



# Mid-infrared quantum cascade laser spectroscopy probing of the kinetics of an atmospherically significant radical reaction, $\text{CH}_3\text{O}_2 + \text{NO}_2 + \text{M} \rightarrow \text{CH}_3\text{O}_2\text{NO}_2 + \text{M}$ , in the gas phase

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**Abstract.** The kinetic parameters of an important atmospheric reaction,  $\text{CH}_3\text{O}_2 + \text{NO}_2 + \text{M} \rightarrow \text{CH}_3\text{O}_2\text{NO}_2 + \text{M}$ , have been recorded by monitoring directly the changes in concentrations of methylperoxy radicals ( $\text{CH}_3\text{O}_2$ ) in the gas phase employing a new mid-infrared quantum cascade laser (QCL)-based apparatus.  $\text{CH}_3\text{O}_2$  radicals in our apparatus have been generated by pulsed UV laser (266 nm) photolysis of  $\text{CH}_3\text{I}$  in a gaseous mixture with oxygen. The absorption band corresponding to the mid-infrared O-O stretching fundamental of the peroxy radical, within a narrow spectral range, 1070–1120  $\text{cm}^{-1}$ , has been recorded by tuning the wavelength of the QCL operated in CW mode. The kinetics of the aforementioned reaction of  $\text{CH}_3\text{O}_2$  with  $\text{NO}_2$  has been followed by analyzing the changes of the infrared (QCL) decay profile of  $\text{CH}_3\text{O}_2$  at 9.1  $\mu\text{m}$  (1098.9  $\text{cm}^{-1}$ ) maintaining a pseudo first order reaction condition. We noticed that the rate constant of the reaction at 298 K varies in the range of  $(1.21\text{--}3.08) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  for changing the total pressure in the range of 75–730 mbar. The absorption cross-section of  $\text{CH}_3\text{O}_2$  at the probe wavelength (1098.9  $\text{cm}^{-1}$ ), has been estimated for the first time to be  $8.3 \pm 0.4 \times 10^{-20} \text{ cm}^2$ .

**Keywords.** Quantum cascade laser; time-resolved mid-infrared spectroscopy; transient absorption; peroxy radicals; absorption cross-section.

## 1. Introduction

Investigation of the kinetics of gas phase radical reactions is important in many interrelated research fields like atmospheric chemistry, combustion chemistry, chemistry of interstellar medium to name but a few.<sup>1–6</sup> Detection of gas phase organic radicals is tricky because of their transient nature with a typical lifetime of a few microseconds to several minutes depending on the environment and nature of the radical. Therefore, time-resolved spectroscopic techniques are required for probing the gas phase transient radicals. Several time-resolved probing methods were developed over the last few decades and the most popular methods available in the literature are based on laser induced fluorescence (LIF) spectroscopy, UV/Visible/NIR absorption spectroscopy, mass spectrometry, etc.<sup>7–16</sup> Each of these probing techniques has its own merits as well as limitations. For example, laser induced fluorescence (LIF) is

known to be an extremely sensitive method for detection of fluorescent radicals, but for non-fluorescent species, viz. organic peroxy radicals, the method cannot be used directly.<sup>7,8,13,14</sup> Likewise, in spite of being extremely sensitive, simple mass spectrometric technique cannot be used in differentiating isomeric species present in a gas mixture. Organic radicals usually have strong absorption bands in ultraviolet, and for that reason UV spectroscopy is being used extensively over a long period for probing different radical species.<sup>9,17–19</sup> However, an inherent problem of the method is that the UV absorption bands are usually broad, and in consequence it can have only limited use when more than one radical species are present in a gas mixture, which display overlapping absorption bands.<sup>9,17</sup>

In recent years, near-infrared (NIR) and visible spectroscopy, coupled with cavity enhanced methods, e.g., cavity ring down spectroscopy (CRDS) have emerged as powerful tools for probing the kinetics of organic radical reactions.<sup>16,20,21</sup> Although NIR spectroscopy provides better selectivity in comparison with the methods based

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on UV absorption, the small NIR absorption cross-sections of organic species is a serious disadvantage in kinetics measurements at low concentrations of the radical species. On the other hand, the characteristic vibrational bands of organic species, particularly in the mid-IR fingerprint region are very sharp with typically more than an order of magnitude absorption cross-section values, and the concerned species in a mixture can be selectively probed for measurement of reaction kinetics.<sup>22–24</sup> Therefore, the band selective mid-infrared spectroscopic method should be superior in studies of kinetics of the radical species produced by laser photolysis method, because typically more than one radical species are produced in more than one dissociation pathways. However, in the literature, the reports on mid-infrared (MIR) probing of organic radicals for kinetics measurements of fast radical reactions are rare, and this is mainly due to non-availability of stable commercial lasers in this optical region.<sup>22,23</sup> Recently, Quantum Cascade Lasers (QCLs) have been introduced as stable and intense laser sources in the MIR spectral region. We describe here a QCL-based methodology in reaction kinetics studies of gaseous radical species. The reaction used to demonstrate the potential of the method is the reaction (Eq. 1), where an adduct between methylperoxy radical ( $\text{CH}_3\text{O}_2$ ) and nitrogen dioxide ( $\text{NO}_2$ ) is produced *via* three body collision involving a background gas molecule ( $\text{M} = \text{N}_2, \text{Ar}, \text{etc.}$ ).



