

## Vibrational analysis of the isotopic methylammonium ions

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MS received 8 May 1978; revised 22 September 1978

**Abstract.** A complete vibrational analysis of the methyl ammonium ion  $\text{CH}_3\text{NH}_3^+$ , and of  $\text{CD}_3\text{NH}_3^+$ ,  $\text{CH}_3\text{ND}_3^+$  and  $\text{CD}_3\text{ND}_3^+$  is described in this paper. A General valence Force Field has been computed for these ions, belonging to  $C_{3v}$  symmetry, using the fundamental frequencies obtained from the infrared spectra. The mean square amplitudes and shrinkage constants and Coriolis coupling constants have also been calculated for the first time and presented here.

**Keywords.** Symmetry; vibrational frequencies; normal co-ordinate analysis; molecular constants.

### 1. Introduction

The four isotopic methylammonium ions possess  $C_{3v}$  symmetry. Their fundamental vibrations can be classified according to the irreducible representations (species) of the point group  $C_{3v}$  :

$$\Gamma = 5a_1 + a_2 + 6e.$$

The torsional mode  $a_2$  is infrared inactive under  $C_{3v}$  symmetry. Assignments of the infrared-active fundamentals have been made for each of these ions by Oxtan *et al* (1977). These assignments were based on the previous infrared studies Waldron (1953), Cabana and Sandorfy (1962) and Theoret and Sandorfy (1967), of the ion with simple halides. The structural parameters, vibrational frequencies and the orthonormal set of symmetry co-ordinates used in the present calculations have been taken from the investigation of Oxtan *et al* (1977). The set of internal co-ordinates chosen, numbering of atoms and the orientation of Cartesian co-ordinates axes of methylammonium ion are shown in figure 1. All the angles were assumed to be tetrahedral.

Normal coordinate analysis has been performed by Oxtan *et al* (1977). However, for continuity, the General Valence Force Field (GVFF) has been repeated in the present investigation. The mean amplitudes of vibration, the generalised mean square amplitudes of vibration and shrinkage constant and Coriolis coupling coefficients have been computed for the first time and reported here.

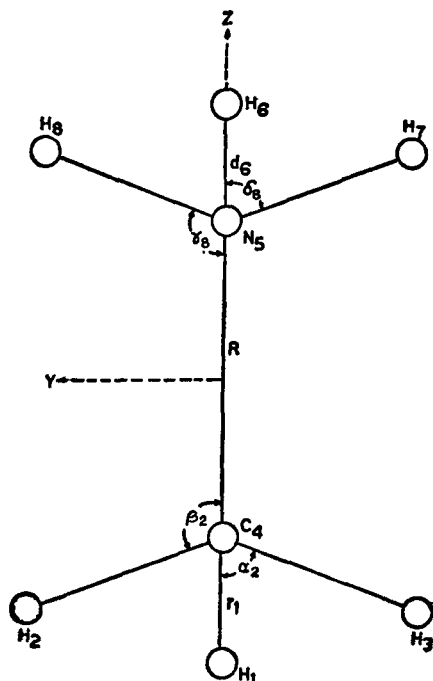


Figure 1. Internal co-ordinates, numbering of atoms and orientation of the principal axes in the isotopic methyl ammonium ions.

## 2. Theoretical considerations

### 2.1. Molecular force field

GVFF is used to obtain the force constants. A reliable set of force constants has been obtained by Wilson's (1939) F-G Matrix method.

### 2.2. Mean-square amplitudes of vibration and shrinkage effect

The symmetrised mean square amplitude matrix  $\Sigma$  has been evaluated using the relation (Cyvin 1960).

$$\Sigma = L \Delta L',$$

where  $L$ 's are characteristic vectors and  $\Delta$  is a diagonal matrix consisting of the mean square values of the normal co-ordinates given by (Block 1932). The generalised mean square amplitudes (Morino and Hirota 1955)  $\langle\langle(\Delta z)^2\rangle\rangle$ ,  $\langle\langle(\Delta y)^2\rangle\rangle$  and  $\langle\langle(\Delta x)^2\rangle\rangle$  and the mean cross products  $\langle\langle(\Delta x)(\Delta y)\rangle\rangle$ ,  $\langle\langle(\Delta y)(\Delta z)\rangle\rangle$  and  $\langle\langle(\Delta z)(\Delta x)\rangle\rangle$  have been evaluated using the  $\Sigma$  matrix from the relation,

$$X = A S \text{ where}$$

$$A = M^{-1} B' G^{-1}.$$

Table 1. Internal force constants of methylammonium ion (m. dyne/Å).

