



Octanuclear $\{Ln_8\}$ complexes: magneto-caloric effect in the $\{Gd_8\}$ analogue

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Abstract. Two neutral, isostructural, octanuclear, Ln^{III} complexes, $[Ln_8(HL)_6(L)_2(\mu_3-OH)_4(\mu_2-OH)_2(H_2O)_4].solv$ ($Ln = Gd(III)$, (1^{Gd}) and $Dy(III)$, (2^{Dy}) have been synthesized using Ln^{III} nitrate salts and a *o*-vanillin-supported multi-dentate ligand, *N'*-(2-hydroxy-3-methoxy-5-nitrobenzylidene)-2-(hydroxyamino)propanehydrazide (H_3L) in the presence of tetramethyl ammonium hydroxide. The complexes were structurally characterized by single crystal X-ray diffraction studies. The complexes are held by the cumulative coordination action of six $[HL]^{2-}$, two $[L]^{3-}$ chelating ligands, H_2O and hydroxide ligands. Magneto-structural analysis of complexes 1^{Gd} and 2^{Dy} reveals the presence of intramolecular antiferromagnetic interactions between the Ln^{III} ions. Magneto-caloric effect was analysed for the complex 1^{Gd} which shows a maximum in the change of molar entropy ($-\Delta S_m$) of magnitude $25.5 \text{ J kg}^{-1} \text{ K}^{-1}$ at $T = 3 \text{ K}$ and an applied field change $\Delta B = 5 \text{ T}$.

1. Introduction

Lanthanide complexes have been of considerable research interest in recent years owing to their interesting magnetic,^{1–10} optical,^{11–20} and catalytic properties.^{21–28} Among these systems, polynuclear complexes are formed usually as a result of the bridging ligands such as the hydroxide ligand.^{29–39} Such complexes are also of interest in the field of molecular magnetism.^{30,31,40–45} However, the synthesis of such complexes is fraught with some challenges. One difficulty is the formation of intractable polymeric complexes because of the ready de-protonation of the coordinated water molecules around the lanthanide ions. The formation of discrete complexes becomes

feasible if controlled de-protonation of the coordinated water molecules can be achieved.²⁹ This is often possible by having an appropriate multi-site coordinating ligand to bind to the lanthanide ions so that the available coordination sites are reduced and the chances for the formation of a discrete complex are increased. In our lab, multi-pocket, multidentate hydrazone ligands have been extensively used for the synthesis of discrete polynuclear lanthanide complexes.^{46–50} In this regard, we have designed a multi-pocket hydrazone ligand which was used to isolate dinuclear Ln^{III} complexes.⁵¹ In these complexes, we have observed uncoordinated free oxime $-OH$ groups which could be potentially deprotonated so that these can proliferate to form multinuclear complexes.

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Accordingly, herein, we report the synthesis, structures and magnetism of two octanuclear complexes, $[\text{Ln}_8(\text{HL})_6(\text{L})_2(\mu_3\text{-OH})_4(\mu_2\text{-OH})_2(\text{H}_2\text{O})_4]\cdot\text{solvents}$ ($\text{Ln}^{\text{III}} = \text{Gd}^{\text{III}}$, (**1^{Gd}**) and Dy^{III} , (**2^{Dy}**)). Interestingly, we did not use any additional co-ligand in isolating these compounds, which are often required to isolate such multinuclear Ln^{III} complexes.^{52–55}

2. Experimental

2.1 Materials and methods

Solvents and other general reagents used in this work were obtained from commercial sources and used without further purification. 3-Methoxy-5-nitrosalicylaldehyde, hydroxylamine hydrochloride and lanthanide nitrate hydrates were obtained from Sigma Aldrich Chemical Co. (U.S.A.), Ethyl pyruvate and tetramethylammonium hydroxide pentahydrate were obtained from Spectrochem Pvt. Ltd., Mumbai, India. All these reagents were used as obtained without further purification. Ethyl (2*E*)-2-(hydroxyimino)-propanoate, (2*E*)-2-(hydroxyimino)-propanehydrazide, and the ligand H_3L were prepared by a previously reported procedure.^{51,56,57}

2.1a Instrumentation: Melting points were measured using a StuartTM SMP10 melting point apparatus and are uncorrected. IR spectra were recorded on a Perkin Elmer FT-IR spectrometer. Mass spectra were recorded with a Bruker micrOTOF-Q II spectrometer. Elemental analyses of the compounds were obtained from a Euro Vector EA instrument (CHNS-O, model EuroEA3000).

2.1b Magnetic measurements: Field dependence of the magnetization at different fields and variable temperature, as well as magnetic measurements in the 2–300 K range were carried out on polycrystalline samples with a Quantum Design SQUID MPMS XL-5 device operating at different magnetic fields. AC susceptibility measurement of **2^{Dy}** were performed using a frequency of 1400 Hz under magnetic fields of zero and 0.1 T. The static experimental susceptibilities were corrected for the diamagnetism of the constituent atoms (using Pascal's constants) and for the sample holder. To avoid any torquing of the microcrystals, a pellet of the sample cut into very small fragments were introduced in the sample holder.

