Tidal and gravity waves study from the airglow measurements at Kolhapur (India)

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Simultaneous photometric measurements of the OI 557.7 nm and OH (7, 2) band from a low latitude station, Kolhapur (16.8°N, 74.2°E) during the period 2004–2007 are analyzed to study the dominant waves present in the 80–100 km altitude region of the atmosphere. The nocturnal intensity variations of different airglow emissions are observed using scanning temperature controlled filter photometers. Waves having period lying between 2 and 12 hours have been recorded. Some of these waves having subharmonic tidal oscillation periods 4, 6, 8 and 12 hours propagate upward with velocity lying in the range 1.6–11.3 m/s and the vertical wave length lying between 28.6 and 163 kms. The other waves may be the upward propagating gravity waves or waves resulting from the interaction of inter-mode tidal oscillations, interaction of tidal waves with planetary waves and gravity waves. Some times, the second harmonic wave has higher vertical velocity than the corresponding fundamental wave. Application of these waves in studying the thermal structure of the region is discussed.

1. Introduction

Atmospheric tides, planetary waves and gravity waves play significant role in dynamic features of the Mesosphere and Lower Thermosphere (MLT) region by transporting energy and momentum horizontally and vertically upward and also providing dynamical linkage between the lower atmosphere and the MLT region. Atmospheric tides are the global response of the atmosphere to the periodic forcing of solar heating; propagate westward following the motion of the Sun (Chapman and Lindzen 1970) with periodicities equal to the

solar day and its sub-harmonics (i.e., 24, 12, 8, 6, 4, 3, 2 hr, etc). Tides are classified as migrating and non-migrating. The non-migrating tides lead to strong longitudinal variations in amplitude and phase of diurnal tide, whereas the migrating tides contain diurnal, semidiurnal, terdiurnal and other harmonics. The diurnal (24 hr period) tide is mainly excited by the direct absorption of the sunlight by water vapour in the troposphere and the stratosphere, whereas the semidiurnal tide (12 hr periods) is caused by ozone heating in the upper stratosphere and the lower mesosphere. Studies of the diurnal and the semidiurnal tides and its

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influences in the mesospheric region at various geographical locations have been made using lidar (Taylor *et al.* 1999; Liu *et al.* 2007), radar (Zhang *et al.* 2003; Forbes *et al.* 2004; Manson *et al.* 2004), airglow emissions (Wiens and Weill 1973; Takahashi *et al.* 1998; Lopez-Gonzalez *et al.* 2005; Won *et al.* 2007) and space-based airglow observations (Abreu and Yee 1989; Shepherd *et al.* 1995, 1998, 2004).

Airglow emissions are an outcome of chemiluminescence property of ambient atomic molecular species in the middle and the upper atmosphere. The variability of emission intensity is dominated by a long period as well as a short period structure that are commonly attributed to the passage of tides, gravity waves, planetary waves and their interplay. A large variation in diurnal behaviour of the airglow emission rates as a function of the year and of the latitude has been observed (Manson *et al.* 2004). Lopez-Gonzalez *et al.* (2005) analyzed airglow observations (the O₂ atmospheric band and the OH Meinel band) during the period 1998–2003 and showed that the amplitudes of the diurnal variation from summer to late autumn decrease by more than a factor of two compared with those in the winter and spring. Both temperatures and emission rates showed the dominance of waves with predominantly semidiurnal modulation throughout the year. However, during spring, the semidiurnal variation changes to a diurnal type. They also observed upward energy propagation during most of the year, with some indication of downward energy propagation close to the equinoxes. Taori and Taylor (2010) analyzing mesospheric O₂ airglow emission intensity and temperature data showed the dominant presence of terdiurnal (8 hr periods) tide-like wave and quasi 5 days wave leading to a possible signature and interplay of tides and planetary waves. Ground observations (Vincent *et al.* 1998) as well as satellite observations (Burrage *et al.* 1995; McLandress *et al.* 1996) showed the existence of a strong semiannual tidal amplitude variation in the mesosphere.