$f_R$	4.6386	$f_{rr}$	0.0416
$f_r$	4.7696	$f_{r\delta}$	-0.0001
$f_d$	5.7320	$f_{d\alpha}$	0.0122
$f_{rr}$	0.0466	$f_{d\beta}$	0.0198
$f_{dd}$	0.0890	$f_{d\gamma}$	0.0697
$f_\alpha - f_{\alpha\alpha}$	0.3581	$f_{d\delta}$	0.0529
$f_\beta - f_{\beta\beta}$	0.4340	$f_{\alpha\beta}$	-0.0918
$f_\gamma - f_{\gamma\gamma}$	0.6521	$f_{\alpha\gamma}$	0.0467
$f_\delta - f_{\delta\delta}$	0.5373	$f_{\alpha\delta}$	0.0020
$f_{Rr}$	0.1002	$f_{\beta\gamma}$	-0.1016
$f_{Rd}$	0.0947	$f_{\beta\delta}$	0.0217
$f_{rd}$	-0.0076	$f_{\gamma\delta}$	0.0287
$f_{r\alpha}$	0.0247	$(1.258 f_{R\alpha} - 1.191 f_{R\beta})$	= -0.4229
$f_{r\beta}$	0.0366	$(1.258 f_{R\beta} - 1.191 f_{R\gamma})$	= -0.3160

Table 2. Mean square amplitudes of vibration of the isotopic methylammonium ions at 300° K in  $10^{-4}$  Å<sup>2</sup>.

	CH <sub>3</sub> NH <sub>3</sub> <sup>+</sup>	CH <sub>3</sub> ND <sub>3</sub> <sup>+</sup>	CD <sub>3</sub> NH <sub>3</sub> <sup>+</sup>	CD <sub>3</sub> ND <sub>3</sub> <sup>+</sup>
$\sigma_R$	25.35	25.16	25.37	25.22
$\sigma_r$	61.46	61.46	44.87	44.64
$\sigma_d$	55.80	40.46	55.74	40.40
$\sigma_{rr}$	-1.13	-1.13	-1.41	-1.29
$\sigma_{dd}$	-1.10	-0.81	-1.11	-0.90
$\sigma_\alpha - \sigma_{\alpha\alpha}$	362.40	413.00	262.20	280.60
$\sigma_\beta - \sigma_{\beta\beta}$	267.90	310.50	196.50	222.80
$\sigma_\gamma - \sigma_{\gamma\gamma}$	208.20	134.80	228.40	157.90
$\sigma_\delta - \sigma_{\delta\delta}$	287.40	206.90	289.00	208.40
$\sigma_{Rr}$	-1.55	-1.55	-2.14	-2.14
$\sigma_{Rd}$	-1.20	-1.71	-1.20	-1.70
$\sigma_{rd}$	0	0	0	0
$\sigma_{r\alpha}$	-7.15	-7.15	-10.75	-10.67
$\sigma_{r\beta}$	-8.14	-8.14	-10.93	-10.85
$\sigma_{r\gamma}$	-5.77	-5.77	-7.75	-7.70
$\sigma_{r\delta}$	0	0	0	0
$\sigma_{d\alpha}$	0	0	0	0
$\sigma_{\alpha\beta}$	-4.39	-5.86	-4.39	-5.87
$\sigma_{d\gamma}$	-6.53	-8.71	-6.52	-8.71
$\sigma_{\alpha\delta}$	-5.43	-8.29	-5.43	-8.30
$\sigma_{\alpha\beta}$	16.33	62.03	18.35	40.48
$\sigma_{\alpha\gamma}$	-28.94	-38.68	-33.00	-39.00
$\sigma_{\alpha\delta}$	0	0	0	0
$\sigma_{\beta\gamma}$	42.66	50.20	51.51	57.63
$\sigma_{\beta\delta}$	-12.26	-15.94	-12.33	-16.03
$\sigma_{\gamma\delta}$	-27.16	-17.68	-21.74	-13.07
$[1.258\sigma_{R\alpha} - 1.191\sigma_{R\beta}]$	27.21	29.10	27.70	28.29
$[1.258\sigma_{R\beta} - 1.191\sigma_{R\gamma}]$	20.25	20.89	20.48	21.01

Table 3. Generalised mean square amplitudes of vibration in  $10^{-4} \text{ \AA}^2$  and shrinkage constants in  $\text{ \AA}$  of the isotopic methylammonium ions.