2.1c X-ray crystallography: The molecular structures of complexes **1^{Gd}** and **2^{Dy}** were analysed

by single-crystal X-ray diffraction studies performed on a Bruker APEX-II CCD diffractometer equipped with an Oxford low-temperature attachment. The crystals were kept at 120 K during data collection. The frames were indexed, integrated, and scaled using the SMART and SAINT software packages.⁵⁸ Absorption correction was performed by a multi-scan method implemented in SADABS.⁵⁹ Space groups were determined using XPREP implemented in APEX-II.⁶⁰ The structures were solved by ShelXT⁶¹ structure solution programme using Intrinsic Phasing and refined with the ShelXL⁶² refinement package using Least Squares minimization on F^2 in the Olex-2 software.⁶³ All the non-hydrogen atoms were refined anisotropically using full-matrix least-square procedures. Hydrogen atoms on all the bridging hydroxides were not observed in the diffraction pattern and therefore omitted entirely, although their oxidation states were confirmed by BVS calculations (see Table S1, SI).^{64,65} The crystal data and the cell parameters for **1^{Gd}** and **2^{Dy}** are summarized in Table 1. The structures also contain heavily disordered solvents of crystallization which could not be modelled satisfactorily. Therefore, PLATON/SQUEEZE was used to get rid of the contributions of the disordered solvent molecules.⁶⁶ All of the mean plane analyses and molecular drawings were produced using the DIAMOND software (version 3.2).⁶⁷

2.1d Synthesis of metal complexes: A general synthetic protocol was utilized for preparation of the octanuclear metal complexes as follows:

To a solution of H_3L in 20 mL of $\text{MeOH}/\text{CH}_3\text{CN}$ (1:1 v/v), tetramethylammonium hydroxide was added. After stirring for half an hour, $\text{Ln}(\text{NO}_3)_3 \cdot x\text{H}_2\text{O}$ ($x = 6$ for **1^{Gd}** and $x = 5$ for **2^{Dy}**) was added to this solution which resulted in a clear yellow coloured solution. After stirring for another 12 h at room temperature the reaction mixture was filtered and kept undisturbed for crystallization. Slow evaporation of the filtrate under ambient conditions resulted in the formation of yellow block-shaped crystals suitable for X-ray diffraction after two weeks. The crystals were collected, washed with MeOH and dried under vacuum for 16 h resulting in polycrystalline powders. The specific quantities of the reactants involved in each reaction, yields of the products, and their characterization data are given below.

$[\text{Gd}_8(\text{HL})_6(\text{L})_2(\mu_3\text{-OH})_4(\mu_2\text{-OH})_2(\text{H}_2\text{O})_4] \cdot 2\text{CH}_3\text{CN} \cdot 7\text{CH}_3\text{OH} \cdot 4\text{H}_2\text{O}$ (**1^{Gd}**). Quantities: H_3L (0.061 g, 0.205 mmol), $\text{Gd}(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$ (0.092 g, 0.205 mmol), Me_4NOH (0.112 g, 0.616 mmol). Yield: 0.066 g (49%

Table 1. Crystallographic details for complexes **1^{Gd}** and **2^{Dy}**.

	1^{Gd}	2^{Dy}
Empirical formula*	C ₈₈ H ₈₆ Gd ₈ N ₃₂ O ₅₈	C ₈₈ H ₈₆ Dy ₈ N ₃₂ O ₅₈
Formula weight*	3777.88	3819.88
Temperature/K	120(2)	120(0)
Crystal system	Monoclinic	Monoclinic
Space group	C2/c	C2/c
a/Å	19.1263(18)	18.9896(19)
b/Å	31.7579(18)	31.583(3)
c/Å	25.4695(16)	25.239(3)
α/°	90	90
β/°	98.198(5)	98.325(5)
γ/°	90	90
Volume/Å ³	15312.4(19)	14978(3)
Z	4	4
ρ ^{calc} g/cm ³	1.639	1.694
μ/mm ⁻¹	3.500	4.027
F(000)	7256.0	7320.0
Crystal size/mm ³	0.21 × 0.15 × 0.11	0.13 × 0.11 × 0.08
Radiation	MoKα (λ = 0.71073)	MoKα (λ = 0.71073)
2θ range for data collection/°	2.504 to 51.332	3.052 to 57.938
Index ranges	– 22 ≤ h ≤ 23, – 38 ≤ k ≤ 38, – 30 ≤ l ≤ 24	– 25 ≤ h ≤ 25, – 42 ≤ k ≤ 41, – 30 ≤ l ≤ 33
Reflections collected	85743	114164
Independent reflections	14325 [R _{int} = 0.1923]	18900 [R _{int} = 0.2075]
Data/restraints/parameters	14325/1/630	18900/2/797
GOF on F ²	1.029	0.948
Final R indexes [I ≥ 2σ (I)]	R ₁ = 0.1065, wR ₂ = 0.2123	R ₁ = 0.0660, wR ₂ = 0.1317
Final R indexes [all data]	R ₁ = 0.2235, wR ₂ = 0.2758	R ₁ = 0.1692, wR ₂ = 0.1690
CCDC Number	2064754	2064755

