X, \mathbb{Z})$  in  $H^1(X, \mathbb{R})$ . Thus upto gauge invariance we can replace  $A_i$  by another so that  $* d * a_i = 0$  and the harmonic part  $\alpha_i$  of  $a_i$  is bounded. Let  $b_i = a_i - \alpha_i$ . The second monopole equation becomes

$$d^+ b_i = \sigma(\Phi_i, \Phi_i) - F_B^+ + \delta_i$$

so that an  $L_k^2$  bound on  $\Phi_i$  will give an  $L_k^2$  bound on  $d^+ b_i$ . But now  $b_i = G * d * d^+ b_i$  by the above construction of  $b_i$ ; here  $G$  is the Green’s operator. Thus we obtain a bound on the  $L_{k+1}^2$  norm of  $b_i$  since  $G$  is 2-smoothing.

Let us write  $\Phi_i = \Psi_i + \phi_i$  where  $D_B \Psi_i = 0$  and  $\phi_i$  is orthogonal to the space of solutions of  $D_B$ . Hence  $\phi = GD_B^* D_B \Phi$  and the first monopole equation becomes

$$D_B \Phi_i = -(b_i + \alpha_i) \circ \Phi_i$$

so that an  $L_k^2$  bound on  $\Phi_i$  and  $b_i$  gives us an  $L_2^{k+1}$  bound on  $\phi_i$ . We also need to find a way to uniformly bound  $\Psi_i$ . We do this by finding a uniform bound for  $\Phi_i$ .

Let  $x_i$  be a point where  $(\Phi_i, \Phi_i)$  attains a supremum. Applying the Weitzenbock formula we see that at  $x_i$  we have

$$0 \geq \Re(\nabla_{A_i}^* \nabla_{A_i} \Phi_i, \Phi_i)(x_i) = -\Re(s\Phi_i, \Phi_i)(x_i) + 2\pi\Im(F_{A_i}^+ \Phi_i, \Phi_i)(x_i).$$

Note that  $F_{A_i}^+$  is a skew-Hermitian endomorphism of  $W_+$  and thus

$$\Im(F_{A_i}^+ \Phi_i, \Phi_i) = (F_{A_i}^+ \Phi_i, \Phi_i).$$

The second monopole equation gives us

$$F_{A_i}^+ \Phi_i = \sigma(\Phi_i, \Phi_i)\Phi_i + \delta_i \Phi_i$$

and the expression for  $\sigma$  gives us

$$\sigma(\Phi_i, \Phi_i)\Phi_i = \frac{i}{2}(\Phi_i, \Phi_i)\Phi_i.$$

Combining the above we obtain

$$(\Phi_i, \Phi_i)(x_i) \leq \max\{0, -s + \|\delta_i\|\}.$$

Thus we uniformly bound  $\Phi_i$  in the  $C^0$ -norm. This gives us uniform bounds for  $\Psi_i$  and  $\phi_i$  in the  $C^0$ -norm. Now  $\Psi_i$  are solutions of the Dirac equation  $D_B \Psi_i = 0$ . Thus the set of  $C^0$ -bounded solutions is a compact set; in particular, we obtain  $L_k^2$  bounds on  $\Psi_i$  for all  $i$ .

The above arguments applied inductively gives the required result. We note that the above arguments also prove that the solutions of the monopole equations are smooth since any solution which is bounded in  $L_k^2$  norm for some  $k$  is actually bounded in all  $L_k^2$  norms as above.

### 6.3 Regular values

Now consider the compact space  $M_0$  of solutions of the unperturbed monopole equations. For each point  $(A, \Phi)$  of  $M_0$  we have a neighbourhood  $U$  in  $\mathcal{N}$  of  $(A, \Phi)$  and a finite dimensional linear space  $H \subset \Omega^{2+}$  such that the composite

$$U \rightarrow \mathcal{N} \rightarrow \mathcal{W}_- \times \Omega^{2+} \rightarrow H^\perp$$

is a submersion. By compactness we can find a common  $H$  and a saturated (for  $\mathcal{G}$ ) open set  $U$  in  $\mathcal{N}$  containing the inverse image of  $M_0$  such that the above composite is a submersion. Since the derivative is a Fredholm map, the fibre over 0 is a finite dimensional manifold  $N$ . We now consider the map of finite dimensional manifolds  $N \rightarrow H$ . By Sard's theorem we have a dense subset of  $H$  which consists of regular values.