The wavy structures observed in the mesosphere night glow intensity variations have also been attributed to the upward propagating gravity waves generated in the lower atmospheric regions (Krassovsky 1972; Hines and Tarasick 1987; Tarasick and Hines 1990; Tarasick and Shepherd 1992a, 1992b; Fagundes *et al.* 1995; Nakamura *et al.* 1999) principally by deep convective activity of thunderstorm (Brown *et al.* 2004; Hecht *et al.* 2004; Snively and Pasko 2008; Vadas *et al.* 2009; Siingh *et al.* 2011, 2012) and orography effects (Tsuda *et al.* 1989; Eckermann and Preusse 1999). Recently Das *et al.* (2011) have proposed heating effect during dust storm also to be one of the probable generation mechanisms of gravity wave in the middle

atmosphere. Multispectral night glow emissions recorded at low latitude stations showed the presence of gravity waves with periods of few hours (Molina *et al.* 1985; Takahashi *et al.* 1985). Short period (tens of minutes) and short horizontal wavelength (tens of kms) gravity waves in the night glow emissions have also been observed (Nakamura *et al.* 1999; Hecht *et al.* 2001; Hecht *et al.* 2004; Pautet *et al.* 2005). Hecht *et al.* (2009) showed the presence of long period (~ 2 hr) and large horizontal wavelength (~ 300 – 400 km) gravity waves as well as short period (15–25 min) and small horizontal wavelength (~ 30 – 45 km) gravity waves in the airglow data over Alice Spring, Australia. The short period waves might have been trapped/ducted by thermal ducts and took several hours to reach the mesopause region. The temperature profile recorded by SABER (Sounding of Atmosphere using Broadband Emission Radiometry) showed the presence of thermal duct (Fadnavis *et al.* 2009). Winds critically control the formation of the duct and hence variations in the wind direction and amplitude may account for sporadic/short-lived observations of the short and the long period waves.

Recently, Narayanan *et al.* (2010) reported the presence of large- and small-scale wave features perpendicularly aligned to each other in the airglow observations, which were associated with gravity wave dynamics. The large scale wave with a horizontal wavelength of ~ 43 km propagated towards the north-west at an apparent phase speed of ~ 48 m/s. These waves near 88 km altitude experienced reflection when the background wind was blowing in an opposite direction to that of wave propagation. The small scale waves with periodicities between 23 and 50 min were continuously appearing and disappearing in different parts of the sky within the field of view of the imaging camera. The small scale waves could have been generated locally by the convective instability (Li *et al.* 2005a, 2005b). The smaller period (less than 12 hr) waves showed a strong wave dissipation throughout the year except during equinox transitions when a large (a factor of about 3–4) amplitude growth was deduced from the analysis of O₂(0–1) and OH(6–2) atmospheric band emissions (Taori *et al.* 2007).

In the middle atmosphere, the thermal structure has been attributed to turbulent heating associated with gravity waves (Fadnavis *et al.* 2009) and tidal waves (Dao *et al.* 1995). The thermal structures so formed depend on wave parameters (Fritts and Alexander 2003). As the wave parameters change from day-to-day and with latitude, we make an attempt to study the characteristic feature of waves present in the airglow emission intensity observed at Kolhapur, a low latitude

station. The data for the years 2004, 2005 and 2007 are used to decipher any possible variation in wave parameters with year. In the present paper, simultaneously measured nightglow emissions data of OI 557.7 nm and OH band are analyzed to study wavy structures in the mesosphere and the lower thermosphere region. The analysis showed presence of terdiurnal and sub-harmonics of tidal oscillations. Waves of shorter periods are also observed. Using the observation of emissions from the two different layers, we determine the horizontal wavelength and velocity of vertical wave propagation. Section 2 deals with observations and analysis of the data. Results and discussions are given in sections 3. Section 4 discusses tidal influences in airglow emissions and brief summary and conclusions are given in section 5.