Methylperoxy radical is an important intermediate in the oxidation of volatile organic compounds (VOCs) in the atmosphere and also in combustion systems.<sup>25–27</sup> From the perspective of atmospheric chemistry, it is an important reaction intermediate, especially with respect to the reactions that occur in polluted environments, where  $\text{NO}_2$  concentrations are high. Recently, this reaction (Eq. 1) has gained much attention, because, similar to much discussed peroxyacetylnitrate (PAN), peroxyethylnitrate ( $\text{CH}_3\text{O}_2\text{NO}_2$ ) can also have a long lifetime, especially at upper troposphere, and thus the species can act as a sink as well as a transporting agent of  $\text{NO}_2$ .<sup>28,29</sup> It is important therefore to estimate the accurate kinetic data of the reaction of  $\text{NO}_2$  with  $\text{CH}_3\text{O}_2$  and other  $\text{RO}_2$  species, because, atmospheric  $\text{NO}_2$  concentration has a large effect on the tropospheric abundances of  $\text{O}_3$  and OH radical. The lifetime of  $\text{CH}_3\text{O}_2\text{NO}_2$  is highly sensitive to temperature, and it is less than a second at room temperature near sea level, but it could be as long as several days in the upper troposphere at low temperatures.<sup>28,29</sup> Recently, Nault *et al.*, measured the concentration of  $\text{CH}_3\text{O}_2\text{NO}_2$  in the upper troposphere

by an aircraft campaign ( $\sim 100$  parts per trillion by volume (pptv) at higher altitudes and  $\sim 35$  pptv at lower altitudes).<sup>29</sup> Thus, because of the enormous atmospheric significances, it is essential to have reliable experimental data concerning accurate values of the rate constants of this reaction in different conditions, and also to verify the rate constant data employing different probing methodologies. Although kinetics of this reaction has been studied before by several methods, we report here the measurement of the same by direct IR absorption monitoring of  $\text{CH}_3\text{O}_2$  radical for the first time.<sup>30–37</sup> Recently, Lee and co-workers measured vibrational spectrum of various peroxy radicals in the mid-infrared by Step-Scan FTIR method at low spectral resolutions, but no kinetic data of those radicals were reported using this method.<sup>38–40</sup> We also have measured for the first time the infrared absorption cross-section of the peroxy radical near the band maximum of its O-O stretching fundamental.

## 2. Experimental

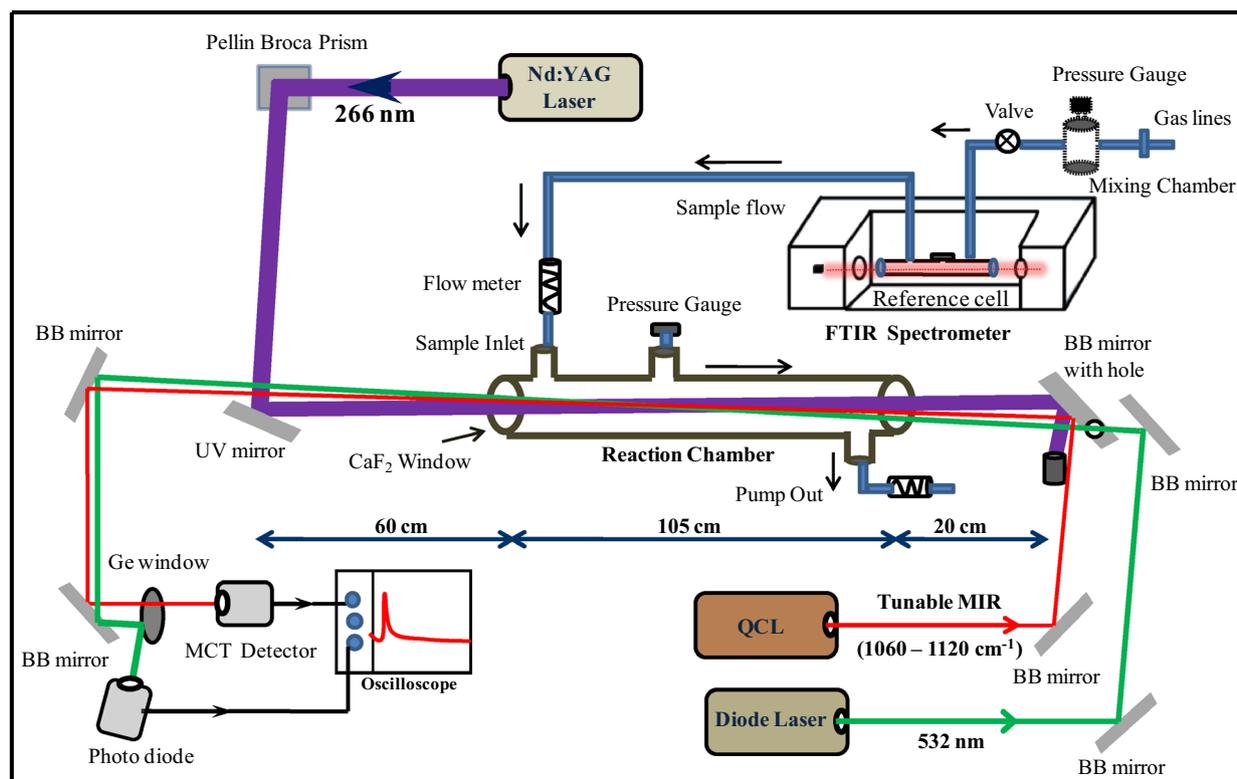
### 2.1 Chemical reactor and gas handling system

A schematic of the experimental setup is depicted in Figure 1.

The reactor is a cylindrical stainless steel tube of length 105 cm and inner diameter 2.5 cm. Two ends of the tube are closed with two  $\text{CaF}_2$  windows and sealed by silicon gaskets.  $\text{CaF}_2$  has high transmittance with respect to UV (266 nm), mid-IR (9–10  $\mu\text{m}$ ) and visible (532 nm) wavelengths, and thus the lasers are introduced into the cell through the windows. The gas inlet and outlet ports are placed at the two opposite ends of the reactor. The distance between these two ports along the length of the cell is  $\sim 100$  cm. A separate mixing chamber of volume  $\sim 4$  L has been used to prepare the reaction gas mixture. The gas mixture ( $\text{CH}_3\text{I}$  vapour diluted in  $\text{N}_2/\text{O}_2$  mixture) is introduced into the reactor as a continuous flow through the sample inlet, and drained out through the outlet port by a root pump (Adixen ACP 15). The flow of the gas mixture is controlled by a precision metering valve (Swagelok) and the flow rate is monitored both at the inlet and outlet by two mechanical flow meters (Rota Yokogawa), and the mean of the two flow meters reading is typically  $\sim 5$ – $25$  L/min. A high pressure gauge (Pfeiffer CMR 361) is attached with a vacuum port located at the middle of the reactor. The temperature within the flow reactor remains almost the same as that of the room temperature (298 K), which has been measured using a thermocouple temperature sensor attached at one of the terminal ports.