*including H atoms of the bridging hydroxyl groups as confirmed by BVS calculations.

based on Gd). M.p.: >250 °C. IR (KBr) (cm⁻¹): 3425(br), 3245(w), 2941(w), 2839(w), 1605(s), 1570(m), 1493(s), 1383(s), 1360(w), 1305(s), 1256(m), 1203(m), 1101(s), 1060 (m), 975(w), 911(s), 842(w), 777(m), 746(m), 725(w), 707(w), 552(w). Calcd. for C₈₈H₈₆Gd₈N₃₂O₅₈ (MW = 3777.82): C, 28.61; H, 3.10; N, 11.46. Found: C, 28.42; H, 2.89; N, 11.17 ESI-MS: m/z = 1868.9176, [Gd₈(HL)₆(L)₂(OH)₆ + Na⁺ + H⁺]²⁺.

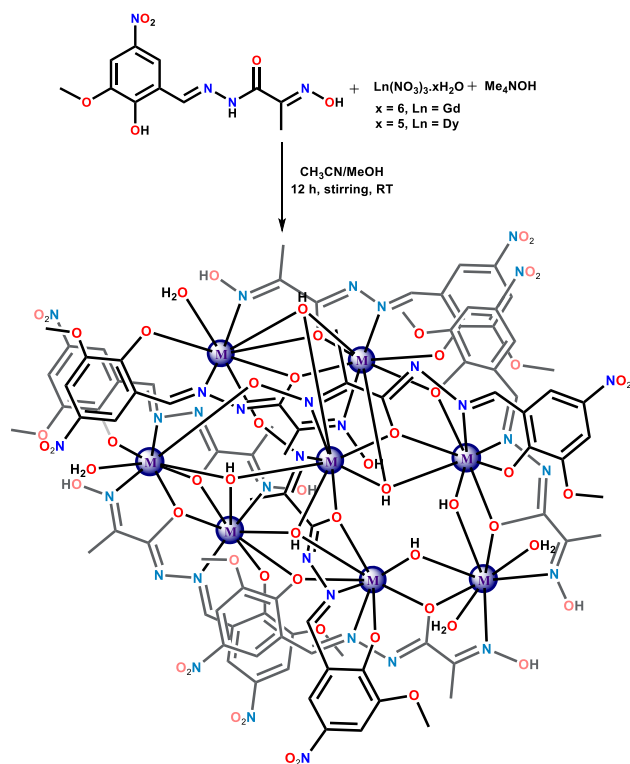
[Dy₈(HL)₆(L)₂(μ₃-OH)₄(μ₂-OH)₂(H₂O)₄].CH₃CN. 5CH₃OH.9H₂O (**2^{Dy}**). H₃L (0.061 g, 0.205 mmol), Dy(NO₃)₃·5H₂O (0.090 g, 0.205 mmol), Me₄NOH (0.111 g, 0.616 mmol). Yield: 0.066 g (49% based on Dy). M.p.: >250 °C. IR (KBr) (cm⁻¹): 3417(br), 3235(w), 2939(w), 2837(w), 1603(s), 1568(m), 1491(s), 1381(s), 1360(w), 1306(s), 1254(m), 1201(m), 1101(s), 1058(m), 975(w), 909(s), 842(w), 777(m), 746(m), 726(w), 705(w), 552(w). Calcd for C₈₈H₈₆Dy₈N₃₂O₅₈ (MW = 3819.82): C, 27.28; H, 3.06; N, 11.05. Found: C, 27.06; H, 2.81; N, 10.77.

ESI-MS: m/z = 1889.9365, [Dy₈(HL)₆(L)₂(OH)₆ + Na⁺ + H⁺]²⁺.

3. Results and Discussion

3.1 Synthetic aspects

The Schiff base hydrazone ligand, *N'*-(2-hydroxy-3-methoxy-5-nitrobenzylidene)-2-(hydroxylamino)propanehydrazide (H₃L), contains six potential coordinating sites capable of binding to multiple lanthanide metal ions in different coordination modes as shown in Scheme 1. The role of the hydroxide base is crucial not only in deprotonating H₃L but also for being a source of hydroxide ligands essential for stitching the multiple metal centres together in the complexes without the need for additional co-ligands. Thus, the reaction of H₃L with hydrated Ln(NO₃)₃·xH₂O and Me₄NOH·5H₂O in the 1:1:3 molar ratio in



Scheme 1. Synthesis of complexes **1^{Gd}** and **2^{Dy}**.

MeOH/CH₃CN (v/v) solvent mixture allowed isolation of [Ln₈(HL)₆(L)₂(μ₃-OH)₄(μ₂-OH)₂(H₂O)₄] complexes (Scheme 1).