2. Observations and analysis techniques

Regular observations of the night airglow emissions, OI 557.7 nm and OH Meinel (683 and 692 nm) bands have been carried out at the low latitude station Kolhapur (16.8°N, 74.2°E) on clear and moonless nights since 1989 using filter scanning photometer having five optical filters with integration time of 30 s each. The half field of view of the photometer is 7.25°. Three Stepper Motors control the movement of five filters in the meridional and the azimuthal directions. Before starting observations, the azimuthal position and the position of particular filters were brought to a home position. The scanning photometer was kept in the north-south direction. Three optical sensors were used for this purpose to get the home position of the filters as well as the meridional and the azimuthal positions. The filters have a band width of 1 nm and their temperature is controlled by a temperature controller at 25°C. The temperature coefficient of filter is 0.011 nm/°C. At 25°C, the transmission efficiency of filters is ~60–70%. The filter scanning photometer consists of a photomultiplier tube (EMI 9658B) as a detector which has a wide spectral response. Data from each filter are stored in the digital form in the hard disk of a PC through RS-232 port after amplification of the signal from the photomultiplier tube (EMI 9658B). The photometer was upgraded with a microprocessor for control and data acquisition and it takes about 2.5 minutes to complete one sequence of observations. The integration time of optical filters is ~30 s.

In the absence of a standard calibration source, we have used relative intensities in arbitrary unit. In order to study the wave features present in the MLT region, we consider only those clear sky days having more than 4 hours of continuous and

simultaneous observations of the OI 557.7 nm, and OH Meinel band emissions. The analysis is carried out for the period 2004–2007. For further analysis, airglow intensity values are averaged at an interval of 15 minutes. In this approach one can decipher wavy structures with periods more than 30 minutes (frequency $\sim 1/2dT$). To retrieve long period waves short period fluctuations are smoothed out. Long period waves mostly represent tidal structures and gravity waves originating from the troposphere and stratosphere.

The presence of the short period and the long period oscillations in the nocturnal airglow emissions have been studied using various mathematical filtering techniques (Daubechies 1988; Mukherjee 2003; Taori and Taylor 2010). In the present work, we developed and modified mat lab program for cosine wave fitting and used it for all dataset. From the raw data mean intensity value is determined, which is subtracted from the individual intensity data and mean intensity deviations are obtained. These intensity deviations are used for the cosine wave fitting and further analysis. The periodicity of waves and residuals are determined. Analyzing residuals, the dominant harmonic wave periodicity and amplitudes are estimated. The phase lags/leads are used to evaluate wave phase velocity.

3. Results and discussion

The nocturnal variations observed on 54 clear nights in OI 557.7 nm and OH (7,2) bands are analyzed using the technique discussed above. The total data length on individual nights is ~9–10 hours. Based on the previous studies, it is assumed that the emission 557.7 nm comes from the 97 km altitude and the OH band is emitted at about 86 km altitude. To measure the wave parameters obtained in the data, a simple best fit wave perturbation cosine model is used (Taori *et al.* 2003, 2005) and periodicity, phase and amplitude of the primary wave and its sub-harmonics are estimated from the intensity of OI and OH emissions independently. The modelled best fit solutions were used to get the amplitude and phase information of the waves present in the intensity data. From all the plots it is evident that the data is dominated by long period wave with short period oscillations superposed on it.