Ion	Atom pair	Parallel		Perpendicular		Mean cross products			Shrinkages
		$\langle(\Delta z)^2\rangle$	$\langle(\Delta x)^2\rangle$	$\langle(\Delta y)^2\rangle$	$\langle(\Delta x)(\Delta y)\rangle$	$\langle(\Delta y)(\Delta z)\rangle$	$\langle(\Delta x)(\Delta z)\rangle$		
$\text{CH}_3 \text{NH}_3^+$	C-N	26.74	8.71	8.71	0	0	0	0	—
	C-H <sub>1</sub>	61.43	145.00	114.90	0	0	0	-2.12	—
	N-H <sub>6</sub>	55.01	101.11	82.79	0	0	0	34.20	—
	H <sub>1</sub> ...H <sub>5</sub>	165.79	266.92	99.73	-9.21	3.69	0	0.16	0.0081
	H <sub>6</sub> ...H <sub>5</sub>	144.73	218.75	84.79	9.45	0	0	0.02	0.0061
	C...H <sub>6</sub>	104.54	82.07	69.09	0	0	0	-37.35	0.0021
	N...H <sub>1</sub>	115.00	85.31	35.41	0	0	0	41.34	0.0004
	N <sub>1</sub> ...H <sub>6</sub>	250.97	125.87	146.66	0	0	0	-2.66	-0.0417
$\text{CH}_3 \text{ND}_3^+$	C-N	25.15	11.81	11.81	0	0	0	0	—
	C-H <sub>1</sub>	61.41	159.02	139.99	0	0	0	-2.22	—
	N-D <sub>6</sub>	40.45	66.16	57.65	0	0	0	26.24	—
	H <sub>1</sub> ...D <sub>5</sub>	180.38	304.23	114.08	-8.31	-0.01	0	0.01	0.0039
	D <sub>6</sub> ...D <sub>5</sub>	101.17	120.02	58.67	3.49	0	0	0.01	0.0064
	C...D <sub>6</sub>	71.57	51.19	46.20	0	0	0	-19.25	0.0093
	N...H <sub>1</sub>	123.27	91.57	103.20	0	0	0	50.02	-0.0038
	H <sub>1</sub> ...D <sub>6</sub>	271.73	91.76	145.25	0	0	0	45.87	-0.0560
$\text{CD}_3 \text{NH}_3^+$	C-N	25.34	13.24	13.24	0	0	0	0	—
	C-D <sub>1</sub>	44.79	96.45	71.27	0	0	0	-2.61	—
	N-H <sub>6</sub>	55.70	104.24	85.51	0	0	0	35.41	—
	D <sub>1</sub> ...D <sub>3</sub>	117.20	166.52	71.48	-2.64	-0.01	0	0.01	0.0065
	H <sub>6</sub> ...H <sub>6</sub>	145.19	243.12	85.21	9.97	0	0	0.03	0.0052
	C...H <sub>6</sub>	101.67	88.76	75.53	0	0	0	-45.83	0.0022
	N...D <sub>1</sub>	83.90	48.85	51.18	0	0	0	25.27	0.0050
	D <sub>1</sub> ...H <sub>6</sub>	209.51	90.23	112.75	0	0	0	-11.60	-0.0372
$\text{CD}_3 \text{ND}_3^+$	C-N	25.26	16.78	16.78	0	0	0	0	—
	C-D <sub>1</sub>	45.23	104.32	92.56	0	0	0	-2.94	—
	N-D <sub>6</sub>	40.37	69.31	60.02	0	0	0	26.15	—
	D <sub>1</sub> ...D <sub>3</sub>	117.86	186.30	72.33	-1.83	-0.04	0	0.01	0.0095
	D <sub>6</sub> ...D <sub>6</sub>	101.62	143.77	59.37	4.06	-0.02	0	-0.05	0.0075
	C...D <sub>6</sub>	77.62	60.29	51.14	0	0	0	-26.36	0.0068
	N...D <sub>1</sub>	89.32	65.92	57.93	0	0	0	37.18	0.0032
	D <sub>1</sub> ...D <sub>6</sub>	173.98	72.86	90.92	0	0	0	-0.54	-0.0295

### 2.3. Coriolis coupling coefficients

The active couplings for these ions are  $a_1 \times e$  and  $e \times e$ , belonging to  $\xi^x$  and  $\xi^z$  respectively. The Coriolis coupling coefficients have been calculated from the relation given by Meal and Polo (1956)

$$\xi^a = L^{-1} C^a (L^{-1})',$$

where  $C^a$  matrix is obtained from the geometry and atomic masses.