Now assume that  $b_2^+$  which is the codimension in  $\Omega^{2+}$  of the  $\delta$ 's of the form  $F_B^+ + d^+b$  is greater than zero. Then the collection of those  $\delta$  for which the fibre is contained in  $\mathcal{M}^*$  is a non-empty open set. If  $b_2^+ > 1$  then this open set is even path-wise connected. Thus in this situation the cobordism class of the fibre is independent of the regular value chosen.

### 6.4 Dependence on the metric

Let  $\mathcal{C}$  denote the space of all metrics  $g$  on  $X$  under which the fixed volume form  $d\tau$  has norm one. We have a natural map  $\mathcal{C} \rightarrow G$  where  $G$  denotes the Grassmannian of rank  $b_2^+$  quotients of  $H^2(X, \mathbb{R})$ . The corresponding tangent level map is

$$\text{Hom}(\Lambda^-, \Lambda^+) = TC \rightarrow TG = \text{Hom}(H_g^{2-}, H_g^{2+}),$$

where a map  $f : \Lambda^- \rightarrow \Lambda^+$  goes to its harmonic projection.

For any class  $c = c_1(L)$  in  $H^2(X, \mathbb{R})$  let  $S_c$  denote the subvariety of  $G$  where the class  $c$  goes to zero in  $H^{2+}$ . At a point of  $S_c$  the tangent space to  $S_c$  is given by the kernel of the evaluation map  $g \mapsto g(c)$ . Consider the composite map

$$\text{Hom}(\Lambda^-, \Lambda^+) \rightarrow \text{Hom}(H_g^{2-}, H_g^{2+}) \rightarrow H_g^{2+}.$$

If we show that this map is surjective, then the space of all metrics under which the class  $c$  becomes  $*$ -anti-invariant will be of codimension  $b_2^+$ . The argument of the previous section will apply to show that the Seiberg–Witten invariant is independent of the metric when  $b_2^+ > 1$ .

To show that the above map is surjective suppose that  $d$  is perpendicular to the image. We will then obtain that  $c \otimes d$  is identically zero. But now if  $c \neq 0$  then it is represented by a harmonic form which cannot vanish on an open set. Thus  $d$  must vanish on an open set. But we represent  $d$  by a harmonic form too. Thus  $d = 0$  as required.

### 7. The case of Kähler manifolds

We now specialize to the case of Kähler surfaces.

#### 7.1 Spin structures

For any four manifold with almost complex structure and (hermitian) metric we have a natural  $\text{Spin}_c$  structure given by taking  $W_+^0 = \wedge^2_{\mathbb{C}} TX \oplus \mathbf{1}$  and  $W_-^0 = TX$ . The inclusion of  $TX$  in  $\text{Hom}_{\mathbb{C}}(W_+^0, W_-^0)$  is the natural one as discussed at the end of §2. Thus any  $\text{Spin}_c$  structure on  $X$  is given by  $W_+ = M \otimes_{\mathbb{C}} \wedge^2_{\mathbb{C}} TX \oplus M$  and  $W_- = TX \otimes_{\mathbb{C}} M$ . For ease of notation we adopt the standard convention  $\wedge T^*X = K_X$ .

#### 7.2 Spin<sub>c</sub> connections

Any  $U(2)$  connection on  $TX$  gives a connection on all associated bundles. In particular we obtain connections on  $W_{\pm}^0$ . However, in order that these be  $\text{Spin}_c$  connections it is necessary that the induced connection on  $TX$  be the Riemannian (torsion-free) connection. This can only happen if the (almost) complex structure is parallel with respect to the Riemannian connection; thus in this case the manifold must be Kähler.

To give a connection in the general  $\text{Spin}_c$  structure we need in addition to give a  $U(1)$  connection on  $M$ .