Figure 1 shows the mean intensity deviations of the emission lines (683, 692 nm) of OH band on the nights of March 12–13, 2007 and April 11–12, 2007 respectively. The wave signature (~6 hr period) seen in the mean intensity deviations could be due to the solar atmospheric tides. It is noted that the minimum mean intensity deviation value is observed around mid-night at about 23–24 hr

Figure 1. The mean intensity deviations of the OH emission lines (683, 692 nm) with time in Indian Standard Time (IST) observed at Kolhapur for the night of March 12–13, 2007 (upper panel) and for the night of April 11–12, 2007 (lower panel).

and, the maximum mean intensity deviation occurs at the evening and early morning hours. The error bars show the range of airglow intensity under each deviation peak. Same nature is also observed on April 11–12, 2007. When we average the intensity data for four days and plot its variation with time, an interesting result emerges. Figure 2(a, b) shows four days (12–15 March 2007) average intensity variation with time for OI 557.7 nm and OH (7, 2) lines respectively, where (a) plots the mean intensity deviations of OI 557.7 nm (normalized to their mean intensity values) and (b) plots the mean intensity deviation of OH (normalized to their mean intensity value). A best fit to OH data is obtained. The periodicity from OI emission is ~ 8.4 hr and that from OH line it comes out to be 8 hr. The residual (obtained by subtracting from this long period oscillations) intensity deviations for OI and OH are shown in figure 2(c) and (d) respectively. The wave periods comes out to be 4.6 and 4.2 hr. Similarly figure 3(a, b) shows the average variation of intensity deviations for OI and OH emissions on 10–13 April, 2007 and obtained period is 8.4 hr for both the lines. However, the residual mean intensity deviations shown in figure 3(c, d) yield wave period 2.2 and 4 hr respectively for OI and OH emissions.

The nocturnal mean intensity deviations of OI 557.7 nm and OH (7, 2) emissions for the night of February 14–15, 2004 are shown in the upper panel, along with the contributions from the first and the second harmonics (9 and 4.5 hr periods) in the lower panels of figure 4. Here it should be noted that although we had data for 9 hr only, the cosine fitting for the data yielded 9 hr of periodicity. The harmonic analyses provide the information of percentage contributions of different harmonics in the total wave structure. In the present case, the maximum percentage contribution comes from the second harmonics in both OI and OH emissions rather than the other harmonics. Because of this we chose the second harmonics for the comparison of OH and OI 557.7 nm emissions. The plots of two emissions are shown in the same figure to get the location of amplitude minima of wavy structure coming from the two different layers and it is inferred that the longer period wave (~ 9 hr) propagates upward. The mean intensity averaging process for 13–16 February, 2004, yield 8 hr period for the OI emission and 9.6 hr for the OH emissions. The wave period for the shorter wave comes out to be 2.4 and 4.4 hr for the OI and OH emissions respectively (table 1). Table 1 summarises the results of analysis based on an average mean intensity deviations of three to four nights depending upon the availability of data.

The mean intensity deviations of OI 557.7 nm and OH emissions for the night of February 15–16, 2004, along with the variation of amplitude of different harmonics (9, 4.5 hr) are shown in figure 5. The phase propagation for the first harmonic seems to be downwards, whereas the second harmonic shows an upward propagation. The data for February 17–18, 2004, was analyzed (figure 6) and phase propagation for both the first and the second harmonics was found to be upward. The analysis of average mean intensity deviations for the period 13–16 February and 17–20 February 2004 show downward phase propagation both for the first and the second harmonics. In normal condition, downward wave propagation is not tenable. Such propagation in localized region may be observed due to the strong wind effect or the influence of multi-wave mode coupling. Analysing OI and OH emission data, Fagundes *et al.* (1995) showed downward phase propagation with velocity in the range $4\text{--}10\text{ ms}^{-1}$ and vertical wavelength lying between 12 and 77 km. Meriwether *et al.* (1998) using Rayleigh lidar observations showed a downward phase propagation of 1 km/hr and explained it through a localized mechanism involving the coupling of gravity waves to the mesopause tidal structure. They suggested that in a critical layer in which phase speed of the propagating gravity wave matches the combined speeds of background mean wind and

Figure 2. Nocturnal variation of mean intensity deviation of the four days average of OI 557.7 nm and OH emissions from 12–15 March 2007 observations are plotted. The upper panels represent the mean deviations in intensity of (a) OI 557.7 nm and (b) OH emission. The bottom panels represent intensity residuals of (c) OI 557.7 nm and (d) OH emission. Solid red line curve in each plot shows the result of best-fit cosine model which is used to estimate the dominant wave periodicity, amplitude and phase of waves present in the data. IST is the Indian Standard Time.