### 2.2 Optical system

**2.2a Photolysis system:** For generation of radicals, the precursor compound ( $\text{CH}_3\text{I}$  in the present case), is photolyzed



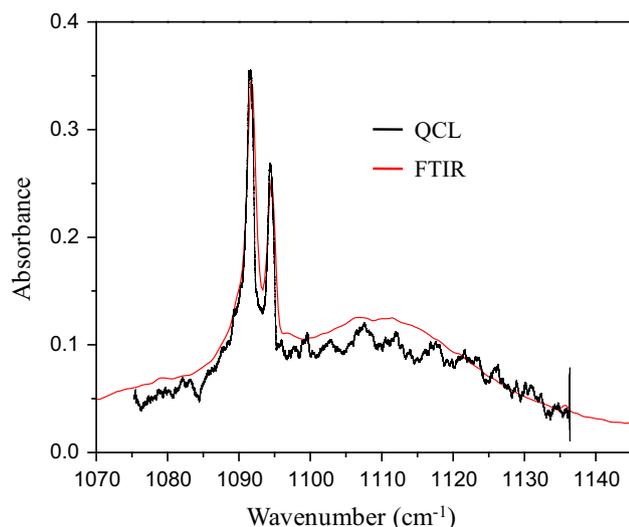
**Figure 1.** Schematic diagram of the pulsed laser photolysis time-resolved IR absorption kinetics measurement apparatus. Overall length of the reactor is 105 cm, and the distance between inlet and outlet ports of reaction gas mixture is 100 cm. Distance between UV mirror and input window of the reactor is 60 cm, while the same in the other side is 20 cm (distances mentioned in the figure are not scaled). The BB (broadband) mirror is coated with protected aluminum (450 nm–20  $\mu$ m). UV and IR beams are aligned at  $\sim 2^\circ$  with respect to each other that results in 100 cm overlap between them within the reactor.

by the 4<sup>th</sup> harmonic UV light pulses (266 nm) of a Q-switched Nd:YAG laser (Spectra Physics Quanta-Ray Lab Series 150). The photolysis laser beam is introduced into the reactor through center of one of the terminal windows, and the beam passes through the reactor along its central axis. The beam shape of the photolysis laser is nearly circular with a diameter of 8 mm. The laser is operated at 10 Hz repetition rate, and typical pulse energy is  $\sim 10$  mJ.

**2.2b Probe laser:** The infrared beam from a QCL laser (Pranalytica Inc., Model: Omnilux), which was operated in continuous wave (CW) mode, is introduced into the reactor in counter propagating direction with respect to the photolysis laser keeping a small angle between the two beams so that they can have overlap within the reactor to the largest possible extent. The beam diameter of the IR laser is  $\sim 2$  mm while the same for the UV laser is  $\sim 8$  mm and this is particularly helpful in keeping the two lasers overlapped over a longer distance. The angle between the two laser beams is  $\sim 2^\circ$ , and for achieving this situation the input windows of the cell are kept far away from the folding mirrors that steer the beams into the reactor. The distance between the UV quartz mirror and starting point of the overlap between the two lasers is approximately 63 cm. This large distance results

in the effective gap of  $\sim 1$  cm between the two laser beams near UV mirror. The QCL has a visible beam guide of HeNe laser, which is extremely helpful to monitor the alignment for beam overlap. At the exit end, the distance between the output window and the broad band (BB) IR mirror is nearly 20 cm. We have estimated that the overlap length between the two lasers within the cell is  $\sim 100$  cm, and this has been considered as the effective reaction path length for absorption of probe wavelength by the peroxy radicals. The range of frequency tunability of the QCL laser for this mode of operation is  $\sim 60$   $\text{cm}^{-1}$  (wavelength maximum, 9.1  $\mu\text{m}$ ) and has a linewidth of  $\sim 0.003$   $\text{cm}^{-1}$  with a power of  $\sim 100$  mW at the peak of the laser emission profile. The wavelength calibration of the QCL is verified by measuring the infrared fundamental band profile of acetone vapor (at room temperature) near 1091  $\text{cm}^{-1}$  and a comparison of the recorded spectrum with the one recorded using a FTIR spectrometer is shown in Figure 2.

**2.2c Reactant concentration monitoring by optical methods:** Rate constants of bimolecular reactions are generally measured adopting pseudo first order conditions where one of the reactants is kept in excess compared to the other, and the same approach is adopted here for measuring the



**Figure 2.** A comparison of the spectral features of the infrared  $1091\text{ cm}^{-1}$  band of acetone vapor at room temperature recorded in transmission mode by continuous scanning of the QCL, with the spectrum obtained using a FTIR spectrometer (Model: IFS66S, Bruker Optics). For the latter, the spectrometer resolution is maintained at  $0.25\text{ cm}^{-1}$ . The depicted comparison ensured the wavelength calibration of the QCL used.

rate constant of reaction in Eq. 1. For accurate measurement of the rate constant under this condition, precise measurements of the concentrations of the reactants within the reactor, particularly for  $\text{NO}_2$  which is taken in excess is important and the following arrangement is made for this purpose. The reaction gas mixture is first flown through a reference cell placed directly inside the sample compartment of a FTIR spectrometer (Bruker IFS 66 s), and this arrangement enables monitoring absolute concentrations of the different reactants used. In addition, an *in situ* arrangement for  $\text{NO}_2$  concentration monitoring has been made by introducing a CW beam of wavelength  $532\text{ nm}$  from a diode laser which is aligned with the photolysis laser (Figure 1). The attenuation of  $532\text{ nm}$  laser intensity was recorded continuously for monitoring the actual concentration of  $\text{NO}_2$  inside the main reactor. Absorption cross-section of  $\text{NO}_2$  at  $532\text{ nm}$  is  $\sigma = 1.50 \times 10^{-19}\text{ cm}^2\text{ molecule}^{-1}$ .<sup>41</sup> Absolute concentrations of  $\text{CH}_3\text{O}_2$  radicals are evaluated from time-resolved mid-IR absorption data. The required absorption cross-section value of the radical at  $9.1\text{ }\mu\text{m}$  is also measured for the first time in the present study.