### 3. Results and discussion

The internal force constants for the methylammonium ion using GVFF are presented in table 1. These values compare very well with those calculated by Oxtan *et al* (1977). In general, the important stretching and bending force constants are found to be the same for all the four isotopic ions, as expected.

The mean-square amplitude quantities ( $\sigma$ ), calculated at 300° K are listed in table 2. The generalised mean-square amplitudes of vibration and the shrinkage constants are given in table 3. The parallel mean-square amplitudes  $\langle(\Delta z)^2\rangle$  compare very well with the mean-square amplitude quantities ( $\sigma$ ). Also for C—H and N—H bonds, the perpendicular mean square amplitudes of vibration are found to be greater than the parallel mean-square amplitudes.

The significant Coriolis coupling coefficients for  $a_1 \times e$  and  $e \times e$  couplings are reported in table 4. The following rules are found to be satisfied for the various couplings:

*a* × *e* coupling

$$\zeta_{a_1 \times e_a}^x = \zeta_{a_1 \times e_b}^y = 0$$

$$\zeta_{a_1 \times e_b}^x = -\zeta_{a_1 \times e_a}^y$$

*e* × *e* coupling

$$\sum_{i=7}^{12} \zeta_{ii}^z = I_A/2I_B$$

where  $I_A$  and  $I_B$  are the principal moments of inertia.

The Coriolis coupling coefficients facilitate the analysis of vibration-rotation spectra.

### 4. Conclusions

It is expected that the analysis presented here would help us in knowing the spectroscopic properties of the ion in an environment of lower symmetry, which is

Table 4. Coriolis coupling coefficients of the isotopic methylammonium ions.

	$\text{CH}_3\text{NH}_3^+$	$\text{CH}_3\text{ND}_3^+$	$\text{CD}_3\text{NH}_3^+$	$\text{CD}_3\text{ND}_3^+$
Coupling $a_1 \times e$				
$\zeta_{1,11}^x$	- 0.6385	- 0.6263	- 0.6256	- 0.6299
$\zeta_{2,12}^x$	0.6054	0.5173	0.6105	0.5229
$\zeta_{4,7}^x$	- 0.7350	0.7213	0.7269	- 0.7152
$\zeta_{5,8}^x$	0.7330	0.6920	0.7308	0.6890
$\zeta_{5,10}^x$	0.5536	0.4840	0.5549	0.4941
Coupling $e \times e$				
$\zeta_{7,7}^z$	0.1015	0.1015	0.1844	0.1844
$\zeta_{7,9}^z$	0.6448	0.6172	0.6220	0.6225
$\zeta_{7,11}^z$	- 0.6877	- 0.7144	- 0.6975	- 0.6946
$\zeta_{8,8}^z$	0.0883	0.1624	0.0883	0.1624
$\zeta_{8,10}^z$	0.8030	0.8345	0.8079	0.8450
$\zeta_{8,12}^z$	- 0.4971	- 0.4262	- 0.4898	- 0.4064
$\zeta_{9,9}^z$	- 0.3294	- 0.3178	- 0.3259	- 0.3307
$\zeta_{10,10}^z$	- 0.3267	- 0.3541	- 0.3265	- 0.3511
$\zeta_{11,11}^z$	0.2907	0.2807	0.2519	0.2494
$\zeta_{12,12}^z$	0.2962	0.2906	0.2808	0.2633

Other couplings are found to be insignificant.

a subject of recent interest. The mean square amplitudes of vibration, the generalised mean square amplitudes and shrinkage constants and Coriolis coupling constants have also been calculated for the first time and presented here. The mean amplitude calculations are useful in the interpretation of electron diffraction studies in the molecular structure determinations and the shrinkage constants are helpful in the refinement of bond lengths obtained experimentally. The Coriolis constants are used in the interpretation of vibrational spectra of molecules.

## References

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