Figure 3. Same as figure 2, but for 10–13 April, 2007.

Figure 4. The nocturnal mean intensity deviations of OI 557.7 nm and OH emissions for the night of February 14–15, 2004 observed at Kolhapur are plotted in the upper panel along with the contributions of the first and second harmonics (9 and 4.5 hr period) in the lower panels. IST is the Indian Standard Time.

the tidal mode, the convective/shear instability causes wave energy to be transferred from gravity wave to mean flow with consequent changes in the tidal wind amplitude and phase. Recently, clear evidence of planetary wave modes penetrating to OH layer heights at different times have been observed (French and Klekociuk 2011) which may interact with gravity waves and tidal modes and change their amplitude and phase. However, in the present case we could not find any clear signature of planetary wave. Another possibility could be that the waves may be moving vertically slow in the presence of strong wind. In that case very large vertical wavelength should be observed. The vertical wavelength in the present case lies between 46 and 113 km (see table 1). The vertical wavelength in the presence of wind is inversely proportional to the difference in phase velocity of the wave and the projection of the background wind on the direction of wave propagation (Akmaev 2001). Thus the vertical wavelength depends on wind direction and magnitude of wind velocity.

Similarly, the mean intensity deviations of OI 557.7 nm and OH emissions for the night of February 10–11, 2005 are shown in the upper panels of figure 7 along with the contributions from the first and the second harmonics (9.4 and 4.6 hr

periods) in the lower panels. The first and the second harmonics show downward phase propagation. Another example of the night airglow mean intensity deviations for March 8–9, 2005 along with the harmonic analysis is shown in figure 8.

Figure 9 shows the corresponding data for the night of March 10–11, 2005. From the figure we notice that the first harmonics propagate upward and the second harmonics are seen to be opposite to each other. From table 1, we find that the analysis of average mean intensity data yield wave period of the first harmonics between 3 and 12 hours. The wave period of the second harmonics varied between 2 and 4.6 hours. These waves may be attributed to the tidal oscillations.

The vertical wavelength of the wave is computed using the relation (Fagundes *et al.* 1995):

$$\lambda(z) = \Delta D \times (360/\Delta\Phi) \quad (1)$$

where ΔD is the separation of the two layers and $\Delta\Phi$ is the phase difference between the waves corresponding to the two emission lines. For waves with periods much larger than the Brunt Vaisala period, T_B (say greater than 1 hr and less than a few hours), the horizontal wave length λ_x and the vertical wavelength λ_z of internal gravity waves are

Table 1. Three/four days average intensity data are used to determine wave period, time shift and phase shift. These are used to determine vertical velocity and vertical wavelengths which are tabulated below*.