### 2.3 Signal detection and data recording

The intensity of the probe IR beam after exiting the reaction chamber is monitored by a fast pre-amplified MCT detector (VIGO PVMI-4TE-10.6). The time-resolved absorption signal from the detector is viewed, averaged and stored in a digital storage oscilloscope (Lecroy wavesurfer 62Xs, 600 MHz, 2.5 GS/s). A  $50\text{ }\Omega$  termination was used at the input of the

oscilloscope. Typically, the decay signals were averaged over 50–100 laser pulses before recording. For further improvement of the signal to noise ratio (S/N), the data are treated with Savitzky-Golay smoothing algorithm.<sup>42</sup> This same smoothing algorithm was used in a previous study for measuring the kinetics of reaction (Eq. 1).<sup>31</sup> The output of  $532\text{ nm}$  diode laser is detected by a photodiode and the signal is sent to the same oscilloscope for records. Since both the beams propagate through the cell along almost the same path, the latter is separated from the former by placing a Ge window in front of the MCT detector, which allows the IR beam to pass through but the wavelength  $532\text{ nm}$  is reflected.

### 2.4 Materials

Methyl iodide ( $\text{CH}_3\text{I}$ ) was procured from Sigma Aldrich and used without further purification. The nitrogen and oxygen gases used were of the purity 99.99%.  $\text{NO}_2$  was prepared by reaction of  $\text{NaNO}_2$  with  $\text{FeSO}_4 \cdot 7\text{H}_2\text{O}$  in presence of excess oxygen, and the gas was dried by using a column of  $\text{Na}_2\text{CO}_3$  before introducing into the mixing chamber.

## 3. Results and Discussion

### 3.1 Measurement of the kinetics of



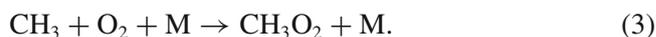
**3.1a Production of  $\text{CH}_3\text{O}_2$  radical:** In the presence of oxygen, methyl radical ( $\text{CH}_3$ ) produced *via* photolysis of methyl iodide ( $\text{CH}_3\text{I}$ ) by  $266\text{ nm}$  UV laser pulses of a Q-switched Nd:YAG laser (4<sup>th</sup> harmonic wavelength), reacts to generate  $\text{CH}_3\text{O}_2$  radical. At this wavelength, the absorption cross-section of  $\text{CH}_3\text{I}$  is very large ( $\sigma \sim 1 \times 10^{-18}\text{ cm}^2\text{ molecule}^{-1}$ ), and photodissociation quantum yield ( $\Phi$ ) for production of  $\text{CH}_3$  radical is effectively unity.<sup>41,43,44</sup> The concentration of the  $\text{CH}_3$  radicals produced is determined in the following way:

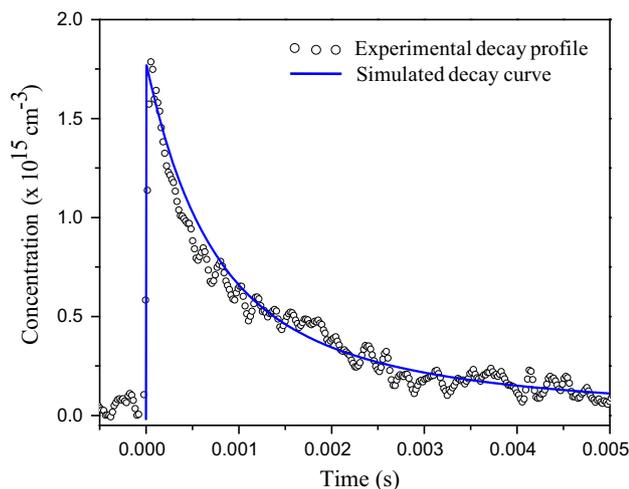
Methyl radical concentration,  $[\text{CH}_3]$

$$= [\text{CH}_3\text{I}]_x \sigma \times \Phi \times F \quad (2)$$

where,  $[\text{CH}_3\text{I}]$  is the concentration of the precursor  $\text{CH}_3\text{I}$ ,  $\sigma$  the absorption cross-section of  $\text{CH}_3\text{I}$  at  $266\text{ nm}$ ,  $\Phi$  the quantum yield of photo-dissociation of  $\text{CH}_3\text{I}$  at  $266\text{ nm}$ , and  $F$  the laser fluence.

In the present experiment for kinetic measurements, concentration of  $\text{CH}_3\text{I}$  being  $\sim 1 \times 10^{16} - 5 \times 10^{16}\text{ molecule cm}^{-3}$  (determined from FTIR absorption data) and laser fluence  $\sim 2 \times 10^{16}\text{ photons cm}^{-2}$ , the estimated concentration of the produced  $\text{CH}_3$  radical is typically  $(2-10) \times 10^{14}\text{ cm}^{-3}$ . Adduct formation of  $\text{CH}_3$  radical with  $\text{O}_2$  is a termolecular reaction (Eq. 3) and the following are the values of the rate constants:

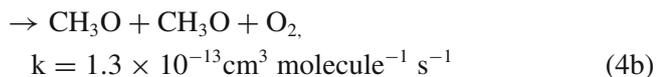
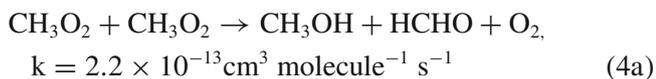