ESI-MS studies of **1^{Gd}** and **2^{Dy}** revealed that all the complexes retain their molecular ion peaks at *m/z* = 1868.9176 and 1889.9365 corresponding to the dicationic fragments [Gd₈(HL)₆(L)₂(OH)₆ + Na⁺ + H⁺]²⁺ and [Dy₈(HL)₆(L)₂(OH)₆ + Na⁺ + H⁺]²⁺. The ESI-MS of **2^{Dy}** is shown in Figure 1 while that of **1^{Gd}** is shown in Figure S1, SI.

3.2 Molecular structures

The molecular structures of the complexes **1^{Gd}** and **2^{Dy}** were confirmed by the single crystal X-ray studies. Both the complexes are charge neutral and crystallized in the monoclinic crystal system in the *C2c* space group (*Z* = 4). The octanuclear complexes comprise of eight Ln^{III} cations, six [HL]²⁻, two [L]³⁻, four [μ₃-OH]⁻, two [μ₂-OH]⁻ anions, and four H₂O molecules (Figure 2). Considering the overall structural similarity of these two complexes, we have chosen **1^{Gd}** to elucidate the common structural features present in them. The molecular structure of **2^{Dy}** is shown in Figure S2 (SI) and the corresponding interatomic distance parameters are given in Table S5 (SI).

The molecular structure of **1^{Gd}** and its asymmetric unit are given in Figure 2. A *C*₂ axis passes through the line connecting the Gd1 and Gd5 which equally divides the whole molecule into two halves generating the asymmetric unit. The formation of the octanuclear complex can be understood as follows. The whole molecular structure can be viewed as being composed of (i) a tetranuclear core (Figure 3(a)) and (ii) two structurally similar dinuclear motifs (Figure 3(b)). In the dinuclear cores, the two [HL]²⁻ ligands hold two Gd^{III} ions in a ‘head-to-tail’ fashion utilizing a tridentate (O, N, O) and a bidentate (N, O) coordinating motif. The enolate oxygen atoms (O2 and O2*) of the ligand bridge the two metal centres affording an approximate rhomboidal-shaped four-membered Gd₂O₂ core. The average Gd⋯Gd distance and average Gd–O_{enol}–Gd angle in the Gd₂O₂ cores are found to be 3.677(19) Å and 103.365(62)°, respectively. Each of the Gd^{III} ions in the dinuclear units is eight-coordinate with an overall 2N, 6O coordination environment. In the tetranuclear unit, one [HL]²⁻ ligand and one [L]³⁻ ligand hold three Gd^{III} ions in a ‘head-to-head’ fashion. The bridging coordination action of the enolate oxygen atoms and the hydroxyl atoms between the four metal centres results in a Gd₄O₈ core where the four metal ions are present in a *kite*-shaped geometry. The average Gd⋯Gd distance and the average Gd–O_{enol}–Gd angle in this core are found to be 3.831(18) Å and 107.071(59)°, respectively. The overall coordination environment around the metal ions in this core is also 2N, 6O similar to the coordination environment observed in the dinuclear units. The cumulative coordination action of the hydroxyl anions and the ligand oxime oxygen atoms connects the two dinuclear units with the tetranuclear unit completing the octanuclear {Gd₈} complex. Selected interatomic distances and bond angle parameters of **1^{Gd}** are given in Table S2 (SI).

A view of the {Gd₈} core with only the bridging atoms is shown in Figure 4 (left). Mean plane analysis reveals that the *kite*-shaped {Gd₄} unit is perfectly planar comprising the Gd1, Gd2, Gd2*, and Gd5 atoms (*plane 1*). Another two planes corresponding to the dinuclear units i.e., *plane 2* (Gd1, Gd3, Gd4) and *plane 3* (Gd1, Gd3*, Gd4*) are at a dihedral angle of 62.97(3)° with respect to *plane 1* while the *plane 2* and *plane 3* bisect one another with a dihedral angle of 33.57(5)° (Figure 4 (right)).

The immediate coordination geometries around all the eight Gd atoms were confirmed by the SHAPE analysis.^{68,69} The geometries were obtained as square antiprism (Gd1, Gd2, Gd2*), Johnson gyrobifastigium (Gd3, Gd3*), biaugmented trigonal prism (Gd4,