Date	Wave period 557.7 (hr)	Wave period OH (hr)	ΔT (min)	$\Delta \Phi$ (deg)	V(m/s)	vw (km)
January 16–19, 2004	8.4	7.2	42.60	35.50	4.30	111.55
January 16–19, 2004	4.4	4.0	92.40	138.60	1.98	28.57
January 23–30, 2004	8.0	8.0	60.30	45.23	3.04	87.56
January 23–30, 2004	3.6	3.6	75.96	126.60	2.41	31.28
February 13–16, 2004	8.0	9.6	114.18	85.5	1.66	46.32
February 13–16, 2004	2.4	4.4	31.10	45.82	5.85	86.40
February 17–20, 2004	8.4	12.0	74.82	52.81	2.45	74.98
February 17–20, 2004	4.4	2.8	16.26	34.84	11.28	113.65
March 17–22, 2004	8.0	7.6	33.00	24.75	5.56	160.00
March 17–22, 2004	3.6	2.8	29.22	62.61	6.27	63.24
December 18–22, 2004	8.0	8.0	32.40	24.30	5.66	162.96
December 18–22, 2004	3.2	2.4	33.60	84.00	5.46	47.14
December 29–31, 2004	12.0	12.0	49.56	24.78	3.70	159.81
December 29–31, 2004	2.8	4.0	57.00	85.50	3.22	46.32
February 3–5, 2005	8.0	8.0	38.52	28.89	4.76	137.07
February 3–5, 2005	4.0	4.0	61.92	92.88	2.96	42.64
February 8–10, 2005	6.2	6.2	57.90	56.03	3.17	70.67
February 8–10, 2005	4.0	3.8	79.02	124.77	2.32	31.74
March 6–8, 2005	6.0	6.2	45.60	42.75	4.02	92.63
March 6–8, 2005	3.2	3.4	38.40	72.00	4.77	55.00
March 9–11, 2005	7.4	3.0	75.60	61.30	2.43	64.60
March 9–11, 2005	4.0	2.0	88.92	133.38	2.06	29.69
March 12–15, 2007	8.4	8.0	53.40	38.14	3.43	103.82
March 12–15, 2007	4.6	4.2	33.60	50.40	5.46	78.57
April 10–13, 2007	8.4	8.4	52.80	39.60	3.47	100.00
April 10–13, 2007	2.2	4.0	49.80	71.14	3.68	55.66

*All the parameters have been calculated from the method of Fagundes *et al.* (1995).

vw - Vertical wavelength.

Figure 5. Same as figure 4, but for February 15–16, 2004.

Figure 6. Same as figure 4, but for February 17–18, 2004.

Figure 7. Same as figure 4, but for February 10–11, 2005.

Figure 8. Same as figure 4, but for March 8–9, 2005.

Figure 9. Same as figure 4, but for March 10–11, 2005.

related through the relation (Hines 1960; Fagundes *et al.* 1995)

$$\frac{\lambda_x}{\lambda_z} = \frac{T_W}{T_B} \quad (2)$$

where T_W is the wave period. This simple relation changes for very long and very short period waves (Fagundes *et al.* 1995). The time period (T in min) and phase difference $|\Delta\phi|$ obtained from the cross correlation analysis of the two emission lines, OI 557.7 nm and OH (7, 2) have been used to compute the vertical wavelength which varies between approximately 28.7 and 113.6 km (short period, S) wave propagation and it varies between 46.3 and 160.9 km (long period, L) wave propagation. Values of the vertical wavelength show significant day-to-day variability (~ 28.5 – 160.9 km) with an average value of ~ 54.6 km for the short period waves and ~ 105 km for the long period wave. The source of these waves may be in the troposphere from which the gravity waves propagate upward in the middle atmosphere. However, *in situ* excitation at the airglow altitudes or ducting in the upper middle atmosphere has also been suggested by observations (Isler *et al.* 1997; Hecht *et al.* 2009) and numerical modelling (Fritts *et al.* 1997; Hertzog *et al.* 2002).

The vertical wavelength obtained by several investigators matches more or less with our observed values. Reddi and Ramkumar (1997) reported significantly higher values of the vertical wavelength (>100 km) with meteor wind observations at Trivandrum. Taori *et al.* (2005) obtained a large variability of vertical wavelength (~ 26 – 110 km) with the mean value ~ 90 km which resembles very well with our results. Contrary to the large variability in the vertical wavelength of tidal oscillations, Deepa *et al.* (2006) using Meteor Radar data collected at Trivandrum showed the vertical wavelength of the terdiurnal tide ~ 16 – 50 km during various seasons with the maximum amplitude at 95 km altitude, a result consistent with the Meteor Radar measurement at the high latitudes (Wu *et al.* 2005). These results clearly show that airglow emission data yield the vertical wavelength consistently higher values than those derived from the Meteor Radar measurements. To resolve this paradox, simultaneous measurements using both the techniques are recommended. Moreover, in the analysis of airglow data one considers emission of a particular line to come from a fixed altitude, which itself may vary due to variation in gas compositions caused by various factors.