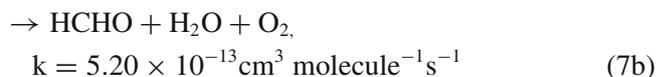
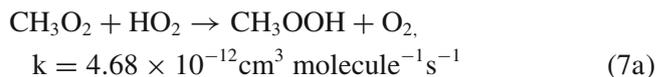
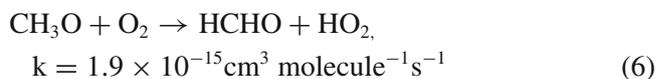
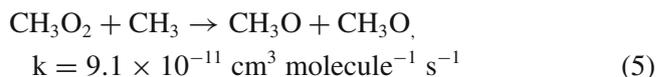
**Figure 3.** A typical decay signal of  $\text{CH}_3\text{O}_2$  radicals (in the absence of  $\text{NO}_2$ ) for probing at  $9.1 \mu\text{m}$ .  $[\text{CH}_3\text{I}] = 4.9 \times 10^{16} \text{ cm}^{-3}$ ,  $[\text{O}_2] = \sim 9.0 \times 10^{17} \text{ cm}^{-3}$  and total pressure of the reaction cell was 520 mbar. The decay curve obtained from modeling (blue trace) is also presented for a comparison with the experimental signal.

Here,  $k_0 = 7.0 \times 10^{-31} [\text{N}_2] \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  and  $k_\infty = 1.8 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ .  $k_0$  and  $k_\infty$  are low and high pressure rate constants, respectively, at 298 K.<sup>43,44</sup> Almost all of  $\text{CH}_3$  radicals produced by such means is converted into the peroxy radicals,  $\text{CH}_3\text{O}_2$ , within a few microseconds. Thus, at 520 mbar total pressure and in the absence of  $\text{NO}_2$ , it has been estimated from a reaction modeling study (described later) that the number of  $\text{CH}_3$  radicals produced on average from each laser shot,  $5.0 \times 10^{14} \text{ cm}^{-3}$ , is converted into  $\text{CH}_3\text{O}_2$  with a number density of  $4.7 \times 10^{14} \text{ cm}^{-3}$  within  $10 \mu\text{s}$ . Depicted in Figure 3 is a typical decaying signal of  $\text{CH}_3\text{O}_2$  radicals (in absence of  $\text{NO}_2$ ) probed at  $9.1 \mu\text{m}$  for a total pressure of 520 mbar. The detection limit of the methylperoxy radical in the current experimental setup is  $\sim 2 \times 10^{14} \text{ cm}^{-3}$  ( $S/N = 2$ ). It should be noted that in Figure 3, the concentration of  $\text{CH}_3\text{O}_2$  radical produced is much higher. However, for kinetic experiments, lower  $\text{CH}_3\text{O}_2$  concentrations have been used.

As reported in previous studies,  $\text{CH}_3\text{O}_2$  radicals could decay predominantly *via* the following self reaction pathways:<sup>16,20,43–45</sup>



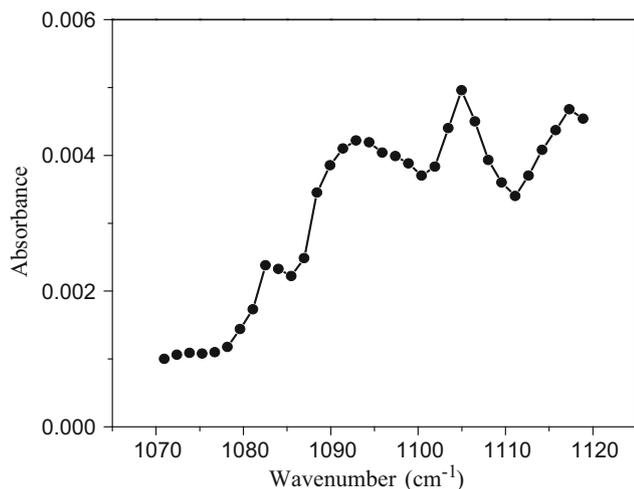
The following reactions also contribute to the decay of  $\text{CH}_3\text{O}_2$  radicals.<sup>43,44,46</sup>



The production and decay of  $\text{CH}_3\text{O}_2$  radicals as obtained from simulation considering reactions 3–7 is also presented in Figure 3 (blue trace), and this has been done using a standard computer program called ACUCHEM.<sup>47</sup> The experimental  $\text{CH}_3\text{O}_2$  concentration profile shows satisfactory agreement with the modeled one. The fate of I atoms formed in the photolysis of  $\text{CH}_3\text{I}$  is as follows. I atoms rapidly combine with  $\text{CH}_3\text{O}_2$  to form  $\text{CH}_3\text{O}_2\text{I}$ , which react very rapidly with I atoms to give rise to  $\text{CH}_3\text{O}_2$  and  $\text{I}_2$ . As a result, there is practically no loss of  $\text{CH}_3\text{O}_2$  radicals in this process.<sup>20,48,49</sup> It is worth mentioning here that the concentrations of the reactants used in our experiments are not too small, and when the reaction with these concentrations is performed in a static cell, we observe accumulation of molecular iodine ( $\text{I}_2$ ) on the walls of the reactor. However, in the flowing reactor, with the flow rates that are typically used, we never noticed  $\text{I}_2$  accumulation.