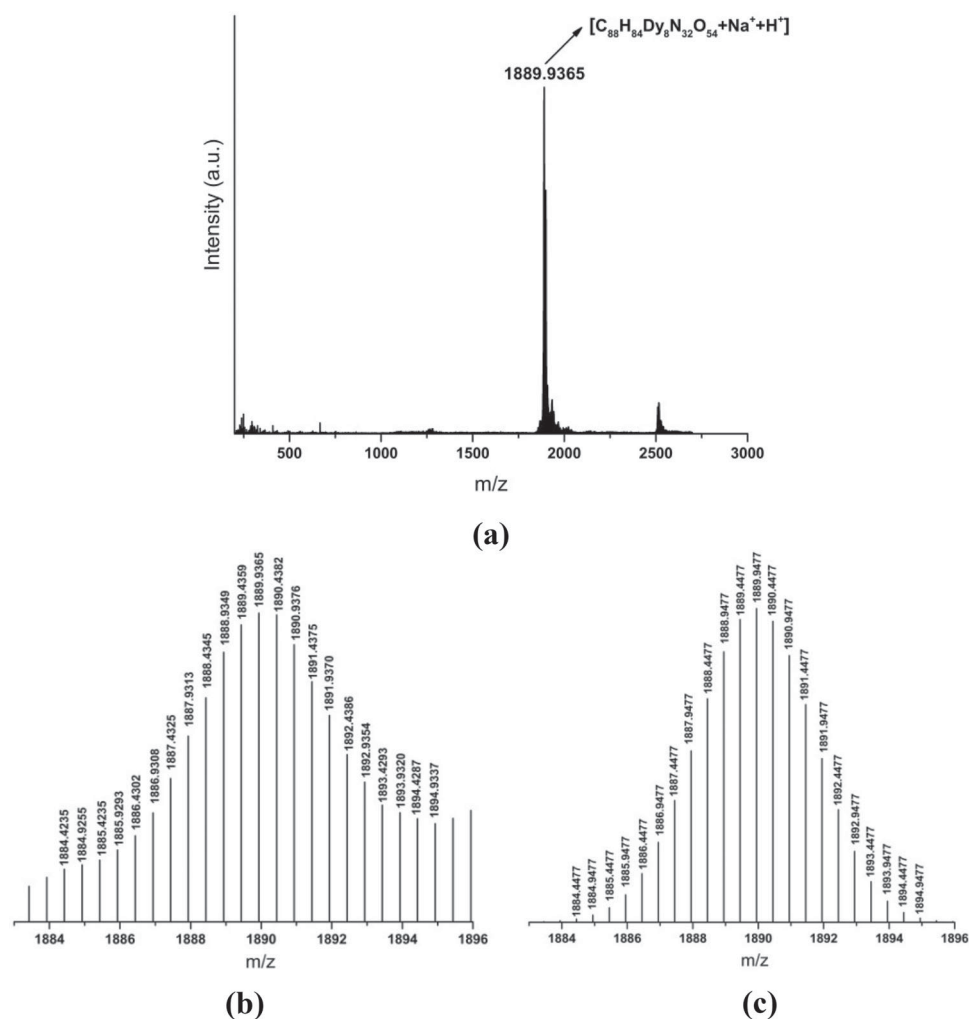


Figure 1. (a) Full range ESI-MS spectrum of complex 2^{Dy} . (b) Experimental and (c) Simulated pattern of $[Dy_8(HL)_6(L)_2(OH)_6 + Na^+ + H^+]^{2+}$.

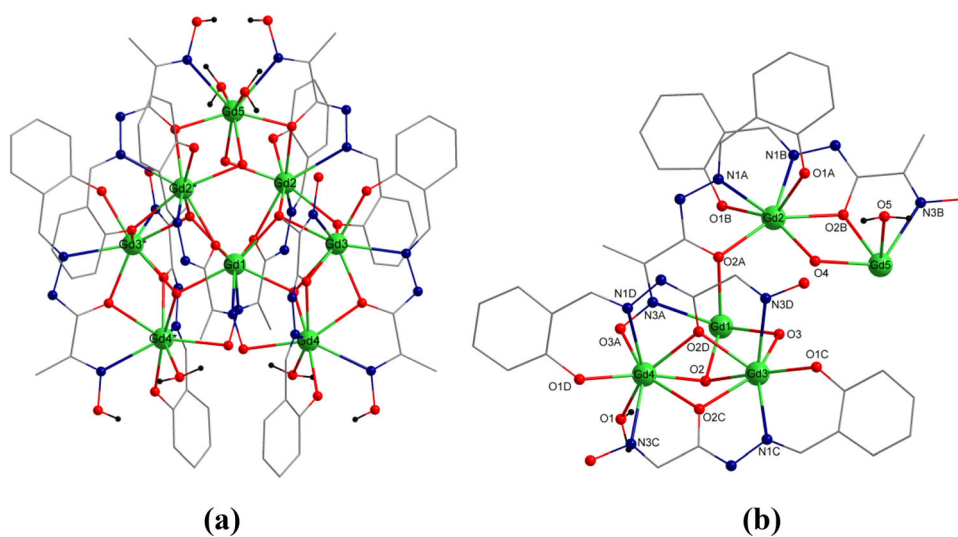


Figure 2. (a) The molecular structure of 1^{Gd} ($-OMe$, $-NO_2$ groups, and H atoms except in the water molecules are omitted for clarity) and (b) the asymmetric unit of the complex Colour scheme: Gd^{III} , light green; O, red; N, blue; C, gray.