We have also computed the correlation coefficients (r) between OI 557.7 nm and OH (7, 2)

emission intensities for the short and the long wave periods observed on 54 nights during the period of 2004–2007. It is noted that there is a high degree of correlation between the two sets of data of two emission lines. The correlation coefficient varies between 0.763 and 0.872. The high values of correlation also predict very high values of the vertical wavelength.

The estimated phase velocities for the short period waves between OH and OI 557.7 nm layers vary between 2.0 and 12.0 m/s and the time lag varies between 16.26 to 92.4 min. The long period wave phase velocities vary between 1.95 and 18.9 m/s and time lag varies between 32.4 and 114.18 min of the two emissions OI 557.7 nm and OH. The data was dominated by a terdiurnal tide-like wave. From table 1, it is noticed that the average vertical velocity of 10 diurnal oscillations is 3.4 m/s in 2004, 3.6 m/s in 2005 and 3.4 m/s in 2007. The corresponding average vertical wavelength for the years 2004, 2005 and 2007 is 107.2, 100.9 and 101.9 km, respectively. Thus the average vertical velocity and wavelength does not show appreciable variation on yearly basis.

4. Tidal influences in airglow

Figure 10(a, b), shows the cosine fit curve of OI 557.7 nm and OH(7, 2) emissions for the night of February 14–15, 2004. Solid line shows the best fit curve along with the 15 min running average of the mean airglow intensity deviations. It is seen that the period of OI 557.7 nm intensity variation is ~ 6 hr and OH emission period is ~ 8 hr. The residuals of mean intensity deviations shown in the lower panels (c, d), which are consistently spread around zero yielding the validation of cosine model. Periodicity of tidal wave at OH emission is more dominant than OI 557.7 nm emission. The same results for February 15–16, 2004 are shown in figure 11 where the solid line shows the cosine fit curve. The period of OI 557.7 nm mean intensity deviation is ~ 4.5 hr and for the OH emission it is ~ 8 hr. The residuals are plotted in the lower panels (c, d). From the figure it is clearly seen the presence of small period waves in OI 557.7 nm data, which is also evident from the plot of residuals. Apart from tidal oscillations, short period oscillation observed in the OI 557.7 nm emissions could be due to gravity wave. Figure 12(a, b) also shows another example of terdiurnal oscillations in the OI 557.7 nm and OH emissions data observed on February 19–20, 2004.

Analyzing nightglow emission data, we have shown the presence of terdiurnal (8 hr) tides during the

Figure 10. Nocturnal variation of mean intensity deviations in OI 557.7 nm and OH emissions for the night of February 14–15, 2004 at Kolhapur are presented. The upper panels represent the mean deviations in intensity of (a) OI 557.7 nm (b) OH emissions. Solid blue line curves in (a) and (b) plots show the result of best cosine fit curve. Long period tide-like oscillation is evident in the data with periodicity of about 6 hr in OI 557.7 nm and 8 hr in OH emissions. The airglow intensity residuals are plotted in the lower panel (c, d). IST is the Indian Standard Time.

Figure 11. Same as figure 10, but for February 15–16, 2004.

Figure 12. Same as figure 10, but for February 19–20, 2004.

month of February, 2004. The observed difference in the periodicity of OH and OI 557.7 nm emissions may be due to the wave-mean wind interaction between these two altitudes. It is observed from the plots that OH is having higher wave activity during the observation period than OI emissions. Thayaparan (1997) reported significant contribution of 8 hr tide during the spring and winter months and less contribution during the summer and fall months. At mid-latitude, the amplitude of 8 hr tides maximizes during spring and fall months (Smith *et al.* 2000). Taori and Taylor (2010) reported dominance of 8 hr tides at the low latitude station Maui, Hawaii during the late winter time, which is close to our findings. Using satellite based wind observations, Zhao *et al.* (2005) showed a quite different seasonal pattern of 8 hr tides.