### 3.2 Mid-IR absorption spectrum and absorption cross-section of $\text{CH}_3\text{O}_2$ radical

The mid-infrared spectrum of  $\text{CH}_3\text{O}_2$  radical has been measured previously by Huang *et al.*, using a step-scan FTIR method at a low spectral resolution of  $4 \text{ cm}^{-1}$ .<sup>38</sup> Using the QCL we have recorded the O-O stretching absorption profile of the radical within  $1070\text{--}1120 \text{ cm}^{-1}$  range with a laser resolution of 100 MHz by the following approach. Time-resolved absorption signals were recorded at discrete wavelengths (maintaining a step size of  $1.5 \text{ cm}^{-1}$ ) operating the laser in CW mode. Absorbance value at each wavelength is measured using the standard relation, Absorbance (O.D.) =  $\ln(I_0/I)$ , where  $I_0$  is the signal level before introducing photolysis laser and  $I$ , the signal after  $\sim 10 \mu\text{s}$  of firing of the photolysis laser. The spectrum, as shown in Figure 4, has been developed by line joining the discrete absorbance values as a function of laser frequency, and within the depicted spectral range it looks nearly similar to the one reported by Huang *et al.*<sup>38</sup> The segment of the band profile within  $1070\text{--}1120 \text{ cm}^{-1}$  range was assigned to a vibrational



**Figure 4.** The infrared absorption spectrum of  $\text{CH}_3\text{O}_2$  radical in the depicted spectral range permitted by the tunability of QCL. The detailed method used to record the spectrum has been discussed in the text.

fundamental corresponding to O-O stretch with a contribution from  $\text{CH}_3$  rocking, and this band was shown to have a maximum at  $1117\text{ cm}^{-1}$ . The same peak is also present in the spectrum shown in Figure 4, although the power of QCL is low at this wavelength. There is another peak at  $\sim 1094\text{ cm}^{-1}$ , and although no assignment for it was discussed in reference 38, on the basis of theoretical prediction it can be assigned to  $\text{CH}_3$  wagging fundamental of the radical<sup>38</sup>. In measuring the decay kinetics, the radical has been probed at the maximum of the laser emission wavelength,  $9.1\text{ }\mu\text{m}$  ( $1098.9\text{ cm}^{-1}$ ), and it is seen that the radical has a considerably large absorbance at this wavelength, which results in a good signal-to-noise ratio.

The absorption cross-section ( $\sigma$ ) at the probe wavelength has been evaluated in the following manner. The cross-section ( $\sigma$ ) is defined as,  $\sigma = \text{Absorbance}/(n l)$ ,  $n$  denotes the concentration of  $\text{CH}_3\text{O}_2$  radical and  $l$ , the optical path length. Here, absorbance is determined from the time-resolved absorption signal in the way mentioned before. Optical path length  $l$  is 100 cm. Precise measurement of  $n$ , i.e., the concentration of  $\text{CH}_3\text{O}_2$  radical is important for accurate measurement of the cross-section ( $\sigma$ ). A coarse measurement of the concentration of  $\text{CH}_3\text{O}_2$  radical, i.e.,  $n$ , is evaluated from laser fluence and  $\text{CH}_3\text{I}$  concentration following the procedure as mentioned previously. For fine measurement of the initial  $\text{CH}_3\text{O}_2$  concentration, following approach has been taken. Recently, Farago *et al.*, measured the absorption cross-section of  $\text{CH}_3\text{O}_2$  radical at NIR region using a reaction model to accurately quantify the initial radical concentration and its time evolution.<sup>20</sup> The same approach has also been adopted previously by

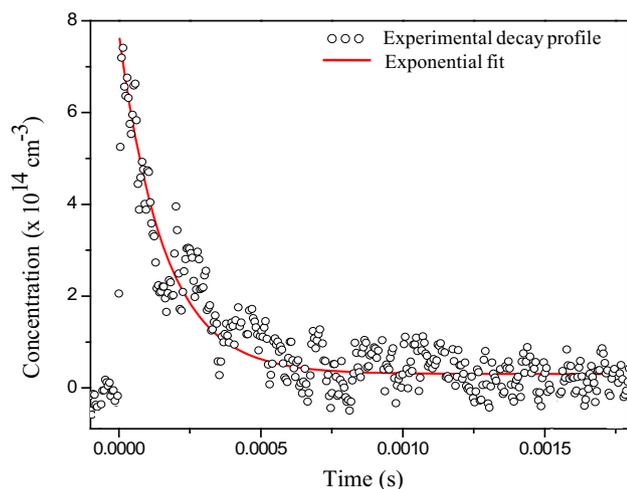
Pushkarsky *et al.*, and Atkinson *et al.*<sup>16,45</sup> In the present study, as shown in Figure 3, a simulation of the decay profile (considering reactions, Eqs. 3–7) of  $\text{CH}_3\text{O}_2$  radicals is carried out. During the construction of the decay profile, the initial  $\text{CH}_3\text{O}_2$  concentration is iterated so as to match the simulated decay profile well with the experimental plot as depicted in Figure 3. This initial  $\text{CH}_3\text{O}_2$  concentration ( $1.8 \times 10^{15}\text{ molecule cm}^{-3}$ ) is then used to estimate the absorption cross-section. Using this concentration data, the absorption cross-section of  $\text{CH}_3\text{O}_2$  radical at  $9.1\text{ }\mu\text{m}$  ( $1098.9\text{ cm}^{-1}$ ) has been estimated to be  $8.3 \times 10^{-20}\text{ cm}^2$ . Another source of uncertainty to the value of  $\sigma$  may come from the consideration of the effective optical path length as 100 cm which may typically vary within 5%. Considering this uncertainty, the value of absorption cross-section  $\sigma$  is taken as  $8.3 \pm 0.4 \times 10^{-20}\text{ cm}^2$ .

### 3.3 Kinetics of $\text{CH}_3\text{O}_2 + \text{NO}_2$ reaction

In presence of  $\text{NO}_2$ ,  $\text{CH}_3\text{O}_2$  reacts with  $\text{NO}_2$  (Eq. 1) apart from the self reactions. The kinetics experiments were performed under pseudo-first order condition where  $[\text{NO}_2] > [\text{CH}_3\text{O}_2]$ .