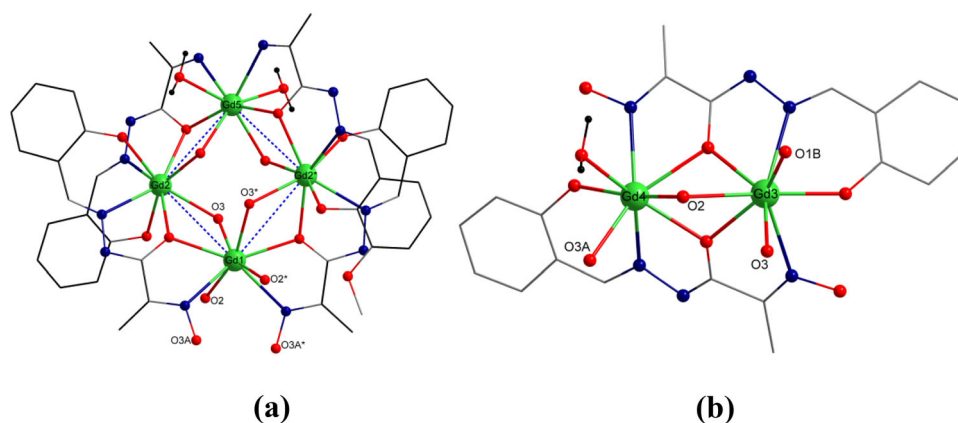


Figure 3. (a) The tetranuclear core and (b) the dinuclear core of 1^{Gd} .

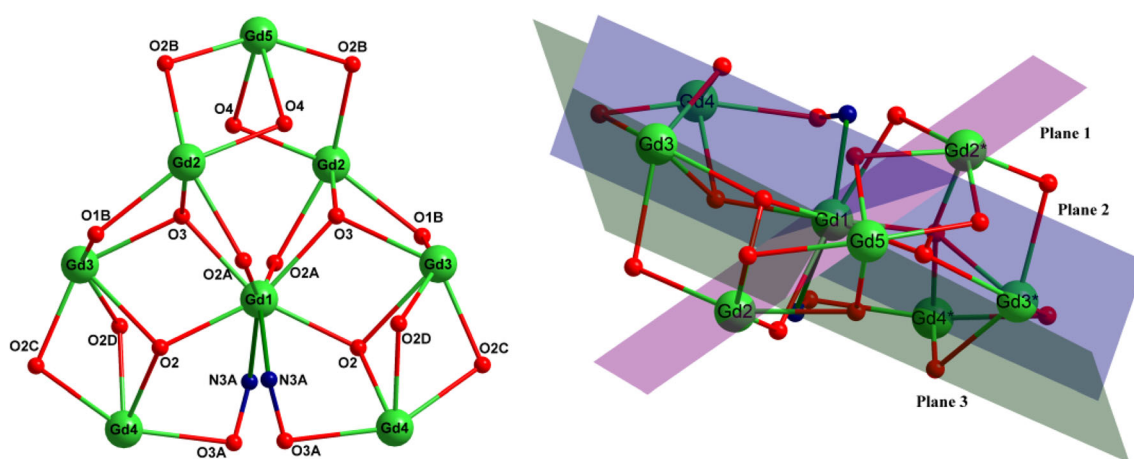


Figure 4. (left) The structure of the $\{\text{Gd}_3\}$ core motif and (right) the mean planes in the structure of complex 1^{Gd} .

Gd4^*), and triangular dodecahedron (Gd5) (Figure S3 and Tables S3 and S4, SI). Similar types of coordination environments were also observed in the complex 2^{Dy} . The packing diagram (Figure S5, SI) shows

that the shortest $\text{Gd}\cdots\text{Gd}$ separation between two adjacent $\{\text{Gd}_3\}$ molecules is 9.215 Å.

3.3 Magnetic properties

The temperature dependence of $\chi_{\text{M}}T$ for complexes 1^{Gd} and 2^{Dy} (χ_{M} is the molar magnetic susceptibility per Ln_8 unit) in the range 2–300 K were measured under a magnetic field of 0.1 T (Figure 5).

The $\chi_{\text{M}}T$ product for 1^{Gd} at room temperature ($62.60 \text{ cm}^3 \text{ K mol}^{-1}$) agrees well with that expected for eight isolated Gd^{III} ions with $g = 2$ and $S = 7/2$ ($63 \text{ cm}^3 \text{ K mol}^{-1}$ for $g = 2$). By lowering the temperature, the $\chi_{\text{M}}T$ product remains almost constant until ~ 100 K and then decreases sharply down to 2 K to reach a value of ($39.40 \text{ cm}^3 \text{ K mol}^{-1}$). The decrease of the χT product at low temperature is mainly due to an overall weak antiferromagnetic interaction between the Gd^{III} ions, combined with very small ZFS of the

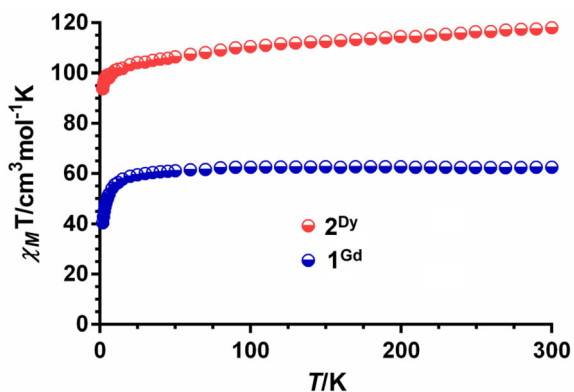


Figure 5. Temperature dependence of the $\chi_{\text{M}}T$ product for complexes 1^{Gd} and 2^{Dy} .