The simultaneous measurements of various airglow emissions emanating from different altitudes in the atmosphere on a given night provide information about the vertical propagation of waves through the region. In these preliminary results we have shown the presence of waves of different periods (table 1) including the determinant terdiurnal tidal modes. Waves with periods of a few hours could be gravity waves. Multispectral mesosphere nightglow emissions have reported gravity waves with periods of a few hours propagating through the upper atmosphere in the low latitude regions (Molina *et al.* 1985; Takahashi *et al.* 1985;

Fagundes *et al.* 1995). Taori and Taylor (2010) have also shown the presence of terdiurnal long period waves and short period waves in the mesospheric O₂ airglow emissions. The wave growth or damping during propagation could not be studied due to the limited availability of data and the absence of multi-instrument observations. The vertical upward propagation of wave suggests that the source region may lie in the lower atmosphere, i.e., in the troposphere or the stratosphere. The source of these waves could be: (a) tidal oscillation caused by absorption of solar radiation in the lower atmosphere, (b) gravity waves generated by different processes in the troposphere, (c) wave-wave or wave-tide interactions and development of instabilities leading to generation and amplification of waves, and (d) interaction between different modes may also give rise to other modes. For example, interaction between semi-diurnal and diurnal tides may result into the generation of terdiurnal waves (Taylor *et al.* 1999). The interaction of tidal wave with gravity waves also generates new wave modes (Thayaparan *et al.* 1995). Non-linear interactions between the stationary planetary waves and migrating tides may also lead to the diurnal tide (Hagan and Roble 2001) and the semidiurnal tides (Angelats i Coll and Forbes 2002). Using modelling studies (Roble and Shepherd 1997; Yudin *et al.* 1998; Zhang *et al.* 2001), the perturbation in airglow emission have also been linked to tidal oscillations.

5. Summary and conclusions

The analysis of night time airglow emission lines OH(7, 2) and OI 557.7 nm showed the presence of terdiurnal and other sub-harmonics of solar tidal oscillations. The observed short period waves may be the manifestation of tidal wave interactions amongst themselves or their interactions with winds, gravity waves, planetary waves, etc. In the present analysis, direction of propagation, wavelength and wave periods are determined. The waves with wave periods of a few hours usually should have origin in the lower atmosphere and wave propagation direction upward. In a few cases we did observe downward wave propagation. Such cases may be the result of local strong wind effect which modified the amplitude and phase of the waves. In the absence of wind data, it is difficult to understand the propagation mechanism. More data analysis and simulation studies with different direction and magnitude of wind are required to resolve this issue. Some of the results are:

- The vertical velocities of the long and short period waves propagating between OH and OI 557.7 nm layers vary randomly on day-to-day basis. On some days, velocity of the long period wave is large and on the other day velocity of the short period wave is large. No clear cut picture could be obtained. Similar results were obtained in 2004, 2005 and 2007.
- The 8 hr terdiurnal component of atmospheric tides are often present in OH (7, 2), and OI 557.7 nm emissions. The average velocity and wavelength did not show appreciable change during the observation period from 2004 to 2007. The observed difference in periodicity of the wave may be due to the wave mean wind interactions between the two airglow layers.
- Good correlations are found between the OH and OI 557.7 nm emission intensity on most of the nights for the short period and the long period variations.
- The wave-wave interactions could be seen in OH emission layer.

In the present study, we have taken the height of OH and OI emission layer to be 86 and 97 km, respectively, which may change from time to time due to altitude variation in atmospheric constituents. The variation will change the derived wave parameters. In the present study we could not decipher the origin of short period waves. Additional observations using different techniques are required to explain the wavy structures observed in the airglow emission data. Hence, further long term multi-instrument observations (using ground-based and space-based), theoretical studies and modelling

are required to deal with this less explored and the most important region of the atmosphere.

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