Before presenting our results, first we review here the studies performed earlier. In the decade of 1980 s, four different research groups carried out measurements independently by probing  $\text{CH}_3\text{O}_2$  radical at an ultraviolet wavelength near 250 nm.<sup>30–33</sup> Using the same method, Bridier *et al.*, studied the reaction in the temperature range of 333–373 K at 730 Torr pressure, and Wallington *et al.*, measured the rate constant values in the pressure range of 0.5–14 atm at 295 K.<sup>34,35</sup> As pointed out before, although the UV absorption cross-section is higher, but UV absorption band is very broad and that limits selectivity. Later, Bacak *et al.*, used chemical ionization mass spectrometry as a probe to study the reaction at low temperatures.<sup>36</sup> Very recently, McKee *et al.*, studied the kinetics of the reaction as well as equilibrium parameters for the process  $\text{CH}_3\text{O}_2 + \text{NO}_2 \leftrightarrow \text{CH}_3\text{O}_2\text{NO}_2$  by probing  $\text{CH}_3\text{O}_2$  radicals indirectly.<sup>37</sup>  $\text{CH}_3\text{O}_2$  was photolyzed to produce  $\text{CH}_3\text{O}$  and  $\text{OH}$  radicals, and these secondary radicals were probed by laser induced fluorescence spectroscopic method.

Figure 5 represents a typical decay profile for IR absorption of  $\text{CH}_3\text{O}_2$  radical at the probing wavelength ( $9.1\text{ }\mu\text{m}$ ) in presence of  $\text{NO}_2$  under the condition of a fast flow of the reaction gas mixture within the reactor. An important merit of the method, in spite of smaller infrared absorption cross-section of the radical at this wavelength in comparison with the previously used UV probing wavelength, is that in the absence of photolysis laser the background absorption of the gas mixture is

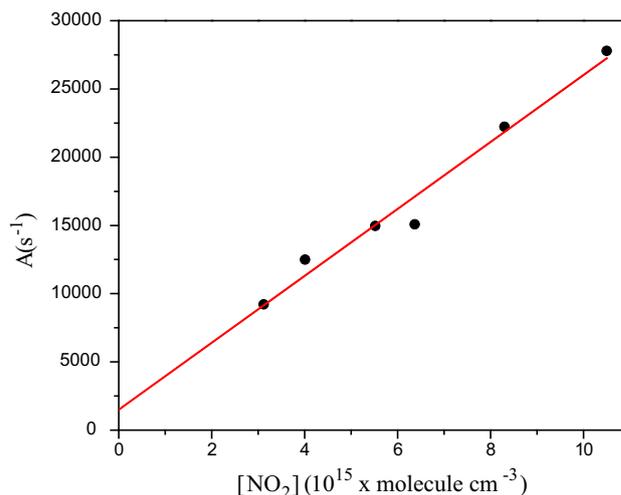


**Figure 5.** Decaying IR absorption signal of  $\text{CH}_3\text{O}_2$  radicals in presence of  $\text{NO}_2$  at 75 mbar total pressure. The initial concentration of  $\text{CH}_3\text{O}_2$  radical and  $\text{NO}_2$  are  $7.30 \times 10^{14} \text{ cm}^{-3}$  and  $5.30 \times 10^{15} \text{ cm}^{-3}$ , respectively. It is seen that the decay can be nicely fitted with a single exponential decay curve.

zero and the adduct formed  $\text{CH}_3\text{O}_2\text{NO}_2$  also does not absorb at this wavelength.<sup>50</sup>

It is notable that for maintaining pseudo first order reaction conditions the concentration ratio,  $[\text{NO}_2]:[\text{CH}_3\text{O}_2]$ , has been varied within the range of  $\sim 8\text{--}40$ . Although this ratio is relatively small, but the single exponential behavior of the plot in Figure 5 indicates that the pseudo first order condition is indeed achieved around this value of the concentration ratio. At still higher concentrations of  $\text{NO}_2$  we have noticed occurrence of dimer ( $\text{N}_2\text{O}_4$ ) formation ( $\text{N}_2\text{O}_4$  was probed at  $750 \text{ cm}^{-1}$  by FTIR). Since the reaction rate of  $\text{N}_2\text{O}_4$  with  $\text{CH}_3\text{O}_2$  is not known, to avoid possible source of errors for such dimer formation,  $\text{NO}_2$  concentration was kept below  $\sim 1 \times 10^{16} \text{ cm}^{-3}$ , and no prominent sign of dimer formation was observed at this concentration of  $\text{NO}_2$ . We would like to point out here that in a number of previous studies, pseudo first order condition was invoked upon use of the same concentration ratio between  $\text{NO}_2$  and  $\text{CH}_3\text{O}_2$ .<sup>30,31</sup> In addition, the linear plot of variation of the decay constant with  $[\text{NO}_2]$  in the reactor as depicted in Figure 6 (discussed below) is a direct support in favor of pseudo 1st order behavior of the reaction. The rate constant of the reaction (Eq. 1) has been measured in the following way.

The associated time constant ( $\tau$ ) for decay of the individual  $\text{CH}_3\text{O}_2$  absorption trace is measured by fitting the decay curve with an exponential function as stated earlier. The decay rate constant ( $A$ ) is the reciprocal of the corresponding time constant ( $A = 1/\tau$ ). The reaction rate constant ( $k$ ) is calculated from decay constant  $A$  using the following relation:



**Figure 6.** Plot depicting the changes of the estimated values of the decay constant ( $A$ ) for different concentrations of  $\text{NO}_2$  (total pressure 305 mbar).

$$A = k_1 + A_d = k[\text{NO}_2] + A_d \quad (8)$$

where,  $A$  = decay constant,  $k_1$  = pseudo first order rate constant,  $k$  = bimolecular rate constant,  $[\text{NO}_2]$  = concentration of  $\text{NO}_2$  ( $\text{molecule cm}^{-3}$ ),  $A_d$  = decay constant of  $\text{CH}_3\text{O}_2$  in the absence of  $\text{NO}_2$ . The slope of the plot of  $A$  vs  $[\text{NO}_2]$  gives the bimolecular rate constant for the reaction. A representative plot of the decay constant ( $A$ ) vs concentration of  $\text{NO}_2$  is presented in Figure 6. The experiment has been performed in  $\text{N}_2$  buffer gas with a total pressure of 305 mbar. The slope of the linear plot, i.e. the rate constant of the reaction is  $k = 2.45 \pm 0.47 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  (error limit:  $2\sigma$ ).