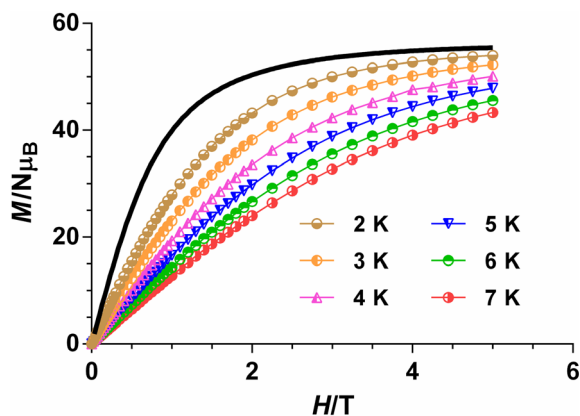


Figure 6. The field dependence of the magnetization plots for 1^{Gd} between 2 and 7 K. The black solid line corresponds to the Brillouin function for eight uncoupled Gd^{III} ions.

ground state (the Gd^{III} ion is an essentially isotropic ion) and Zeeman effects. The magnitude of the antiferromagnetic exchange interaction in 1^{Gd} could not be determined by diagonalization matrix methods because of the extremely high dimension of the matrices to be diagonalized for a Gd_8 system.

The existence of antiferromagnetic interactions between the Gd^{III} ions is supported by the field dependence of the magnetization at 2 K which is well below the Brillouin function for eight non-interacting Gd^{III} ions (Figure 6). At high field, the saturation of the magnetization is almost complete at 5 T ($54 \mu_{\text{B}}$), reaching values that agree well with the theoretical saturation values for eight Gd^{III} ions ($56 \mu_{\text{B}}$).

Experimental and theoretical studies^{70–72} carried out on oxygen-bridged Gd_2 complexes (alkoxido, phenoxido and carboxylate) have suggested that J becomes more antiferromagnetic as the Gd-O-Gd angle (θ), and consequently the $\text{Gd}\cdots\text{Gd}$ distance, decrease, and the Gd-O distances become more equal. The crossing point between antiferromagnetic and ferromagnetic interactions occurs approximately at θ and $\text{Gd}\cdots\text{Gd}$ values of 112° and 4.0 \AA , respectively. The global antiferromagnetic interaction observed in complex 1^{Gd} is not unexpected because, except for one angle (117.39°) and one distance (4.076 \AA), in the asymmetric unit all the θ angles (in the $3.677\text{--}3.887 \text{ \AA}$ range) and $\text{Gd}\cdots\text{Gd}$ distances (in the $96.70\text{--}109.02^\circ$) are below the above values.

It is worth mentioning that low temperature molecular magnetic coolers (MMCs) based on Gd^{III} complexes can exhibit higher magneto-caloric effects (MCEs), that is to say, an important change of magnetic entropy upon application of a magnetic field.⁷³ This effect is of great interest because Gd^{III} complexes with high MCE could be potentially employed for refrigeration applications *via*

adiabatic demagnetization.⁷³ We have studied the MCE properties of 1^{Gd} because of the following reasons. (i) The Gd^{III} ion exhibits negligible anisotropy due to the absence of orbital contribution and the largest single-ion spin ($S = 7/2$) arising from the $4f^7$ electron configuration. Both factors (small anisotropy and large spin ground state) favour an enhanced MCE. (ii) The weak antiferromagnetic interactions between the Gd^{III} ions give rise to multiple low-lying excited and field-accessible spin states, very close in energy with respect to each other, each of which can contribute to the magnetic entropy of the system.

In view of the above considerations, a significant magneto-caloric effect is expected for 1^{Gd} . The magnetic entropy changes ($-\Delta S_{\text{m}}$) that characterize the magneto-caloric properties of 1^{Gd} were calculated from the experimental isothermal field dependent magnetization data (Figure 7) by making use of the Maxwell relation:

$$\Delta S_{\text{M}} = (T, \Delta B) = \int_{B_i}^{B_f} \left[\frac{\partial M(T, B)}{\partial T} \right]_{\text{B}} dB$$

where B_i and B_f are the initial and final applied magnetic fields. The values of $-\Delta S_{\text{m}}$ for 1^{Gd} (Figure 7) under all magnetic fields increase with decreasing temperature from 7 to 3 K. The maximum value of $-\Delta S_{\text{m}}$ achieved for 1^{Gd} is $25.5 \text{ J kg}^{-1} \text{ K}^{-1}$ at $T = 3 \text{ K}$ and applied field change $\Delta B = 5 \text{ T}$ (Figure 6). Despite the antiferromagnetic interactions between the Gd^{III} ions in 1^{Gd} , there is an important change in $-\Delta S_{\text{m}}$, which is due to the easy spin polarization at a relatively low magnetic field. The extracted $-\Delta S_{\text{m}}$ value at $T = 3 \text{ K}$ is lower than that calculated for the full magnetic entropy content per mole $nR \ln(2S_{\text{Gd}} + 1) = 16.6 \text{ R} = 35.1 \text{ J kg}^{-1} \text{ K}^{-1}$. This is essentially due to the existence of