#### 7.3 The first monopole equation

Consider a  $\text{Spin}_c$  connection as above. We then obtain a Dirac operator on  $M \oplus M \otimes K_X^{-1}$ . By the above discussion we note that the restriction of this to  $M$  is the composite

$$M \rightarrow M \otimes_{\mathbb{R}} T^*X = M \otimes_{\mathbb{C}} T^*X \oplus M \otimes_{\mathbb{C}} \overline{T^*X} \rightarrow M \otimes_{\mathbb{C}} TX.$$

Here we have used the identification of  $T^*X$  with  $\overline{TX}$  given by the hermitian structure. The first map in the above composite is the  $U(1)$  connection on  $M$ . Thus we see that the restriction of the Dirac operator to  $M$  is  $\nabla^{(0,1)}$ . We similarly show that the restriction of the Dirac operator to  $M \otimes K_X^{-1}$  is also  $\nabla^{(0,1)}$  for the induced  $U(1)$  connection on this line bundle.

#### 7.4 The second monopole equation

Following §2 we compute that the  $(2, 0)$  part of  $\sigma(\Phi, \Phi)$  for  $\Phi = (\alpha, \beta)$  is  $\bar{\alpha}\beta$  and the  $(1, 1)$  part of it is  $\frac{1}{2}(\|\beta\|^2 - \|\alpha\|^2)$ . Thus the second monopole equation becomes

$$(F_A^+)^{(2,0)} = \bar{\alpha}\beta \text{ and } (F_A^+)^{(1,1)} = \frac{1}{2}(\|\beta\|^2 - \|\alpha\|^2)\omega.$$

### 7.5 The Weitzenbock formula

We next apply the Weitzenbock formula for any pair  $(A, \Phi)$

$$D_A^* D_A \Phi = \nabla_A^* \nabla_A \phi + s\Phi - 2\pi i F_A^+ \Phi$$

to obtain an equality of global inner products

$$(D_A \Phi, D_A \Phi)_X = (\nabla_A \Phi, \nabla_A \Phi)_X + (s\Phi, \Phi)_X + 2\pi \Im(F_A^+ \Phi, \Phi)_X.$$

On the other hand we compute the global norm of  $F_A^+ - \sigma(\Phi, \Phi)$  as follows

$$\|F_A^+ - \sigma(\Phi, \Phi)\|_X^2 = \|F_A^+\|_X^2 + \|\sigma(\Phi, \Phi)\|_X^2 - 2\Re(F_A^+, \sigma(\Phi, \Phi))_X.$$

The last term is computed by the integral of the function

$$\Re \text{Trace}_{w_\tau}(F_A^+ \circ \sigma(\Phi, \Phi)) = -\Im(F_A^+ \Phi, \Phi) + \frac{1}{2} \text{Trace}(F_A^+) \|\Phi\|^2.$$

Now  $\text{Trace}(F_A^+)$  is 0. Thus adding the above two identities we obtain

$$\|D_A \Phi\|_X^2 + \|F_A^+ - \sigma(\Phi, \Phi)\|_X^2 = \|\nabla_A \Phi\|_X^2 + (s\Phi, \Phi)_X + 2\pi(\|F_A^+\|_X^2 + \|\sigma(\Phi, \Phi)\|_X^2).$$

We note that the right hand side is equal to

$$\|\nabla_A \alpha\|_X^2 + \|\nabla_A \beta\|_X^2 + (s\alpha, \alpha)_X + (s\beta, \beta)_X + 2\pi\|F_A^+\|_X^2 + 2\pi(\|\alpha\|_X^2 + \|\beta\|_X^2)^2$$

which is invariant under a change of sign for  $\alpha$  or  $\beta$ .

Now suppose that  $(A, \Phi)$  solve the monopole equations and consider the pair  $(A, \Phi_1)$  where  $\Phi_1 = (\alpha, -\beta)$ . By the above discussion we see that  $(A, \Phi_1)$  is also a solution for the monopole equations. But then we must have

$$(F_A)^{(2,0)} = \bar{\alpha}\beta = -\bar{\alpha}\beta = 0.$$

Thus we obtain the fact that  $F_A$  is a holomorphic connection on  $M^{\otimes 2} \otimes K_X$ . Moreover, by ellipticity of the Dirac operator (and its components) we must have that either  $\alpha$  or  $\beta$  is zero according as  $(F_A^+)^{(1,1)}$  is a positive or negative multiple of  $\omega$ . By the first monopole equation it then follows that  $\alpha$  and  $\beta$  are holomorphic sections of the corresponding line bundles.

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