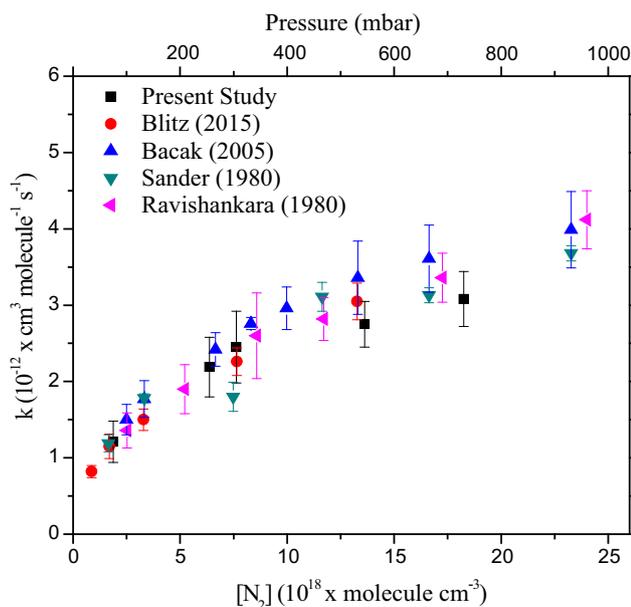
It is worth mentioning about a possible error source concerning depletion in concentration of  $\text{NO}_2$  as a result of UV photolysis. The UV absorption cross-section of  $\text{NO}_2$  at 266 nm is typically low,  $\sim 2.2 \times 10^{-20} \text{ cm}^2$ , and although the quantum yield for photodissociation of  $\text{NO}_2$  at this wavelength is not known, but for the sake of discussion we have considered it to be 1. For  $\text{NO}_2$  concentration of  $5.30 \times 10^{15} \text{ cm}^{-3}$  (Figure 5), the extent of its loss for each laser pulse has been calculated (following Eq. 2) and the value is  $4.10 \times 10^{12} \text{ cm}^{-3}$ , which is quite low. Moreover, the actual concentration of  $\text{NO}_2$  inside the reactor during the laser photolysis is probed by a 532 nm CW diode laser, and the change of its signal in the presence and absence of the photolysis laser is also not significant.

### 3.4 Effect of total pressure on rate constant

The reaction was studied at different total pressures within the range of 75–730 mbar at 298 K by varying

**Table 1.** Bimolecular rate constant values estimated as a function of total pressure using N<sub>2</sub> as the buffer gas (error is 2  $\sigma$ ).

Pressure (mbar)	Bimolecular rate constant $k$ ( $10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ )
730	$3.08 \pm 0.36$
545	$2.75 \pm 0.30$
305	$2.45 \pm 0.47$
255	$2.19 \pm 0.39$
75	$1.21 \pm 0.27$

**Figure 7.** Plot depicting the variation of the rate constant for reaction between NO<sub>2</sub> and CH<sub>3</sub>O<sub>2</sub> due to change in N<sub>2</sub> buffer gas pressure. The observed results [■] are compared with similar studies performed earlier.

the partial pressure of N<sub>2</sub> buffer gas, and the rate constant values estimated are shown in Figure 7 as well as in Table 1. As expected for an adduct formation reaction, the rate constant decreases steadily with decrease in total pressure.

In Figure 7, a comparison of the changes of the estimated values of the rate constants with total pressure in the present study with those reported by other workers by use of different methods is shown. It is seen that the agreements are very much satisfactory in lower pressure range. At higher pressures, although the rate constant values estimated here by direct infrared probing of CH<sub>3</sub>O<sub>2</sub> radical is a little smaller compared to the UV absorption probing, the values fall within the error range.

## 4. Conclusions

We have demonstrated here that a quantum cascade laser-based time-resolved mid-infrared absorption spectroscopic method can be used as an effective means for measuring the kinetic parameters of radical reactions in the gas phase. The performance of the apparatus has been demonstrated measuring the reaction kinetics of a tropospherically important reaction between CH<sub>3</sub>O<sub>2</sub> radical and NO<sub>2</sub>. Our estimated rate constant of the reaction for the range of total pressure between 75 and 730 mbar varies between  $(1.21-3.08) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  which agree quite satisfactorily with the previous reports. The major merit of the method, in comparison to much used UV probing of the radicals, is selectivity with respect to specific reactant or product of the reaction, and this is possible due to narrower spectral widths of the infrared bands in comparison with typically broad UV spectral bands. On the other hand, larger absorption cross-sections of the UV bands of the radicals in comparison with mid-infrared bands, and also easy availability of intense and narrow bandwidth UV light source are the primary factors for the popularity of the UV-based spectroscopic methods. We would also like to make a remark that for having only a single pass of optical path (path length  $\sim 100$  cm), the detection sensitivity of this apparatus appears smaller compared to other much used methods, e.g. cavity ring down spectroscopy (CRDS), and partly also due to relatively lower absorption cross-sections of radicals in mid-IR spectral range. For methylperoxy radical, the detection limit of our apparatus has been shown to be  $\sim 2 \times 10^{14} \text{ cm}^{-3}$ . On the other hand, typical optical path length for cavity ring down spectroscopy (CRDS) is  $\sim 1$  km, which results in detection limit of  $\sim 10^{11}-10^{12} \text{ cm}^{-3}$ . The laser induced fluorescence (LIF) or mass spectrometric techniques also have similar detection limit. In near future, we plan to replace the linear reactor by a Herriott type multi-pass absorption cell, which will enable us in detecting the peroxy radicals with concentration as low as  $\sim 5 \times 10^{12} \text{ cm}^{-3}$ . Furthermore, we intend to investigate the reactions of NO and NO<sub>2</sub> with other atmospherically important peroxy radicals.

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