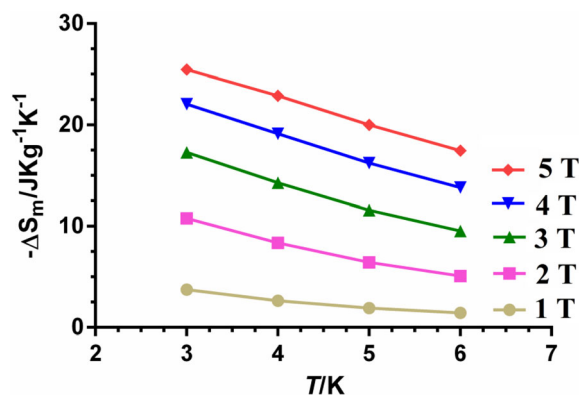


Figure 7. The magnetic entropy changes ($-\Delta S_{\text{m}}$) calculated using the magnetization data for 1^{Gd} from 1 to 5 T and temperatures from 3 to 7 K.

antiferromagnetic interactions between the Gd^{III} ions. It should be noted that the extracted $-\Delta S_m$ value at 5 T for **1**^{Gd} is close to that found for alkoxido/hydroxo bridged Gd₈ complexes with antiferromagnetic interaction between the Gd^{III} ions and possessing similar molecular weight (M_w),^{74–76} but lower than that found for other much more magnetically dense Gd^{III} complexes (M_w/N , where N = number of Gd³⁺ ions).^{77–83}

It is well known that Dy^{III} complexes are good candidates to exhibit slow relaxation of the magnetization because Dy^{III} is a Kramers ion with a bistable ground state (electronic structure composed of Kramers doublets, KDs) and a large magnetic moment. In addition to this, to observe slow magnetization relaxation, the ground state must be axial, to possess the largest value of M_J ($\pm 15/2$). This KD can be stabilized by an axial crystal field because the repulsive interactions between the ligands and its disc shape charge cloud are minimized.⁸⁴ It is worth noting that the axial ligand field can be attained by serendipity in low symmetry Dy^{III} complexes and, therefore, they often exhibit slow relaxation of the magnetization and SMM behaviour.

In view of the above considerations, we have performed ac magnetic susceptibility measurements under zero- and small-applied *dc* magnetic fields to know if **2**^{Dy} exhibits slow relaxation of the magnetization. The results of these measurements demonstrate that compound **2**^{Dy} does not exhibit any maximum above 2 K in the out-of-phase (χ''_M) signals even at the highest used frequency of 1400 Hz (Figure S6, SI). This behaviour could be due to a fast relaxation of the magnetization through quantum tunneling (QTM) and/or to the existence of a very small effective energy barrier that is not able to trap the magnetization in one of the two equivalent orientations at temperatures above 2 K. When the ac measurements were carried out in the presence of a small external *dc* field of 1000 Oe (Figure S6, SI), to fully or partly suppress the possible fast quantum tunneling relaxation, the temperature dependence of χ''_M for **2**^{Dy} at 1400 Hz did not significantly change. This either supports a very small height of anisotropic energy barrier (typically below 5 K) or suggests that the strong QTM process, leading to apparently lower U_{eff} values, is not effectively suppressed by the field and therefore should have its origin in hyperfine and intramolecular/intermolecular magnetic interactions.

4. Conclusions

In conclusion, we have demonstrated a *cluster expansion* synthetic method to assemble discrete multinuclear Ln^{III} complexes [Ln^{III} = Gd^{III}/Dy^{III}]. The

structural topology of these complexes is quite unprecedented among the known {Ln₈} topologies. The two complexes are also quite stable in the solution phase as confirmed by their ESI-MS studies. We have studied the magnetic properties of the Gd^{III} and Dy^{III} derivatives. Although the Dy^{III} derivative shows a slow magnetic relaxation at very low temperature and at a high frequency in the ac susceptibility measurements a clear maxima was not observed. The Gd^{III} derivative being isotropic and the Gd^{III} ions interacting in a weakly antiferromagnetic manner, the magneto-caloric effect in this complex was studied from the magnetization vs field data at different temperatures. This study reveals a maximum in the change of molar entropy ($-\Delta S_m$) of magnitude 25.5 J kg⁻¹ K⁻¹ at $T = 3$ K and an applied field change $\Delta B = 5$ T.

Supplementary Information (SI)

Scheme S1, Tables S1-S5, Figures S1-S6 are available at www.ias.ac.in/chemsci.

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